



The distribution of the number of isolated nodes in the 1-Dimensional soft random geometric graph

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ABSTRACT

We study the number of isolated nodes in a soft random geometric graph whose vertices constitute a Poisson process on the torus of length L (the line segment $[0, L]$ with periodic boundary conditions), and where an edge is present between two nodes with a probability which depends on the distance between them. Edges between distinct pairs of nodes are mutually independent. In a suitable scaling regime, we show that the number of isolated nodes converges in total variation to a Poisson random variable. The result implies an upper bound on the probability that the random graph is connected.

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1. Introduction

Random geometric graphs (RGGs) were introduced in Gilbert (1961) as a model for communication networks with limited connection range and have subsequently been widely used to model networks with a spatial element; see, e.g., Penrose (2003) and references therein. In Gilbert's model, nodes or vertices are randomly distributed within some space, typically \mathbb{R}^d or a bounded subset of it, and an edge is placed between nodes if their mutual distance is smaller than a specified threshold. We shall refer to this model as the hard RGG, as the connection probability has a hard cutoff as a function of the distance between nodes. A generalisation is to allow the edge between a pair of nodes separated by distance r to be present with some probability $H(r)$, independent of all other edges. This model has been termed the Soft RGG (Penrose, 2016), the random connection model (RCM) (Meester and Roy, 1996), or a Waxman graph (Waxman, 1988). In this work we refer to them as Soft RGGs and term $H(\cdot)$ the connection function. The hard RGG is a special case, obtained by setting $H(r) = 1$ for $r \leq r_c$ and $H(r) = 0$ otherwise; $r_c > 0$ is a parameter of the model.

The one-dimensional version of this model is motivated by vehicular ad-hoc networks (VANETs), which are expected to be essential for autonomous vehicles; these will be fitted with on-board radios to transmit information such as location, velocity, and hazard warnings between vehicles. The road is represented by a line, the nodes represent vehicles and an edge between two nodes indicates that the two vehicles can communicate directly. One key question is connectivity: When is every vehicle in the network able to communicate with every other, either via a single- or multi-hop path? A necessary condition for full connectivity is the absence of isolated nodes, namely nodes that are not connected to any other in the graph. In the 2-D and 3-D versions of the soft RGG model, it has been shown in a suitable asymptotic regime that the soft RGG is connected if and only if there are no isolated nodes (Mao and Anderson, 2012; Penrose, 2016). It was further shown in Mao and Anderson (2012) that the number of isolated nodes in a 2-D soft RGG can be well approximated by a Poisson distribution. The result was extended in Penrose (2016) to dimension three and greater; also see Ganesh and Xue (2007) for an analogous result for small-world networks. There has been little work to date on the 1-D model, which

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differs in important respects from those in two or more dimensions. In particular, the dominant reason for disconnection in 1-D hard RGGs is the presence of uncrossed gaps rather than of isolated nodes. Soft RGGs in 1-D were studied in [Wilsher et al. \(2020\)](#), where it was shown that isolated nodes dominate uncrossed gaps as a cause of disconnection. The threshold for the emergence of isolated nodes was established in [Wilsher et al. \(2020\)](#), and a Poisson limit was conjectured for the number of isolated nodes at the threshold (Conjecture 4.1). We prove this conjecture here.

We now describe the model studied in this paper and our assumptions. We consider a sequence of networks indexed by a parameter $L \in \mathbb{R}_+$, which tends to infinity along the sequence. The nodes or vertices of the network constitute a Poisson point process (PPP) of unit intensity on $[0, L]$, which we denote \mathcal{P}_L . If two nodes are located a distance r apart, the edge between them is present with probability $H^L(r) = H(r/R_L)$, independent of all other edges; here, $H : \mathbb{R}^+ \rightarrow [0, 1]$ is a given connection function, and $R_L > 0$ is a scaling parameter to be specified.

In order to avoid inessential technicalities associated with the boundaries, we shall study the model with periodic boundary conditions. In other words, we turn $[0, L]$ into a circle or torus by identifying 0 with L . We denote the circular distance on $[0, L]$ by ρ^L , i.e., $\rho^L(x, y) = \min\{|x - y|, L - |x - y|\}$. Finally, we define the connection probabilities to be $h^L(x, y) = H^L(\rho^L(x, y))$. We denote by $\mathcal{G}_{h^L}(\mathcal{P}_L)$ the graph with vertex set \mathcal{P} , and independent edges generated with connection probabilities h^L . This is the graph model we study. We believe that the same results hold if $[0, L]$ is treated as a line segment rather than a circle, and that this can be established by analysing boundary effects separately, as in [Dettmann and Georgiou \(2016\)](#).

