t^{2/3}-scaling of current variance in interacting particle systems

Joint with Júlia Komjáthy and Timo Seppäläinen

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Stochastic Analysis Seminar, Oxford, June 9, 2008.

The models

Asymmetric simple exclusion process

Zero range

Hydrodynamics

Characteristics

Tool: the second class particle

Single

Many second class particles

Results

Normal fluctuations
Abnormal fluctuations

Proof

Upper bound

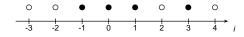
Lower bound

Coupling results

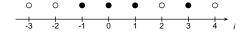
Other models

Linear models

Nonconvex, nonconcave

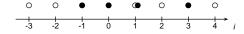


Bernoulli(ϱ) distribution; $\omega_i = 0$ or 1.



Particles try to jump

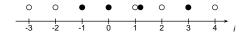
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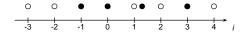
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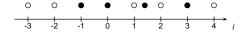
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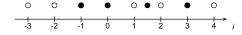
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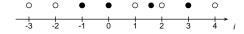
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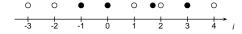
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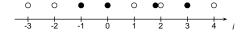
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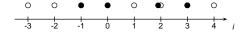
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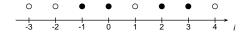
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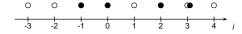
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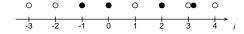
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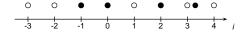
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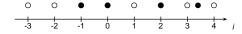
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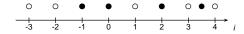
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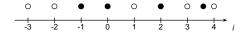
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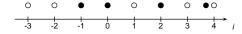
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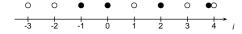
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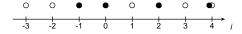
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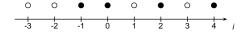
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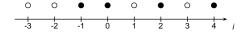
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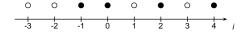
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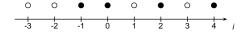
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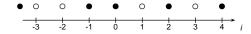
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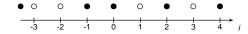
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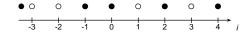
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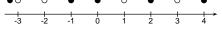
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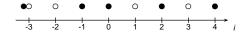
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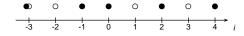
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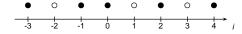
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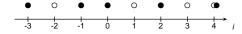
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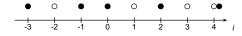
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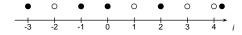
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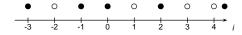
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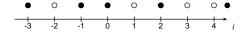
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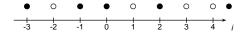
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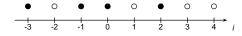
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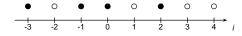
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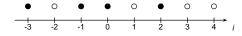
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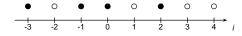
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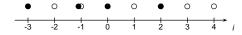
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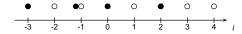
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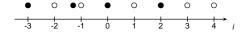
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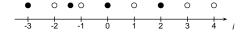
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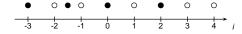
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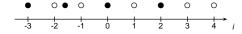
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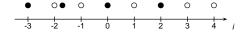
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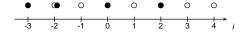
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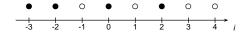
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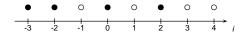
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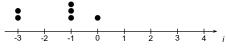
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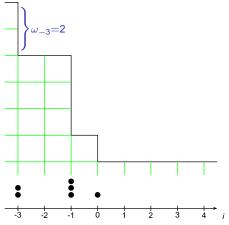
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The jump is suppressed if the destination site is occupied by another particle.

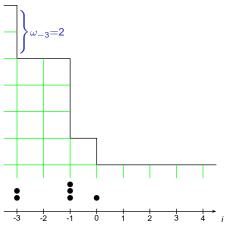
The Bernoulli(ρ) distribution is time-stationary for any $(0 < \rho < 1)$. Any translation-invariant stationary distribution is a mixture of Bernoullis.



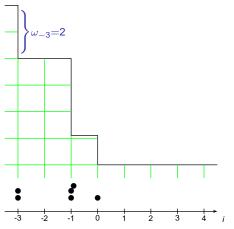
Poisson-type distribution; $\omega_i \in \mathbb{Z}^+$.



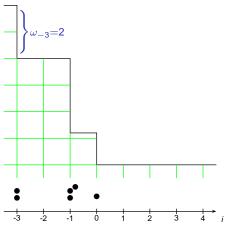
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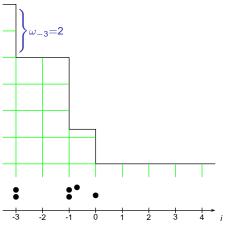
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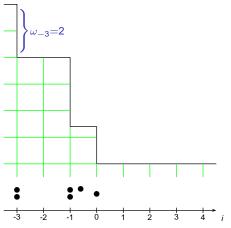
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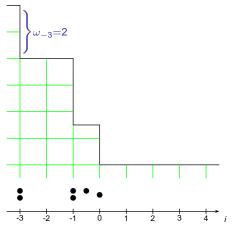
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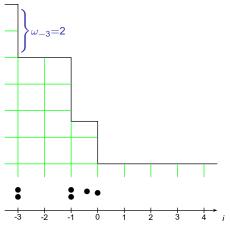
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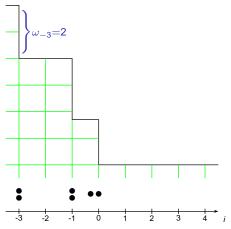
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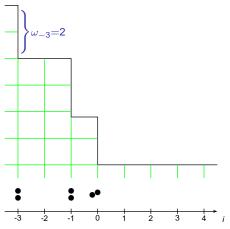
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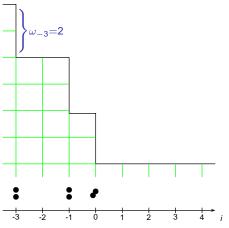
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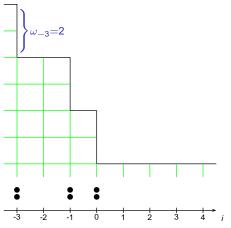
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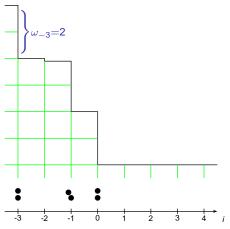
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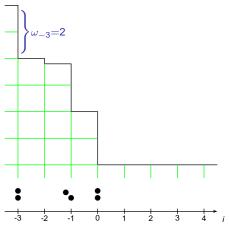
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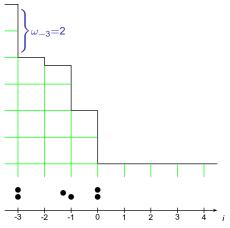
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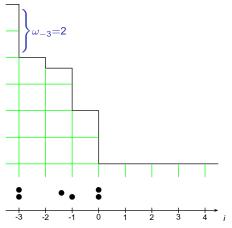
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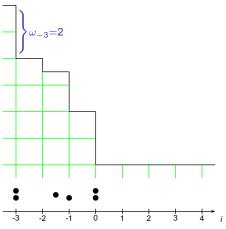
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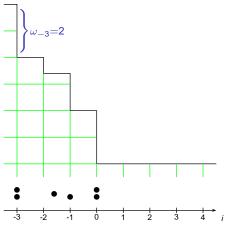
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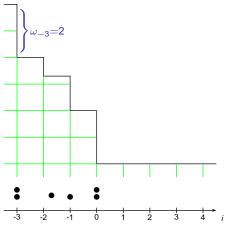
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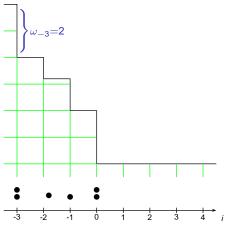
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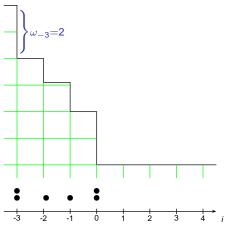
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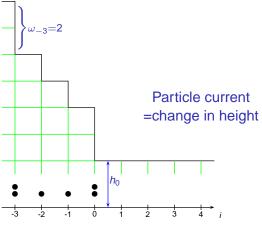
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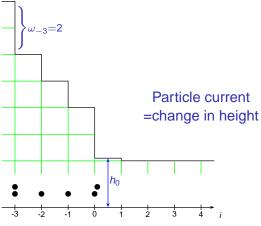
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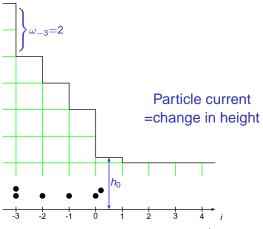
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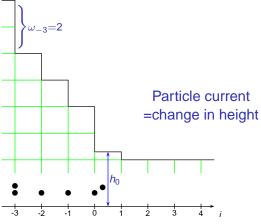
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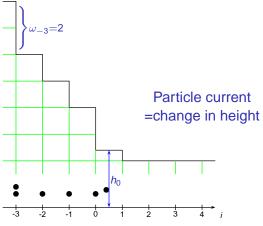
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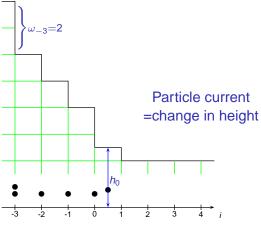
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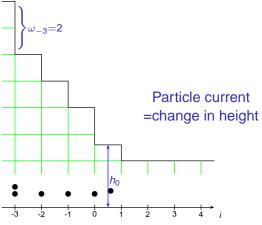
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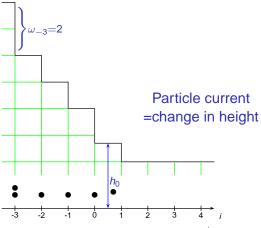
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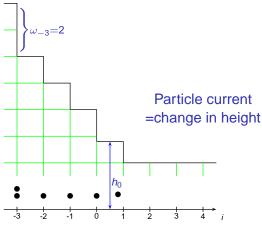
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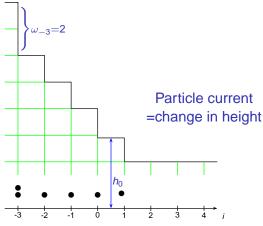
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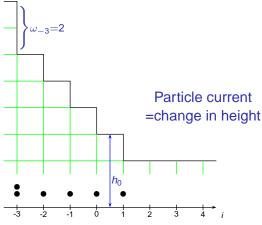
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The model

