On the distinguishability of random quantum states

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Distinguishing quantum states This talk

Distinguishing quantum states

Question

Consider a known ensemble \mathcal{E} of *n* quantum states $\{|\psi_i\rangle\}$ with a priori probabilities p_i . Given an unknown state $|\psi_i\rangle$, picked at random from \mathcal{E} , what is the optimal probability P^{opt} of identifying $|\psi_i\rangle$? That is,

$$P^{opt} = \max_{M} \sum_{i} p_i \langle \psi_i | M_i | \psi_i
angle$$

where we maximise over all POVMs $M = \{M_i\}$.

- Considered by many authors under titles like "quantum hypothesis testing", "quantum detection", etc.
- In general, producing an analytic expression for *P*^{opt} appears to be intractable (although good numerical solutions can be found)

Introduction

Two lower bounds Random quantum states Oracle identification Summary

Distinguishing quantum states This talk

This talk

I will discuss:

- Two analytic lower bounds recently obtained for this optimal probability.
- The application of one of them to distinguishing *random* quantum states.
- An application to the "oracle identification problem" in quantum computation.

Methods The pairwise inner product bound The eigenvalue bound Comparison with previous bounds

Methods

The lower bounds are obtained by putting a lower bound on the probability of success of a specific measurement that can be defined for any ensemble of states, the *Pretty Good Measurement* (PGM). Set $\rho = \sum_{i} p_i |\psi_i\rangle \langle \psi_i |$. Then the PGM is defined by the set of measurement operators $\{|\mu_i\rangle \langle \mu_i|\}$, where $|\mu_i\rangle = \sqrt{p_i} \rho^{-1/2} |\psi_i\rangle$.

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Key fact

Let *G* be the rescaled Gram matrix of the ensemble \mathcal{E} , $G_{ij} = \sqrt{p_i p_j} \langle \psi_i | \psi_j \rangle$. Then the probability of success of the PGM is

$$P^{pgm}(\mathcal{E}) = \sum_{i} p_i |\langle \psi_i | \mu_i \rangle|^2 = \sum_{i} (\sqrt{G})_{ii}^2$$

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The pairwise inner product bound

- The first lower bound is based on the pairwise distinguishability of the states in \mathcal{E} .
- The strategy is to put a lower bound on the square root function by an "easier" function (a parabola), and then optimise the parabola.
- Works because $\sqrt{x} \ge ax + bx^2 \Rightarrow (\sqrt{G})_{ii} \ge aG_{ii} + b\sum_j |G_{ij}|^2$.

Pairwise inner product bound

Let \mathcal{E} be an ensemble of *n* states $\{|\psi_i\rangle\}$ with a priori probabilities p_i .

Then
$$P^{pgm}(\mathcal{E}) \ge \sum_{i=1}^{n} \frac{p_i^2}{\sum_{j=1}^{n} p_j |\langle \psi_i | \psi_j \rangle|^2}$$

Methods The pairwise inner product bound **The eigenvalue bound** Comparison with previous bounds

The eigenvalue bound

- The second lower bound is based on a global measure of distinguishability of the states in \mathcal{E} : the eigenvalues of the Gram matrix *G*.
- Using a Cauchy-Schwarz inequality, we can show the following:

Eigenvalue bound

Let *G* be the Gram matrix of an ensemble \mathcal{E} of *n* states and let *G* have eigenvalues $\{\lambda_i\}$. Then

$$P^{pgm}(\mathcal{E}) \ge \frac{1}{n} \left(\sum_{i} \sqrt{\lambda_i}\right)^2 = \frac{1}{n} \operatorname{tr}(\sqrt{G})^2$$

Methods The pairwise inner product bound The eigenvalue bound **Comparison with previous bounds**

Comparison with previous bounds

- Previous authors (e.g. Burnashev and Holevo¹) have used bounds based on similar principles.
- But the bounds here are stronger, especially for low values of $P^{pgm}(\mathcal{E})$, and always give a non-trivial value.

¹M. V. Burnashev and A. S. Holevo, On reliability function of quantum communication channel, quant-ph/9703013

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Comparison with previous bounds

- Previous authors (e.g. Burnashev and Holevo¹) have used bounds based on similar principles.
- But the bounds here are stronger, especially for low values of $P^{pgm}(\mathcal{E})$, and always give a non-trivial value.
- Assuming the states in \mathcal{E} have equal probabilities:

Comparison of bounds

Previously known lower bound $P^{pgm}(\mathcal{E}) \ge 1 - \frac{1}{n} \sum_{i \ne j} |\langle \psi_i | \psi_j \rangle|^2$ $P^{pgm}(\mathcal{E}) \ge \frac{2}{\sqrt{n}} \operatorname{tr}(\sqrt{G}) - 1$

 $\begin{array}{c} \textbf{New lower bound} \\ P^{pgm}(\mathcal{E}) \geq \frac{1}{n} \sum_{i=1}^{n} \frac{1}{\sum_{j=1}^{n} |\langle \psi_i | \psi_j \rangle|^2} \\ P^{pgm}(\mathcal{E}) \geq \frac{1}{n} \mathrm{tr}(\sqrt{G})^2 \end{array}$

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Random quantum states The Marčenko-Pastur law Technical issues The finished lower bound Comparison with numerical results

Random quantum states

- What is an ensemble of random quantum states?
- Here, we mean a set of *n d*-dimensional pure states whose components (in some basis) are i.i.d. complex random variables with mean 0 and variance 1/d.
- This is a quite general notion of randomness that includes pure states distributed uniformly at random (according to Haar measure), in which case the components (in any basis!) are Gaussians.