We make the following assumptions about the connection function.

Assumptions. Let $\|H\|_1$ and $\|H\|_2$ denote the L^1 and L^2 norms of $H : \mathbb{R}_+ \rightarrow [0, 1]$, i.e.,

$$\|H\|_1 = \int_0^\infty H(x)dx, \quad \|H\|_2^2 = \int_0^\infty H(x)^2 dx.$$

We assume that $\|H\|_1 < \infty$ and $\|H\|_2^2 < \|H\|_1$.

The first assumption says that H is integrable, and is required for the mean number of neighbours of a node to remain bounded as L tends to infinity. We expect this to hold in real-world networks. Next, observe from the definition that $\|H\|_2^2 \leq \|H\|_1$, since $H(x) \in [0, 1]$ for all x as it is a probability. Hence, the second assumption, which asserts that this inequality is strict, is a mild one. It is satisfied whenever the set, $\{x \in \mathbb{R}_+ : 0 < H(x) < 1\}$, where the connection probability is strictly between 0 and 1, has positive Lebesgue measure. Nevertheless, the assumption excludes the connection function of a hard RGG, which is $\{0, 1\}$ -valued.

We now state our main result, which resolves a conjecture in [Wilsher et al. \(2020\)](#).

Theorem 1. Fix $\tau \in \mathbb{R}_+$ and take $R_L = \ln(\tau L)/2\|H\|_1$. Let N_{iso} denote the number of isolated nodes in the soft RGG $\mathcal{G}_{h^L}(\mathcal{P}_L)$. Its dependence on L and τ has been suppressed for notational convenience. As L tends to infinity, the random variable N_{iso} converges in total variation distance to a Poisson random variable with mean $1/\tau$. In particular, $\mathbb{P}(N_{iso} = 0)$ tends to $e^{-1/\tau}$.

As noted above, the assumption of the theorem excludes the hard RGG model. This appears to be an artifact of the proof technique, where the assumption is needed (only) for the proof of [Lemma 6](#), which bounds the correlation of isolation between nearby nodes. The result in the theorem is already known for the hard RGG model ([Penrose, 2003](#)), in which the connection region is of the form $[-r, r]$; it is not known if the connection region is of the general form $[a, b]$ or if it is a disconnected set, such as a union of countably many intervals. An approach to extending the analysis to these cases might involved bounding the overlap between the connection regions of nodes at different locations.

Before turning to the proof of the theorem, we compare our model and results in greater detail with those in [Penrose \(2016\)](#), as our approach is very similar. Note that our model has 3 parameters: the length L of the domain, the intensity of the Poisson process (taken to be 1) and the scale R_L of the connection function. These have dimensions of length or its inverse, so that there are only two dimensionless parameters in the model (which is invariant to changes in the units in which length is measured). The scaling regime in [Penrose \(2016\)](#) is superficially different as it keeps the domain length fixed, while scaling the connection range and the intensity of the Poisson process. In fact, these scalings are equivalent as they use the same two dimensionless parameters.

The analysis in [Penrose \(2016\)](#) is restricted to dimensions $d \geq 2$, and makes strong assumptions about the connection function such as that the connection probability either decays as a stretched exponential or has compact support. These restrictions are needed in order to show that connectivity of the random graph is equivalent to the absence of isolated vertices, with high probability. As connectivity is the primary object of interest, that paper does not seek Poisson approximations for the number of isolated nodes under weaker assumptions. Here, we are able to establish a Poisson approximation under much weaker assumptions. The analysis should generalise easily to dimensions $d \geq 2$, and is indeed adapted from that context in [Penrose \(2016\)](#), but we have not attempted to provide a detailed generalisation within this short note.

As the soft RGG is disconnected if there are any isolated nodes, an immediate corollary of the theorem is that the probability that $\mathcal{G}_{h^L}(\mathcal{P}_L)$ is connected is asymptotically bounded above by $e^{-1/\tau}$, in the scaling regime considered in the theorem. This upper bound would be tight if isolated nodes were the dominant cause of disconnection in this random graph model, as conjectured in [Wilsher et al. \(2020\)](#) under the mild additional condition that the connection function has

unbounded support. Resolving this conjecture remains an open problem, as does extending the analysis of this paper to point process models other than the Poisson process. Unfortunately, our analysis in this paper only pertains to isolated nodes, and does not shed light on other possible ways in which the graph can be disconnected.