$$\begin{pmatrix} \omega_i \\ \omega_{i+1} \end{pmatrix} \rightarrow \begin{pmatrix} \omega_i - 1 \\ \omega_{i+1} + 1 \end{pmatrix} \qquad \text{with rate } p(\omega_i, \, \omega_{i+1}),$$

$$\begin{pmatrix} \omega_i \\ \omega_{i+1} \end{pmatrix} \rightarrow \begin{pmatrix} \omega_i + 1 \\ \omega_{i+1} - 1 \end{pmatrix} \qquad \text{with rate } q(\omega_i, \, \omega_{i+1}), \text{ where }$$

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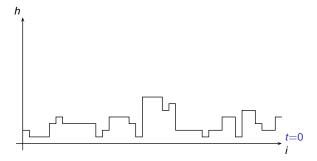
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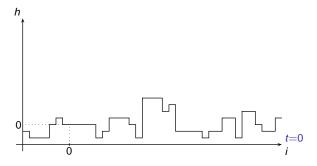
- p and q are such that they keep the state space (ASEP, ZRP),
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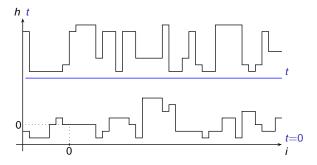
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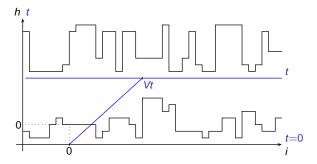
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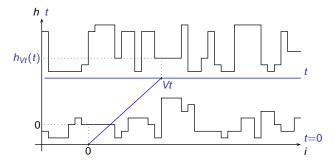
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- they satisfy some regularity conditions to make sure the dynamics exists.











 $h_{Vt}(t)$ = height as seen by a moving observer of velocity V. = net number of particles passing the window $s \mapsto Vs$.

(Remember: particle current=change in height.)

The question

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... is the properties of $h_{Vt}(t)$ under the time-stationary evolution.

▶ $\mathbf{E}(h_{Vt}(t)) = t \cdot \mathbf{E}(\text{growth rate})$ is easily computed with martingales.

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Hydrodynamics (very briefly)

The density $\varrho:=\mathbf{E}(\omega)$ and the hydrodynamic flux $H:=\mathbf{E}[\text{growth rate}]$ both depend on a parameter of the stationary distribution.

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- If the process is locally in equilibrium, but changes over some *large scale* (variables $X = \varepsilon i$ and $T = \varepsilon t$), then

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▶ The characteristics is a path X(T) where $\rho(T, X(T))$ is constant.

Characteristics (very briefly)

$$\partial_T \varrho + \partial_X H(\varrho) = 0$$

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 $\partial_T \varrho + \mathbf{H}'(\varrho) \cdot \partial_X \varrho = 0$ (while smooth)

$$\begin{split} \partial_{\mathcal{T}} \varrho + \partial_{X} \boldsymbol{H}(\varrho) &= 0 \\ \partial_{\mathcal{T}} \varrho + \boldsymbol{H}'(\varrho) \cdot \partial_{X} \varrho &= 0 \qquad \text{(while smooth)} \\ &\qquad \qquad \frac{\mathrm{d}}{\mathrm{d}\mathcal{T}} \varrho(\mathcal{T}, \, X(\mathcal{T})) = 0 \end{split}$$

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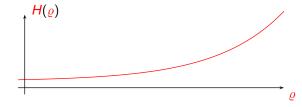
$$R = \frac{H(\varrho) - H(\lambda)}{\rho - \lambda}.$$

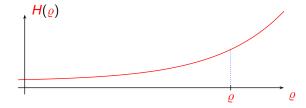
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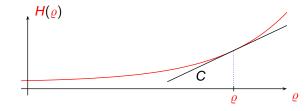
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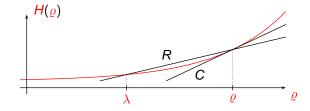
This would be the speed of a shock of densities ρ and λ .



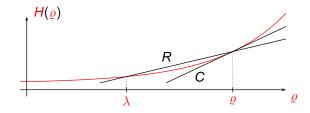




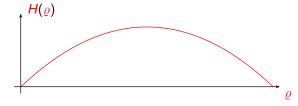
$$C = H'(\varrho)$$

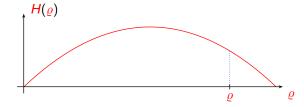


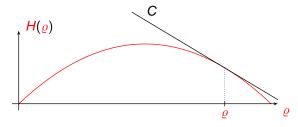
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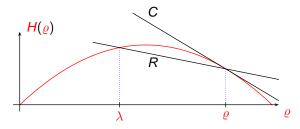
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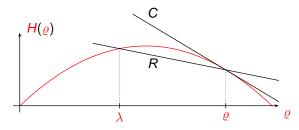




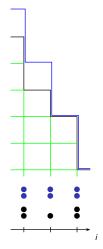
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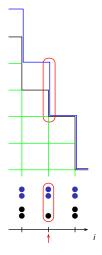


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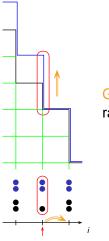


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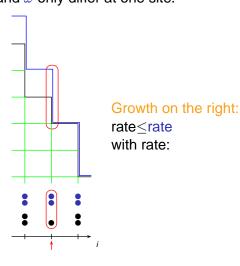


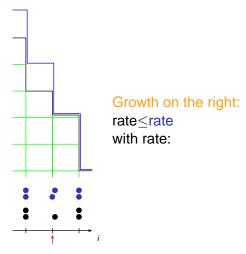


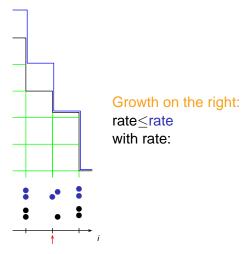
States ω and ω only differ at one site.

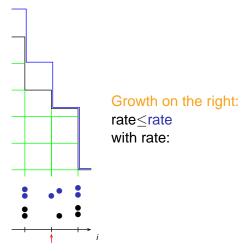


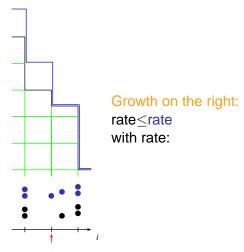
Growth on the right: rate<rate

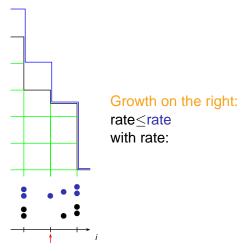


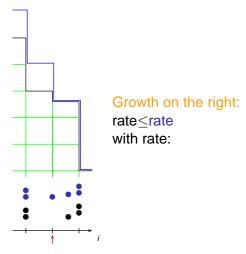


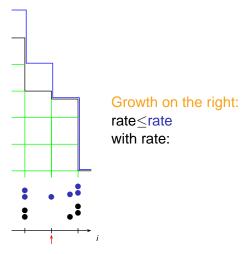


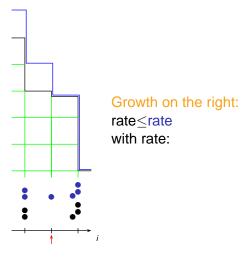


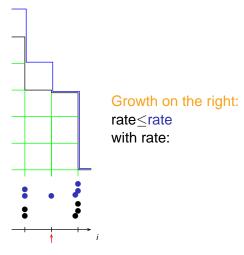




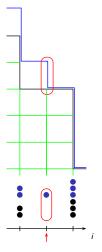








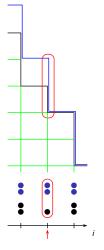
States ω and ω only differ at one site.