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- This is a quite general notion of randomness that includes pure states distributed uniformly at random (according to Haar measure), in which case the components (in any basis!) are Gaussians.
- The pairwise inner product bound (above) can be applied to random quantum states directly, but we can get better results from the eigenvalue bound.
- In order to apply this bound, we need a powerful result from random matrix theory.

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The Marčenko-Pastur law

- If the states in \mathcal{E} are random and $p_i = 1/n$ for all *i*, the Gram matrix *G* is known to statisticians (since the 1930s!) as a rescaled complex *Wishart matrix*.
- The density of the eigenvalues of *G* is known and is given by the Marčenko-Pastur law.
 - This is the equivalent of the famous Wigner semicircle law for random Hermitian matrices...
- This allows us to calculate a lower bound on the expected probability of success for the PGM!

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Technical issues

- The Marčenko-Pastur law can be applied to random states in the asymptotic regime where:
 - The number of states *n* and the dimension *d* approach infinity.
 - The ratio n/d approaches a constant, r.
- We need to modify the Marčenko-Pastur law slightly.
 - It gives the density of the eigenvalues of the Gram matrix; we need the density of the square roots of the eigenvalues.
- The lower bound we get for $\mathbb{E}(P^{pgm}(\mathcal{E}))$ turns out to be given by an intractable elliptic integral.
- However, a good lower bound may be proven on this integral, giving the main result...

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The finished lower bound

Main theorem

Let \mathcal{E} be an ensemble of *n* equiprobable *d*-dimensional quantum states $\{|\psi_i\rangle\}$ with $n/d \to r \in (0, \infty)$ as $n, d \to \infty$, and let the components of $|\psi_i\rangle$ in some basis be i.i.d. complex random variables with mean 0 and variance 1/d. Then

$$\mathbb{E}(P^{pgm}(\mathcal{E})) \ge \begin{cases} \frac{1}{r} \left(1 - \frac{1}{r} \left(1 - \frac{64}{9\pi^2}\right)\right) & \text{if } n \ge d\\ 1 - r \left(1 - \frac{64}{9\pi^2}\right) & \text{otherwise} \end{cases}$$

and in particular $\mathbb{E}(P^{pgm}(\mathcal{E})) > 0.720$ when $n \leq d$.

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• Concentration of measure results may be used to show that almost all states obey this lower bound!

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Comparison with numerical results (1) $(0 \le n \le 2d)$



Figure: Asymptotic bound on $P^{pgm}(\mathcal{E})$ vs. numerical results (averaged over 10 runs) for ensembles of n = 50r 50-dimensional uniformly random states.

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Comparison with numerical results (2) $(0 \le n \le 10d)$



Figure: Asymptotic bound on $P^{pgm}(\mathcal{E})$ vs. numerical results (averaged over 10 runs) for ensembles of n = 50r 50-dimensional uniformly random states.

Oracle identification

Oracle identification

Problem

Given an unknown Boolean function f, picked uniformly at random from a set S of N Boolean functions on n bits, identify f with the minimum number of uses of f.

- This is a particular case of the oracle identification problem studied by Ambainis et al².
- We consider the case where we must identify *f* with a bounded probability of error.

²A. Ambainis et al, Quantum identification of Boolean oracles, quant-ph/0403056

Oracle identification

Oracle identification

- Consider the following single-query "algorithm":
 - Create the state $|\psi_f\rangle = \sum_x (-1)^{f(x)} |x\rangle$.
 - Apply the PGM.
- When S is a random set of functions, the states $\{|\psi_f\rangle\}$ are random quantum states.
- So the results here can be used to put the same lower bound on the probability of success of distinguishing these states.
- Concentration of measure is used again to show that this bound holds for almost all sets of functions.
- When the probability of success is a constant > 1/2, we can repeat the algorithm a constant number of times for an arbitrarily good probability of success.

Summary

- Good lower bounds have been obtained on the probability of distinguishing pure quantum states.
- These bounds can be applied to distinguishing random quantum states.
- Asymptotically, *n* random states in *n* dimensions can be distinguished with probability > 0.72.
- Almost all sets of 2ⁿ Boolean functions on *n* bits can be distinguished with a constant number of quantum queries.

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- Asymptotically, *n* random states in *n* dimensions can be distinguished with probability > 0.72.
- Almost all sets of 2ⁿ Boolean functions on *n* bits can be distinguished with a constant number of quantum queries.
- Further reading: quant-ph/0607011
- Thank you for your time!