2. Proofs

Denote the total variation distance between probability distributions μ and ν on \mathbb{R} by $d_{TV}(\mu, \nu)$. With a slight abuse of notation, we shall write $d_{TV}(X, Y)$ for random variables X and Y to denote the total variation distance between their probability distributions. The following bound on the total variation distance between random variables X and Y defined on the same probability space is well-known and elementary:

$$d_{TV}(X, Y) \leq \mathbb{P}(X \neq Y). \tag{1}$$

The proof of [Theorem 1](#) proceeds through a sequence of lemmas. Our first result approximates the number of isolated nodes in our model with one in which the connection function is truncated. This step is not needed if the connection function has bounded support.

Lemma 1. Fix $\alpha > 0$, and define the truncated connection function $\tilde{h}^L(x, y) = h^L(x, y)\mathbb{1}\{\rho^L(x, y) \leq R_L^{1+1/\alpha}\}$. Denote by N_{iso} and \tilde{N}_{iso} the number of isolated nodes in $\mathcal{G}_{h^L}(\mathcal{P}_L)$ and $\mathcal{G}_{\tilde{h}^L}(\mathcal{P}_L)$ respectively. Then,

$$\lim_{L \rightarrow \infty} d_{TV}(N_{iso}, \tilde{N}_{iso}) = 0.$$

Proof. We can couple $\mathcal{G}_{h^L}(\mathcal{P}_L)$ and $\mathcal{G}_{\tilde{h}^L}(\mathcal{P}_L)$ by first generating $\mathcal{G}_{h^L}(\mathcal{P}_L)$, and then removing any edges of length at least $R_L^{1+1/\alpha}$. Observe that $\tilde{N}_{iso} \geq N_{iso}$, since removing edges cannot reduce the number of isolated nodes. Therefore, it follows from Markov's inequality that

$$\mathbb{P}(|N_{iso} - \tilde{N}_{iso}| \geq 1) \leq \mathbb{E}[|N_{iso} - \tilde{N}_{iso}|] = \mathbb{E}[\tilde{N}_{iso}] - \mathbb{E}[N_{iso}]. \tag{2}$$

Using the expression for the expected number of isolated nodes in a soft RGG derived in [Wilsher et al. \(2020, eqn. \(9\)\)](#), we get

$$\mathbb{E}[N_{iso}] = L \exp\left(-2R_L \int_0^{L/2R_L} H(r)dr\right) = L \exp\left(-\frac{\ln(\tau L)}{\|H\|_1} \int_0^{L/2R_L} H(r)dr\right).$$

Now, $L/2R_L$ tends to infinity as L tends to infinity, and so $\int_0^{L/2R_L} H(r)dr$ tends to $\|H\|_1 = \int_0^\infty H(r)dr$. Hence, it follows from the above that

$$\mathbb{E}[N_{iso}] \xrightarrow{L \rightarrow \infty} \frac{1}{\tau}. \tag{3}$$

Similarly, since $R_L = \frac{\ln(\tau L)}{2\|H\|_1}$ tends to infinity with L , we have

$$\begin{aligned} \mathbb{E}[\tilde{N}_{iso}] &= L \exp\left(-2 \int_0^{L/2} \tilde{h}^L(0, r)dr\right) = L \exp\left(-2 \int_0^{R_L^{1+1/\alpha}} H\left(\frac{r}{R_L}\right)dr\right) \\ &= L \exp\left(-\frac{\ln(\tau L)}{\|H\|_1} \int_0^{R_L^{1/\alpha}} H(r)dr\right) \xrightarrow{L \rightarrow \infty} \frac{1}{\tau}. \end{aligned} \tag{4}$$

It follows from (3) and (4) that $\mathbb{E}[\tilde{N}_{iso}] - \mathbb{E}[N_{iso}]$ tends to zero as L tends to infinity. The claim of the lemma follows from (1) and (2). \square