Growth on the right:

rate < rate with rate:

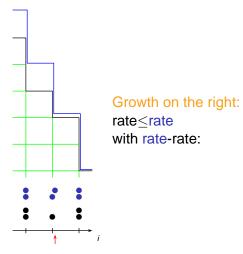
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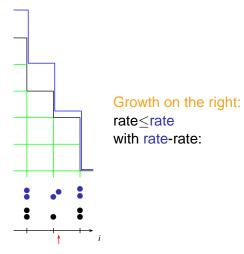


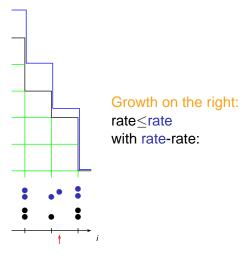
Growth on the right:

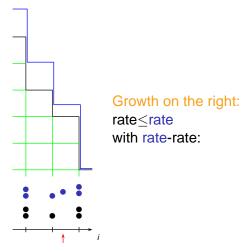
rate < rate

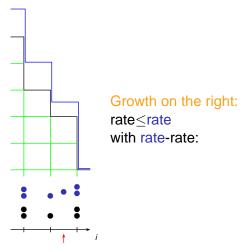
with rate-rate:

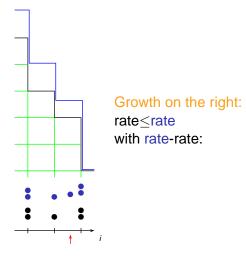


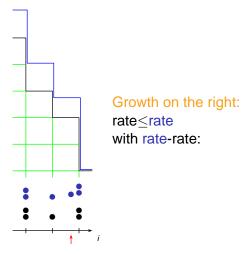


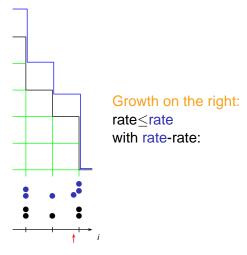


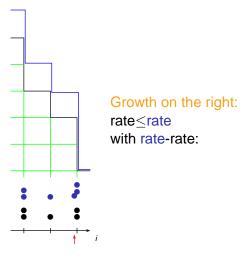




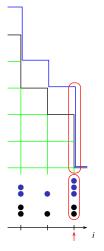








States ω and ω only differ at one site.



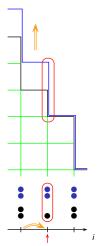
Growth on the right:

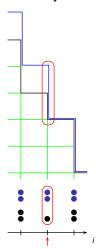
rate < rate

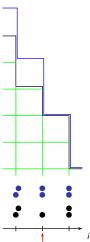
with rate-rate:

States ω and ω only differ at one site.

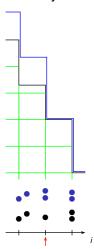
Growth on the left: rate>rate



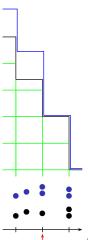


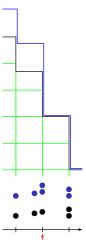


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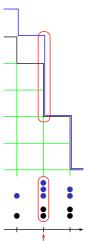
States ω and ω only differ at one site.



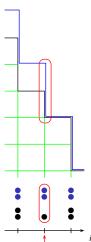


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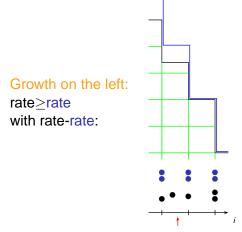
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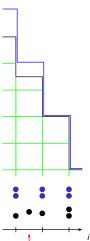
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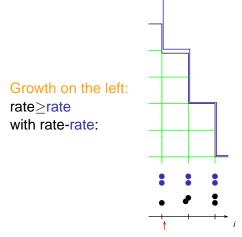
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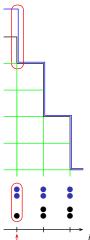
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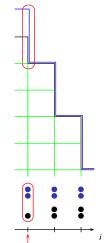
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Growth on the left: rate>rate with rate-rate:

A single discrepancy, the second class particle, is conserved. Its position at time t is Q(t).

Theorem (B. - Seppäläinen; also ideas from B. Tóth, H. Spohn and M. Prähofer)

Started from (almost) equilibrium,

$$\mathsf{E}(\mathsf{Q}(t)) = C \cdot t$$

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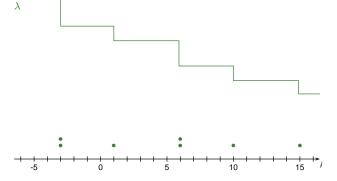
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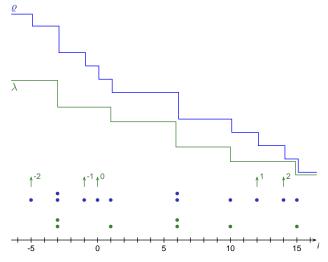
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The second class particle follows the characteristics, people have known this for a long time.

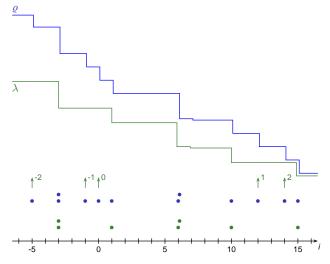
Many second class particles

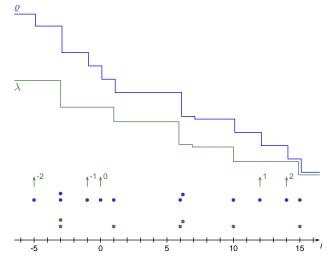


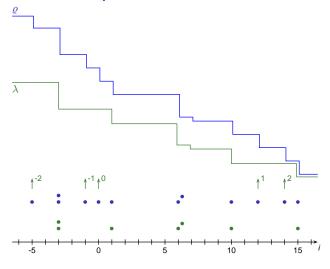
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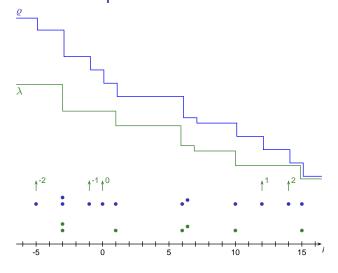


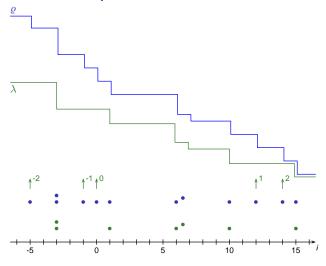
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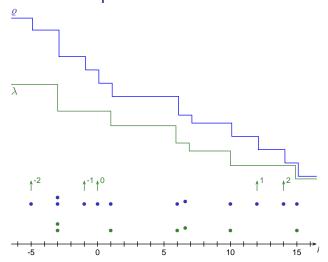


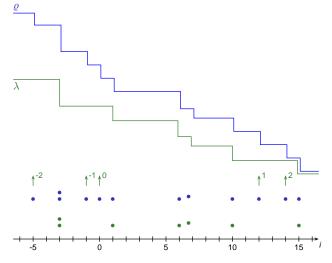


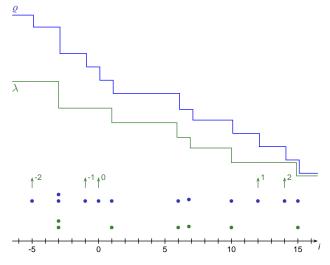


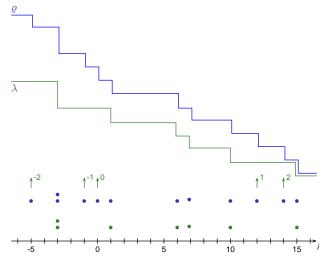


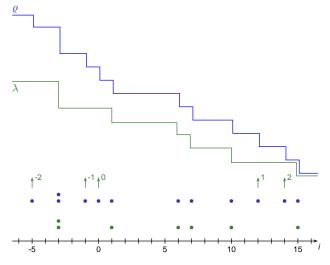


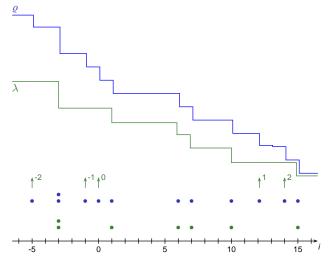




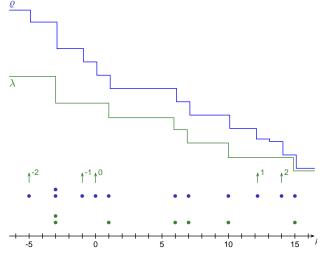


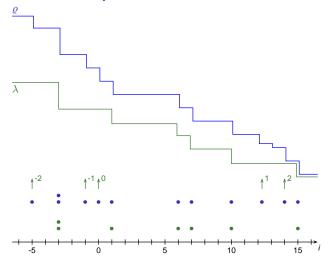


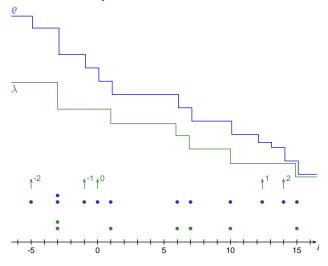


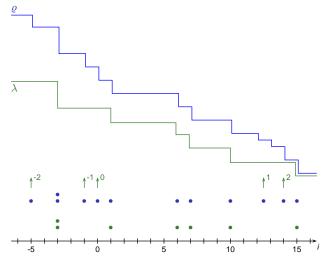


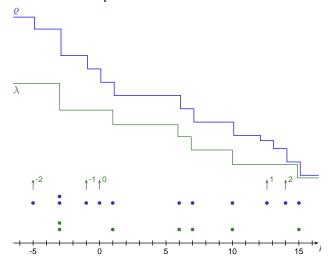


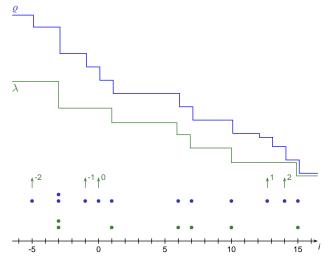


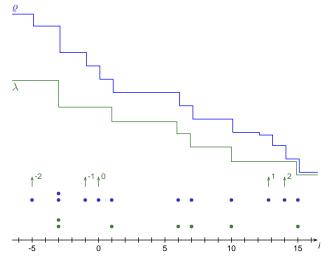


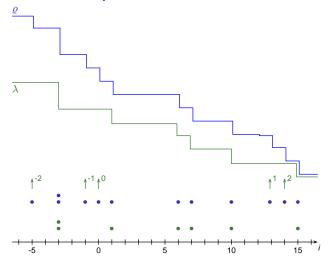


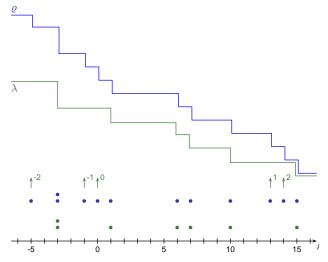


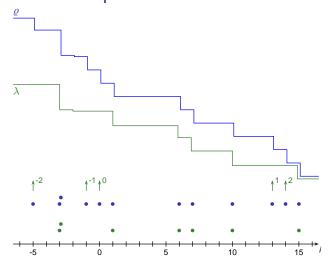


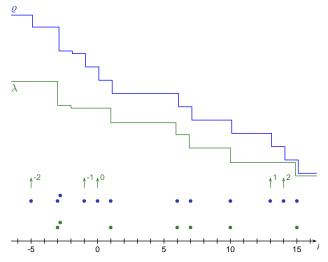


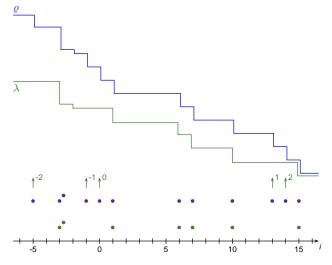


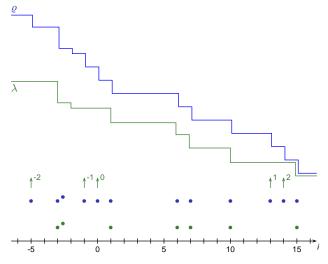


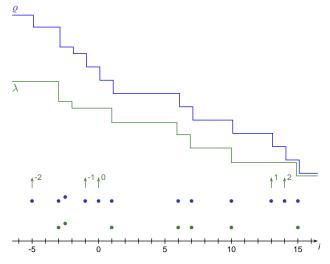


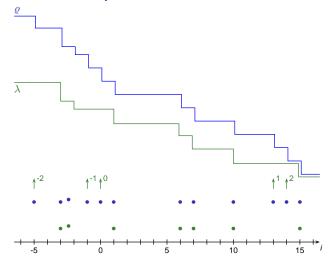


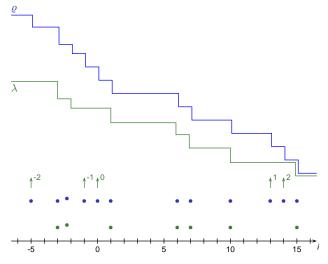


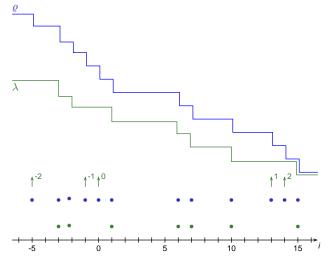


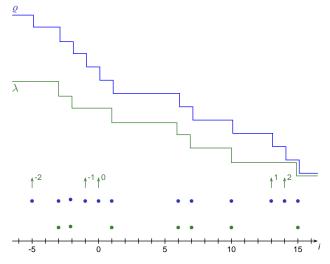


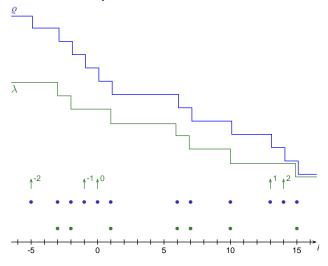


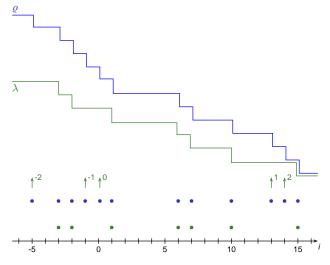


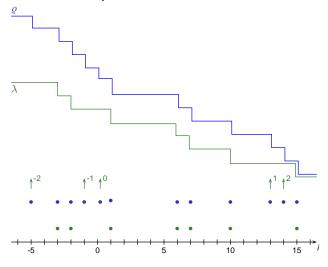


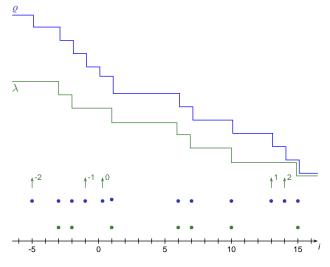


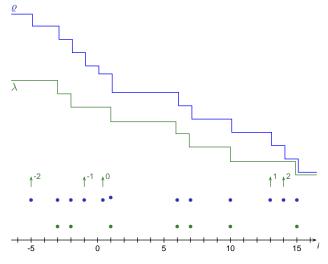


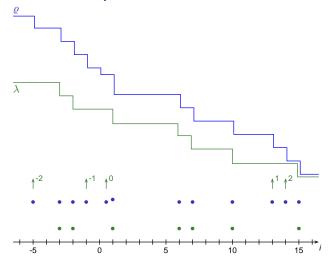


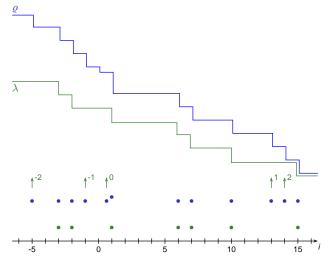


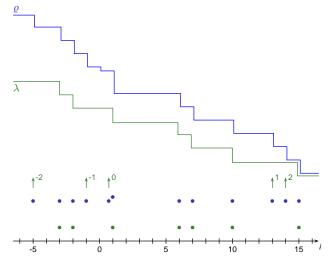


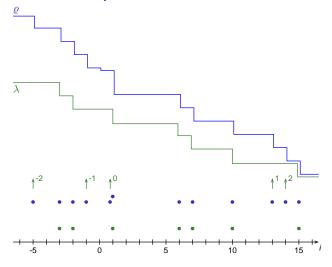


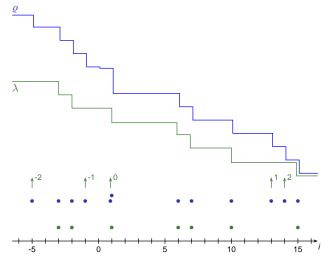


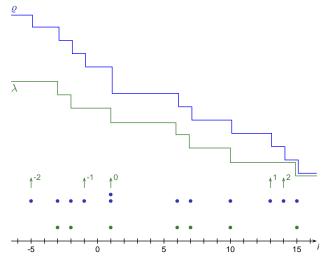