Henceforth, we shall work with $\mathcal{G}_{\tilde{h}^L}(\mathcal{P}_L)$, the soft RGG with truncated connection function. We shall use the Chen–Stein method for Poisson approximation described in [Barbour et al. \(1992\)](#); as it requires a discrete index set, we fix $m \in \mathbb{N}$ and discretise the torus $[0, L]$ into mL segments of width $1/m$. (Assume for convenience that L is an integer. Otherwise, we need one segment of a different width, which does not fundamentally alter the analysis.) Denote the i th segment by A_i , where i takes values in the index set $\Gamma = \{1, 2, \dots, mL\}$. Denote by \mathcal{P}_L and \mathcal{I} the vertex set and the set of isolated nodes of the random graph $\mathcal{G}_{\tilde{h}^L}(\mathcal{P}_L)$, and by $\mathcal{P}_L(A_i)$ and $\mathcal{I}(A_i)$ the subsets of these nodes lying within A_i . Define

$$\begin{aligned} J_i &:= \mathbb{1}\{|\mathcal{P}_L(A_i)| = 1\}, \\ I_i &:= \mathbb{1}\{|\mathcal{P}_L(A_i)| = 1 \text{ and } |\mathcal{I}(A_i)| = 1\}, \\ W_{m,L} &:= \sum_{i \in \Gamma} I_i, \end{aligned} \tag{5}$$

for $i \in \Gamma$. Denote the centre of the segment A_i by x_i . (Although A_i, I_i, J_i , and x_i all implicitly depend on m , this dependence is suppressed for notational convenience.) Our next result states that the number of isolated nodes is well-approximated by $W_{m,L}$ when m is large.

Lemma 2. Let I_i denote the indicator that the i th segment of $[0, L]$ contains exactly one node, and that node is isolated in $\mathcal{G}_{\tilde{h}_L}(\mathcal{P}_L)$. Let $W_{m,L}$ denote the sum of these indicators, as defined in (5). Let \tilde{N}_{iso} denote the total number of isolated nodes in $\mathcal{G}_{\tilde{h}_L}(\mathcal{P}_L)$. Then, for any fixed $L > 0$,

$$d_{TV}(\tilde{N}_{iso}, W_{m,L}) \rightarrow 0 \text{ as } m \rightarrow \infty.$$

Proof. Notice that the random variable $W_{m,L}$ is exactly the same as the number of isolated nodes unless there is a segment containing two or more nodes. Now, the number of nodes in a segment of length $1/m$ has a Poisson distribution with mean $1/m$. Hence, it follows from (1) and the union bound that

$$\begin{aligned} d_{TV}(\tilde{N}_{iso}, W_{m,L}) &\leq \mathbb{P}(\exists i \in \Gamma : |\mathcal{P}_L(A_i)| \geq 2) \\ &\leq \sum_{i \in \Gamma} \left(1 - e^{-1/m} - \frac{1}{m}e^{-1/m}\right) \\ &= mL \left(1 - e^{-1/m} - \frac{1}{m}e^{-1/m}\right), \end{aligned}$$

which tends to zero as m tends to infinity, as claimed. \square

In light of the above lemma, it suffices to establish a Poisson approximation for the random variables $W_{m,L}$. To this end, we define the ‘‘neighbourhood’’ of an index $i \in \Gamma$ by

$$B_i = \left\{j \in \Gamma : \rho^L(x_i, x_j) \leq 3R_L^{1+1/\alpha} \text{ for some } x_i \in I_i, x_j \in I_j\right\}. \tag{6}$$

We also define the quantities

$$\begin{aligned} p_i &:= \mathbb{E}[I_i], \quad p_{ij} := \mathbb{E}[I_i I_j], \\ b_1 &:= \sum_{i \in \Gamma} \sum_{j \in B_i} p_i p_j, \quad b_2 := \sum_{i \in \Gamma} \sum_{j \in B_i \setminus \{i\}} p_{ij}, \quad b_3 := \sum_{i \in \Gamma} \mathbb{E} \left[|\mathbb{E}[I_i | (I_j : j \in \Gamma \setminus B_i)] - p_i| \right]. \end{aligned} \tag{7}$$

We shall use of the following result on Chen–Stein approximation.

Lemma 3 ((Barbour et al., 1992, Theorem 1.A)). Let $Po(\lambda)$ denote a Poisson random variable with mean λ . Let $W_{m,L}$ be defined as in (5) and b_1, b_2 and b_3 as in (7). Then,

$$d_{TV}(W_{m,L}, Po(\mathbb{E}[W_{m,L}])) \leq \min\left(1, \frac{1}{\mathbb{E}[W_{m,L}]}\right)(b_1 + b_2 + b_3).$$

Thus, in order to establish a Poisson approximation for $W_{m,L}$, we need to bound the terms b_1, b_2, b_3 .