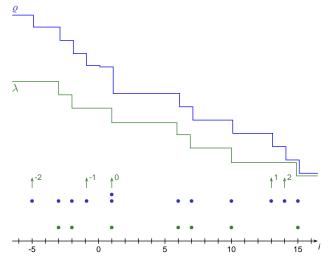


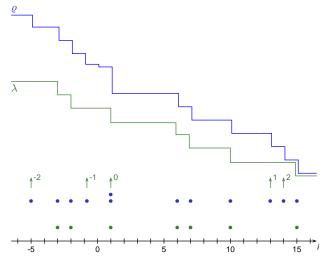


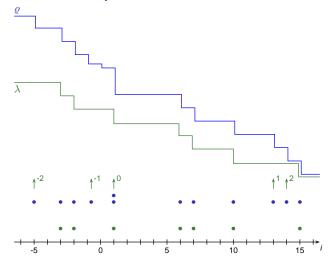


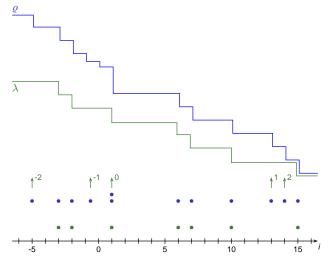


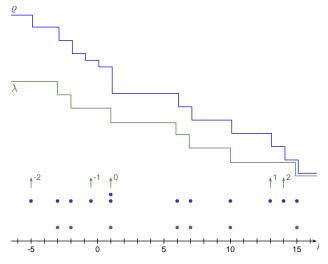


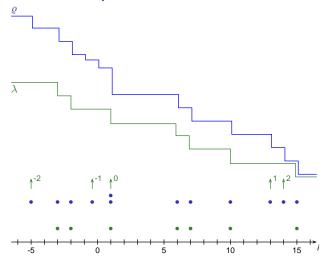


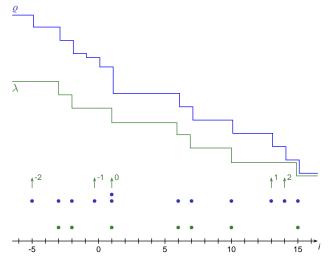


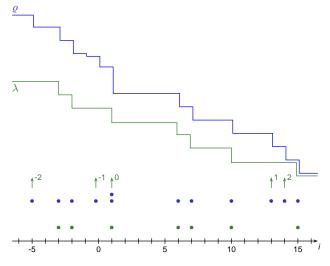


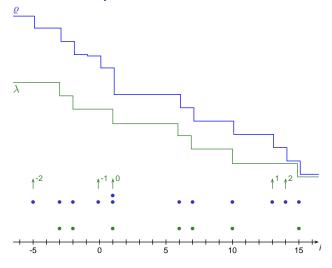


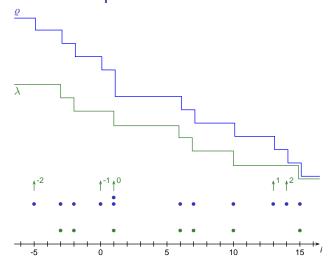


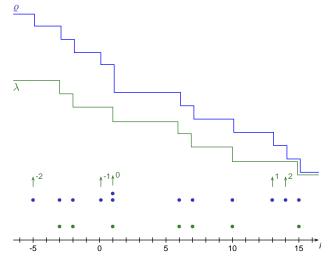


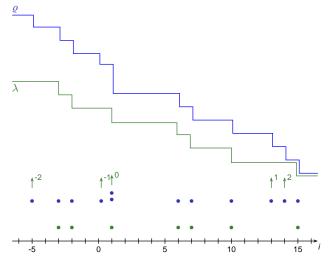


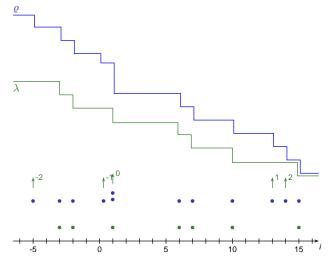


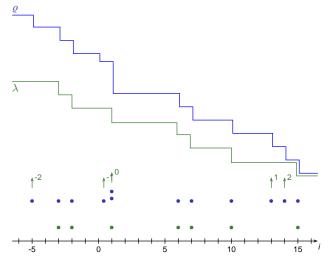


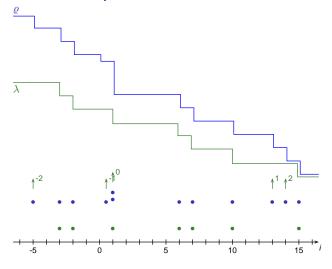


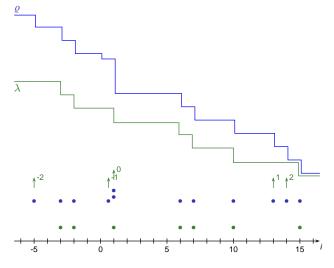


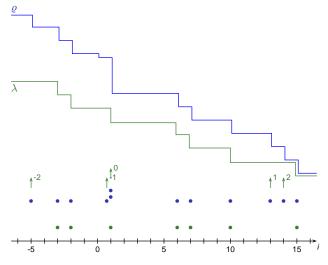


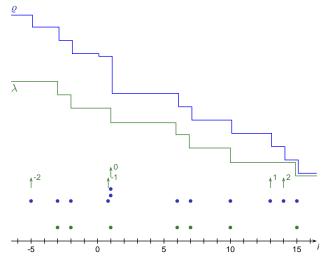


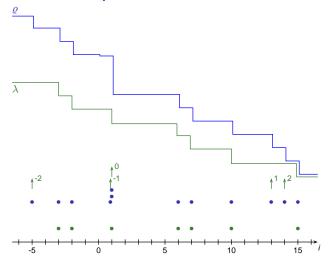


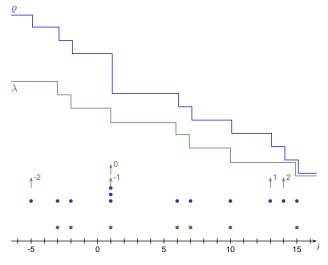


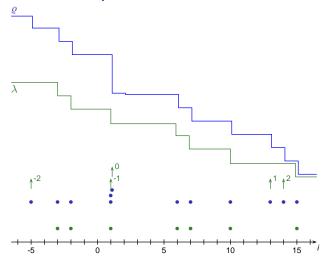


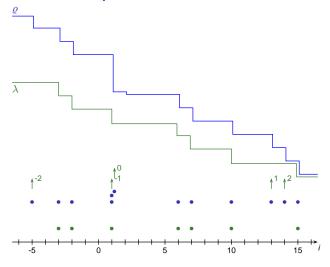


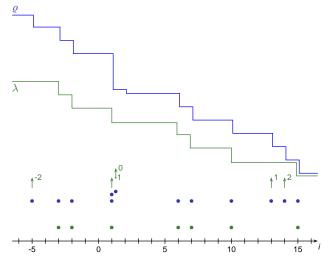


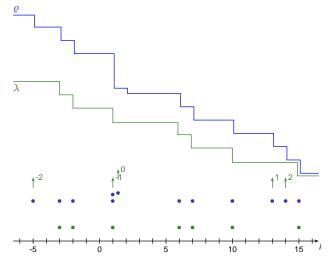


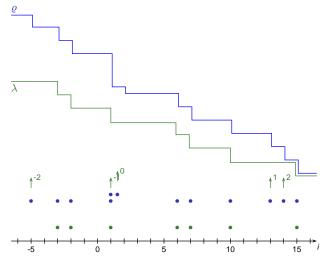


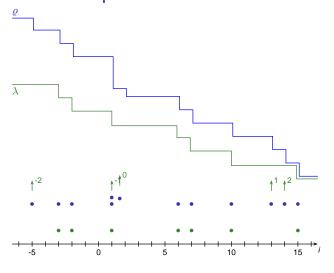


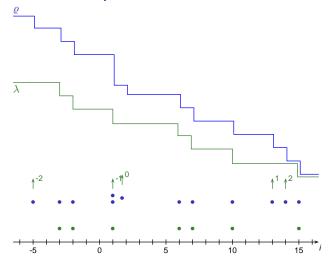


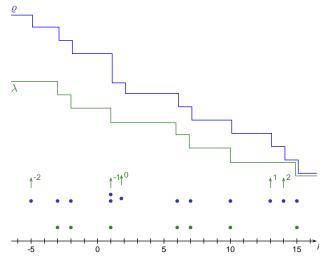




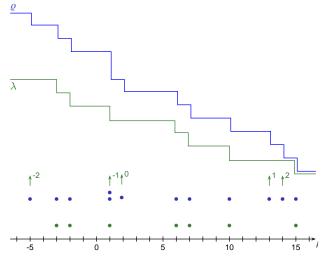


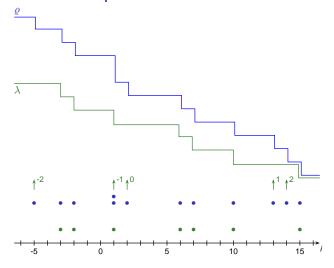


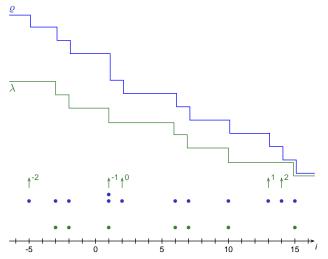










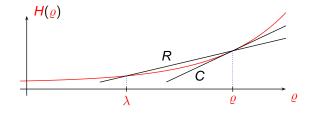


Picture:

The position X(t) of \uparrow^0 follows the Rankine-Hugoniot speed R.

Characteristics (very briefly)

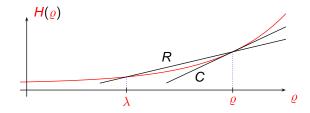
Convex flux (some cases of AZRP):



Recall
$$C = H'(\varrho) > R = \frac{H(\varrho) - H(\lambda)}{\varrho - \lambda}$$

Characteristics (very briefly)

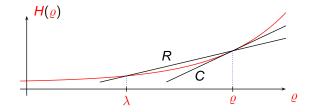
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Do we have $Q(t) \stackrel{?}{\geq} X(t)$

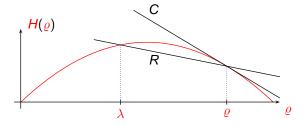
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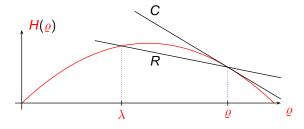
Do we have $Q(t) \stackrel{?}{\geq} X(t) + \text{tight error}$

Concave flux (ASEP, AZRP):



$$C = H'(\varrho) < R = \frac{H(\varrho) - H(\lambda)}{\varrho - \lambda}$$

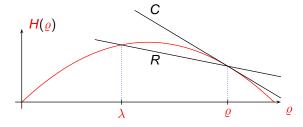
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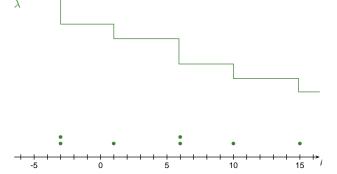
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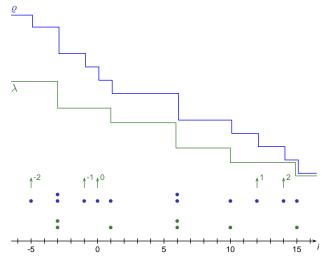


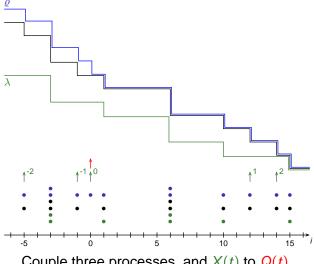
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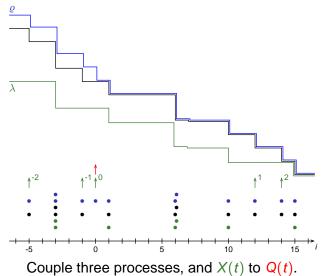
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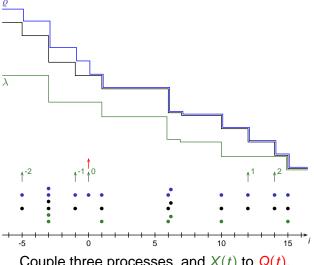




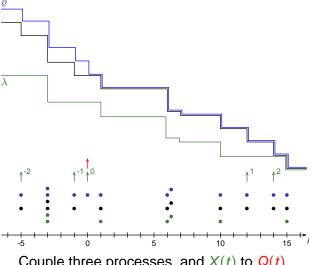


Couple three processes, and X(t) to Q(t).

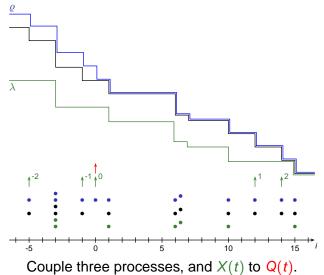


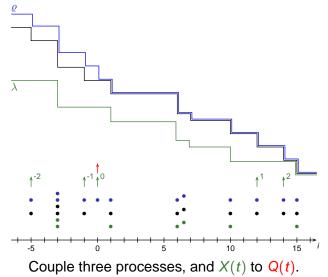


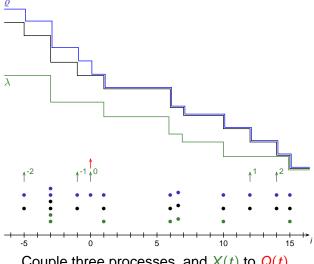
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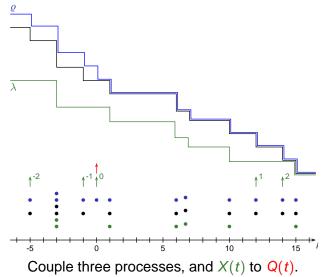
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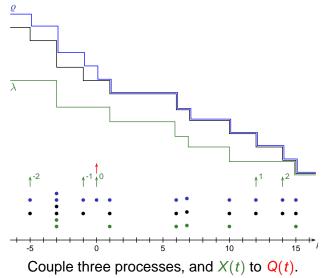


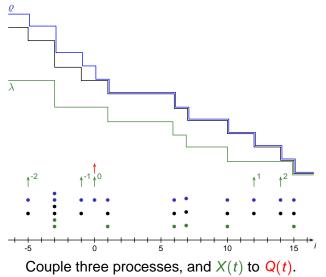


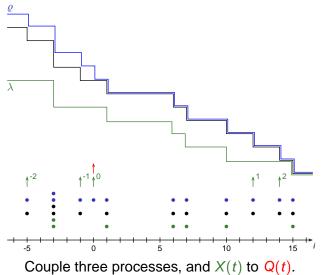


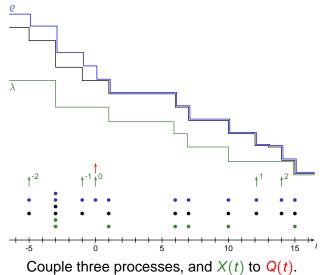
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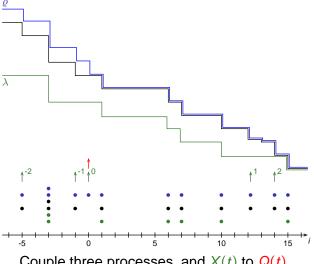




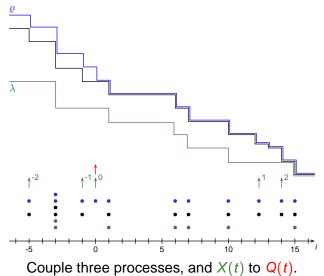


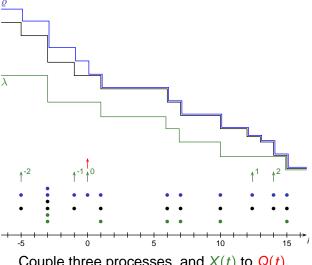




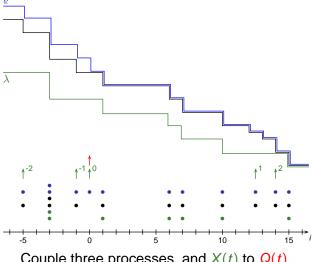


Couple three processes, and X(t) to Q(t).