Lemma 4. Let B_i be defined as in (6) and b_1, b_2, b_3 as in (7). Then, $b_3 = 0$ for all m sufficiently large.

Proof. The connection function \tilde{h}_L is defined by truncating h_L at $R_L^{1+1/\alpha}$. Hence, the event for which I_i is the indicator depends only on nodes within distance $R_L^{1+1/\alpha}$ of the segment A_i . By the same reasoning, the isolation or otherwise of nodes at distance greater than $2R_L^{1+1/\alpha}$ from this segment is independent of the Poisson point process in the set $\{x : \min_{y \in A_i} \rho^L(x, y) \leq R_L^{1+1/\alpha}\}$. Hence, I_i is jointly independent of I_j for all j such that $\min_{x \in A_i, y \in A_j} \rho^L(x, y) > 2R_L^{1+1/\alpha}$. In particular, if m is large enough that $\frac{1}{m} < R_L^{1+1/\alpha}$, then I_i is independent of the collection of random variables, $\{I_j, j \in \Gamma \setminus B_i\}$. Hence, $b_3 = 0$ for all m sufficiently large. \square

Lemma 5. Let B_i be defined as in (6) and b_1, b_2, b_3 as in (7). Then, b_1 tends to zero as we let m tend to infinity, followed by L .

Proof. Observe from the definition of I_i that

$$\begin{aligned} p_i &= \mathbb{P}(J_i = 1)\mathbb{P}(I_i = 1 | J_i = 1) \\ &= \frac{1}{m}e^{-1/m}\mathbb{P}(v \in \mathcal{P}_L(A_i) \text{ is isolated in } \mathcal{G}_{\tilde{h}_L}(\mathcal{P}_L) | J_i = 1), \end{aligned} \tag{8}$$

where $v \in \mathcal{P}_L(A_i)$ exists and is unique since $J_i = 1$. Say v is located at $(i/m) + x$, where $0 \leq x < 1/m$. Now, the set of nodes to which v is connected form an inhomogeneous Poisson process on $[0, L] \setminus A_i$, of intensity $\tilde{h}^L(i/m + x, y)$; v is isolated only if this set is empty. Thus, by translation invariance of the connection function, the probability that v is isolated is given by

$$\mathbb{P}(v \text{ is isolated}) = e^{-\kappa x}, \text{ where}$$

$$\kappa_x = \int_{[0,L] \setminus [0,1/m]} \tilde{h}^L(x, y) dy \xrightarrow{m \rightarrow \infty} \int_0^L \tilde{h}^L(x, y) dy,$$

since \tilde{h}^L is bounded above by 1. We have not made the dependence of κ_x on m and L explicit in the notation. Suppose L is large enough that $L/2R_L \geq R_L^{1+1/\alpha}$. Then, using the translation invariance of \tilde{h}^L once more, we have

$$\lim_{m \rightarrow \infty} \kappa_x = \int_0^L \tilde{h}^L(0, y) dy = 2 \int_0^{R_L^{1+1/\alpha}} H\left(\frac{y}{R_L}\right) dy = 2R_L \int_0^{R_L^{1/\alpha}} H(y) dy,$$

which does not depend on x . Substituting in (8), we get that

$$\lim_{m \rightarrow \infty} mp_i = \lim_{m \rightarrow \infty} e^{-1/m} e^{-\kappa_x} = e^{-2R_L \int_0^{R_L^{1/\alpha}} H(y) dy}.$$

Substituting in the definition of b_1 , and noting that p_i does not depend on i , we see that

$$\begin{aligned} \lim_{m \rightarrow \infty} b_1 &= \lim_{m \rightarrow \infty} |\Gamma| |B_i| p_i^2 = \lim_{m \rightarrow \infty} 6m^2 L R_L^{1+1/\alpha} p_i^2 \\ &= 6L \left(\frac{\ln(\tau L)}{2\|H\|_1}\right)^{1+1/\alpha} \exp\left(-\frac{2\ln(\tau L)}{\|H\|_1} \int_0^{R_L^{1/\alpha}} H(y) dy\right). \end{aligned}$$

Now, $\int_0^{R_L^{1/\alpha}} H(y) dy$ tends to $\|H\|_1$ as L , and hence R_L , tends to infinity. So, it follows that

$$\lim_{L \rightarrow \infty} \lim_{m \rightarrow \infty} b_1 = \lim_{L \rightarrow \infty} \frac{6L}{(\tau L)^2} \left(\frac{\ln(\tau L)}{2\|H\|_1}\right)^{1+1/\alpha} = 0.$$

This completes the proof of the lemma. \square

Lemma 6. Let B_i be defined as in (6) and b_2 as in (7). Then, b_2 tends to zero as we let m tend to infinity, followed by L .