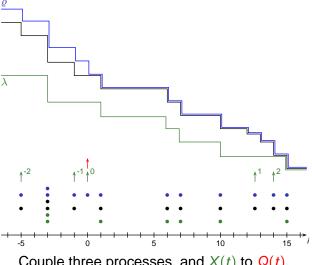




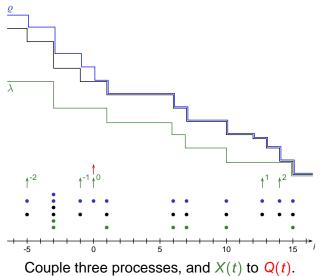
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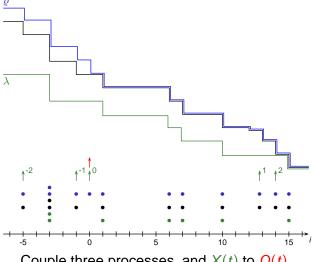


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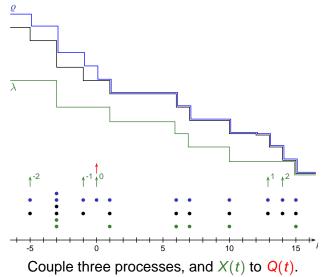


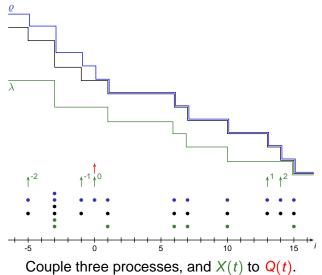
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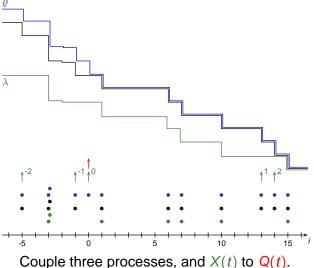


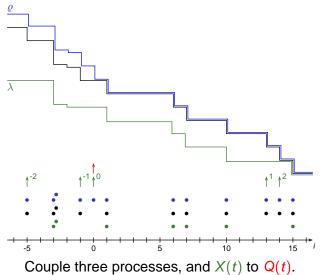


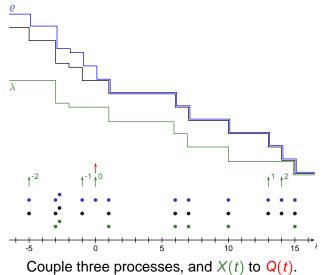
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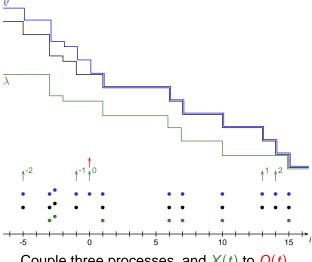




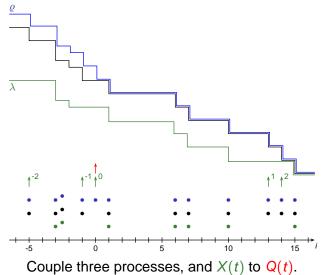


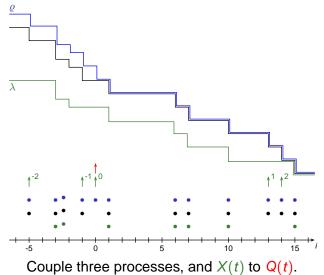


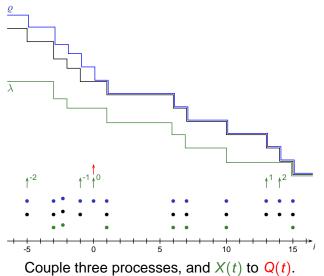


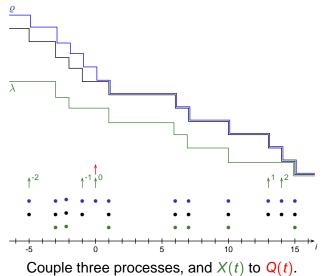


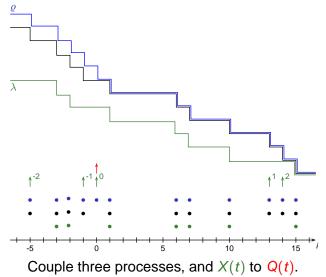
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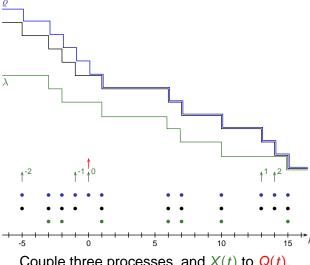




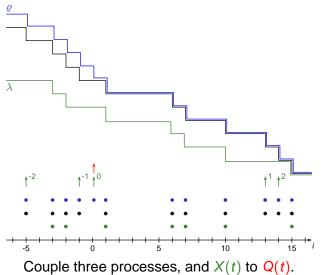


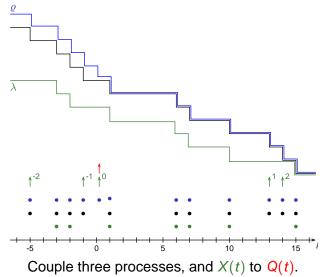


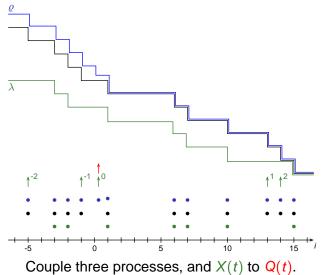


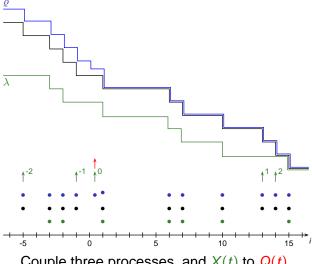


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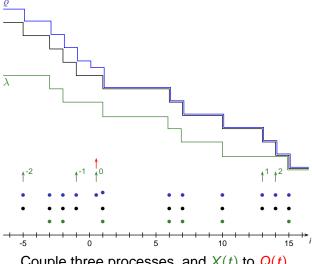




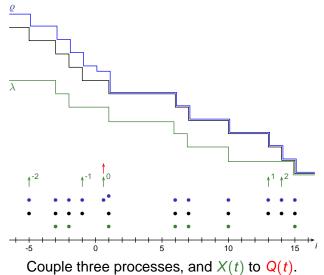


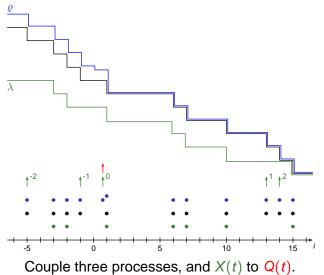


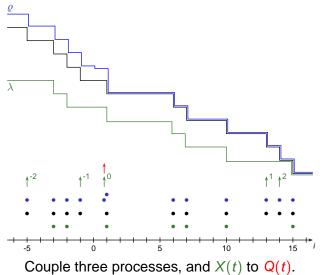
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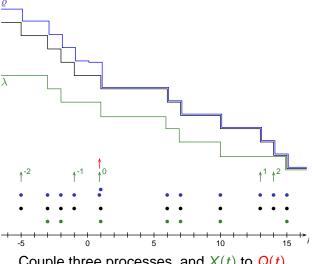


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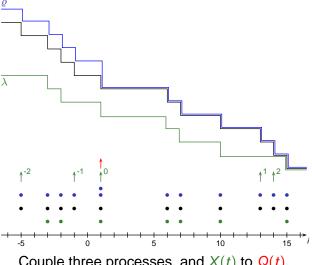




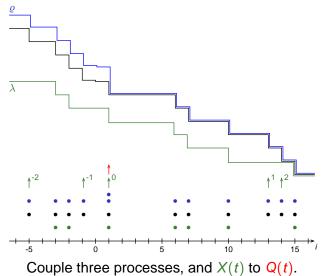


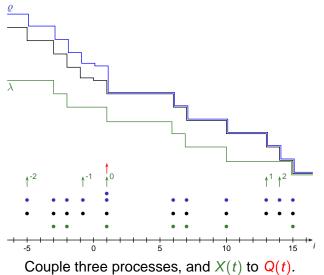


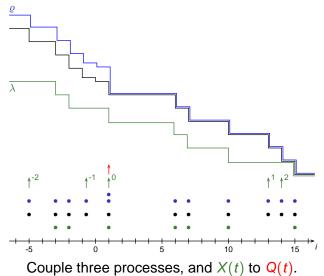
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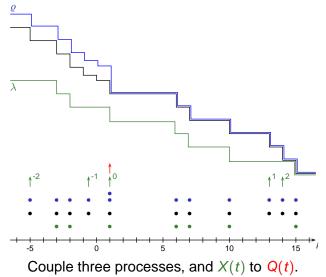


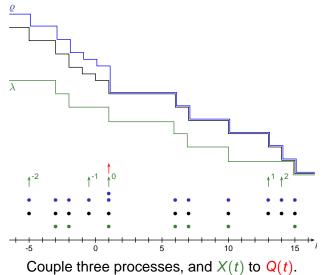
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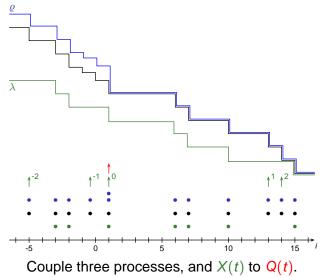