Proof. Consider two nodes located at $x, y \in [0, L]$. The set of all other nodes to which at least one of them has an edge constitutes an inhomogeneous PPP on $[0, L]$, with intensity $\phi(\cdot, \{x, y\})$ given by

$$\phi(z, \{x, y\}) = 1 - (1 - \tilde{h}^L(z, x))(1 - \tilde{h}^L(z, y)). \tag{9}$$

Fix $i, j \in \Gamma, i \neq j$, and condition on the event that $J_i = 1$ and $J_j = 1$, i.e., that there is a unique point of the PPP, \mathcal{P}_L , on each of the segments A_i and A_j . Denote their positions by x and y . The set of nodes to which these might be connected, besides each other, constitutes a PPP on $[0, L] \setminus (A_i \cup A_j)$ with intensity ϕ defined above. Hence, the probability that both these nodes are isolated is given by

$$(1 - \tilde{h}^L(x, y)) \exp\left(-\int_{[0,L] \setminus (A_i \cup A_j)} \phi(z, \{x, y\}) dz\right).$$

Now, by well-known properties of Poisson point processes, the unique points of the homogeneous PPP on the segments A_i and A_j are independently and uniformly distributed within them. Hence,

$$\begin{aligned} p_{ij} &= \mathbb{P}(J_i = 1, J_j = 1) \mathbb{P}(I_i = 1, I_j = 1 | J_i = 1, J_j = 1) \\ &= \frac{1}{m^2} e^{-2/m} \int_{x \in A_i} \int_{y \in A_j} m^2 (1 - \tilde{h}^L(x, y)) \exp\left(-\int_{[0,L] \setminus (A_i \cup A_j)} \phi_{\tilde{h}^L}(z, \{x, y\}) dz\right) dx dy \\ &\leq \int_{x \in A_i} \int_{y \in A_j} \exp\left(-\int_{[0,L] \setminus (A_i \cup A_j)} \phi_{\tilde{h}^L}(z, \{x, y\}) dz\right) dx dy. \end{aligned}$$

Substituting in the definition of b_2 , we obtain that

$$b_2 \leq \int_{x \in [0,L]} \int_{y: \rho^L(x,y) \leq 3R_L^{1+1/\alpha}} \exp\left(-\int_{[0,L] \setminus (A_i(x) \cup A_j(y))} \phi_{\tilde{h}^L}(z, \{x, y\}) dz\right) dx dy,$$

where we write $A_i(x)$ and $A_j(y)$ to denote the segments in which x and y lie; allowing them to lie in the same segment yields an upper bound, as does dropping the $(1 - \tilde{h}^L(x, y))$ term. Now,

$$\int_{[0,L] \setminus (A_i(x) \cup A_j(y))} \phi_{\tilde{h}^L}(z, \{x, y\}) dz \xrightarrow{m \rightarrow \infty} \int_{[0,L]} \phi_{\tilde{h}^L}(z, \{x, y\}) dz,$$

since $\phi_{\tilde{h}^L}$ is bounded above by 1. Hence, we conclude using the translation invariance of \tilde{h}^L that

$$\begin{aligned} \limsup_{m \rightarrow \infty} b_2 &\leq \int_{x \in [0, L]} \int_{y: \rho^L(x, y) \leq 3R_L^{1+1/\alpha}} \exp\left(-\int_{[0, L]} \phi_{\tilde{h}^L}(z, \{x, y\}) dz\right) dx dy \\ &= 2L \int_{y=0}^{3R_L^{1+1/\alpha}} \exp\left(-\int_0^L \phi_{\tilde{h}^L}(z, \{0, y\}) dz\right) dy. \end{aligned} \tag{10}$$

Substituting for $\phi_{\tilde{h}^L}$ from (9), we have

$$\int_0^L \phi_{\tilde{h}^L}(z, \{0, y\}) dz = \int_0^L (\tilde{h}^L(0, z) + \tilde{h}^L(y, z) - \tilde{h}^L(0, z)\tilde{h}^L(y, z)) dz. \tag{11}$$

Now, by the Schwarz inequality,

$$\left(\int_0^L \tilde{h}^L(0, z)\tilde{h}^L(y, z) dz\right)^2 \leq \left(\int_0^L (\tilde{h}^L(0, z))^2 dz\right) \left(\int_0^L (\tilde{h}^L(y, z))^2 dz\right). \tag{12}$$

As $\tilde{h}^L(y, \cdot)$ is just a circular shift of $\tilde{h}^L(0, \cdot)$, we also have

$$\int_0^L \tilde{h}^L(y, z) dz = \int_0^L \tilde{h}^L(0, z) dz, \quad \int_0^L (\tilde{h}^L(y, z))^2 dz = \int_0^L (\tilde{h}^L(0, z))^2 dz. \tag{13}$$

Substituting (12) and (13) into (11), we get

$$\int_0^L \phi_{\tilde{h}^L}(z, \{0, y\}) dz \geq 2 \int_0^L \tilde{h}^L(0, z) dz - \int_0^L (\tilde{h}^L(0, z))^2 dz, \tag{14}$$

Now,

$$\begin{aligned} \frac{1}{R_L} \int_0^L \tilde{h}^L(0, z) dz &= \frac{2}{R_L} \int_0^{R_L^{1+1/\alpha}} H\left(\frac{z}{R_L}\right) dz \\ &= 2 \int_0^{R_L^\alpha} H(x) dx \xrightarrow{L \rightarrow \infty} 2\|H\|_1, \end{aligned} \tag{15}$$

since R_L tends to infinity as L does. A similar calculation yields

$$\frac{1}{R_L} \int_0^L (\tilde{h}^L(0, z))^2 dz \xrightarrow{L \rightarrow \infty} 2R_L \|H\|_2^2. \tag{16}$$

Substituting (15) and (16) in (14), we get

$$\liminf_{L \rightarrow \infty} \frac{1}{R_L} \int_0^L \phi_{\tilde{h}^L}(z, \{0, y\}) dz \geq 4\|H\|_1 - 2\|H\|_2^2.$$

Recall that, by assumption, there exists $\delta > 0$ such that $\|H\|_2^2 \leq \|H\|_1 - \delta$. Hence, it follows that

$$\liminf_{L \rightarrow \infty} \frac{1}{R_L} \int_0^L \phi_{\tilde{h}^L}(z, \{0, y\}) dz \geq 2\|H\|_1 + 2\delta.$$

Since $R_L = \ln(\tau L)/2\|H\|_1$, we now obtain from (10) that

$$\begin{aligned} \limsup_{L \rightarrow \infty} \limsup_{m \rightarrow \infty} b_2 &\leq \limsup_{L \rightarrow \infty} 6LR_L^{1+1/\alpha} e^{-\ln(\tau L) - 2\delta R_L} \\ &= \limsup_{L \rightarrow \infty} \frac{6}{\tau} R_L^{1+1/\alpha} e^{-2\delta R_L} = 0, \end{aligned}$$

since R_L tends to infinity with L . \square

Proof of Theorem 1. Observe from Lemmas 3, 4, 5 and 6 that

$$\lim_{L \rightarrow \infty} \lim_{m \rightarrow \infty} d_{TV}(W_{m,L}, Po(\mathbb{E}[W_{m,L}])) = 0.$$

It follows from Lemmas 1 and 2, and the triangle inequality, that

$$\lim_{L \rightarrow \infty} \lim_{m \rightarrow \infty} d_{TV}(N_{iso}, Po(\mathbb{E}[W_{m,L}])) = 0.$$

It is straightforward to check that $\mathbb{E}[W_{m,L}] = mLp_i$ tends to $1/\tau$. Finally, a straightforward calculation shows that, if a sequence λ_n converges to λ , then $Po(\lambda_n)$ converges to $Po(\lambda)$ in total variation distance. Hence, invoking the triangle

inequality once more, we obtain that

$$d_{TV}(N_{\text{iso}}, Po(1/\tau)) \rightarrow 0 \text{ as } m, L \rightarrow \infty.$$

This completes the proof of the theorem. \square

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