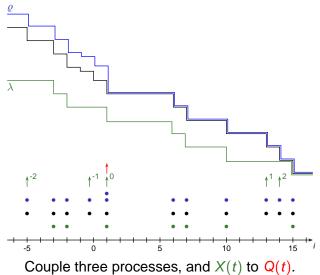


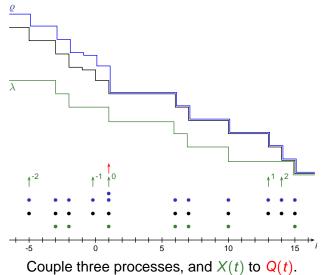


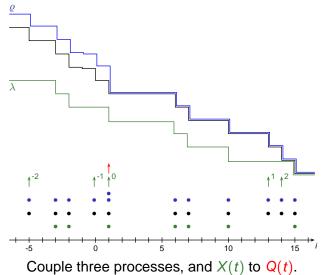


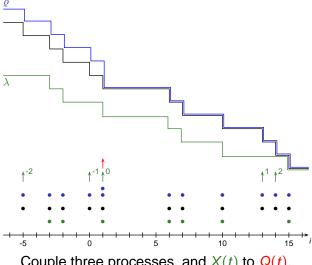




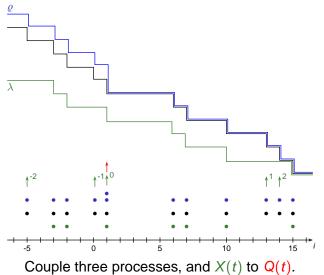


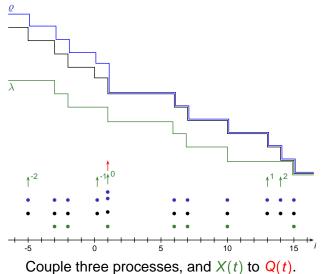


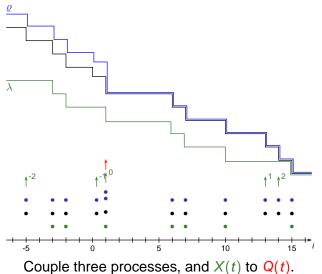


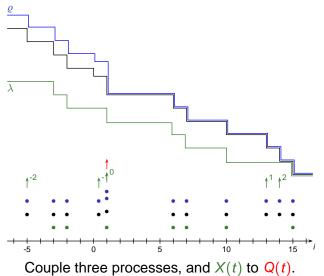


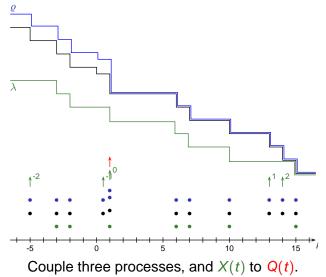
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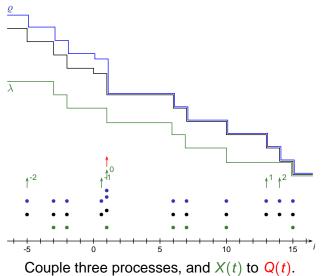


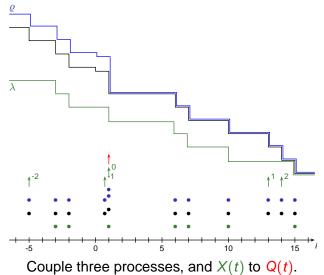


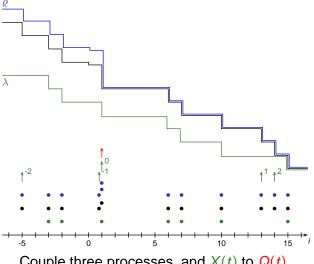




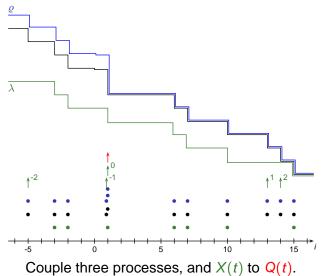


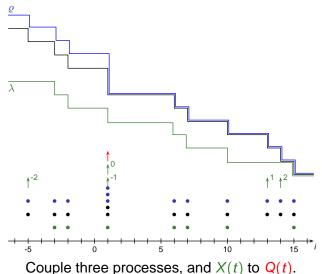


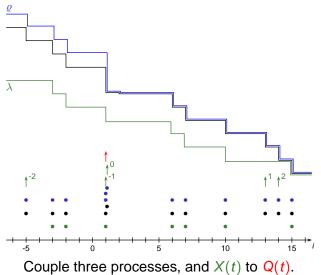


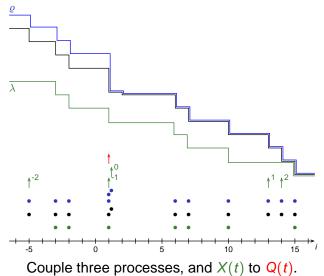


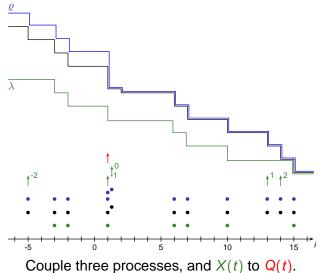
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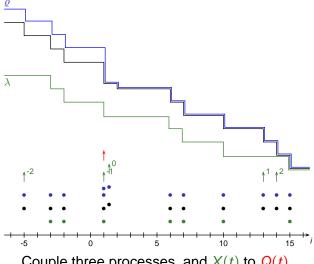




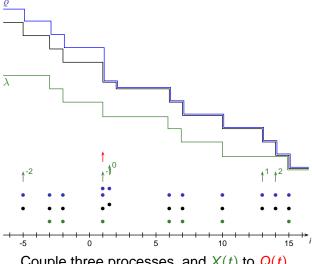




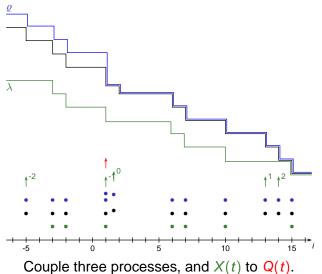


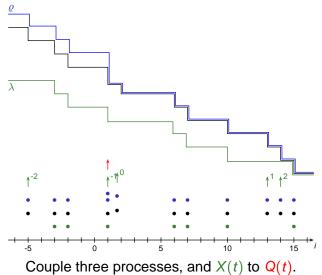


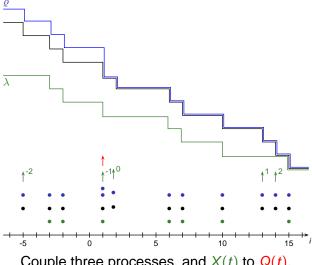
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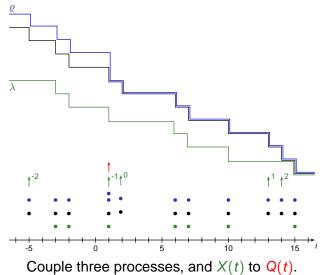
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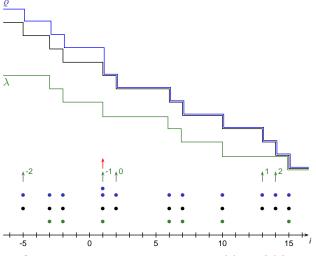






Couple three processes, and X(t) to Q(t).





Couple three processes, and X(t) to Q(t). We'll assume $Q(t) \le X(t)$ can be achieved.

Normal fluctuations:

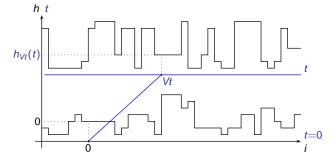
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$$\lim_{t \to \infty} \frac{\mathsf{Var}(h_{\mathit{Vt}}(t))}{t} = \mathsf{Var}(\omega) \cdot |C - V|$$

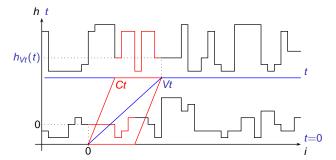


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Initial fluctuations are transported along the characteristics on this scale.

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Theorem (B. - Komjáthy - Seppäläinen (ASEP, some special TAZRP's so far, but...))

$$0<\liminf_{t\to\infty}\frac{\mathsf{Var}(h_{Ct}(t))}{t^{2/3}}\leq\limsup_{t\to\infty}\frac{\mathsf{Var}(h_{Ct}(t))}{t^{2/3}}<\infty.$$

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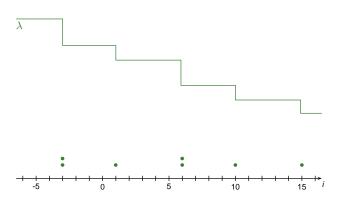
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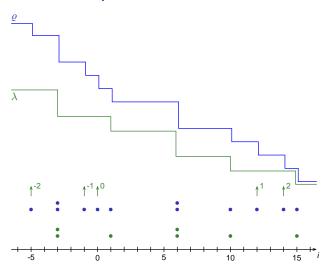
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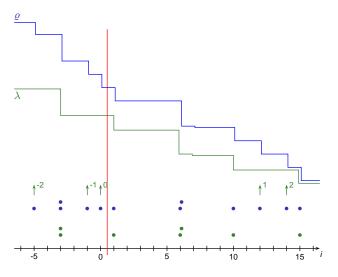
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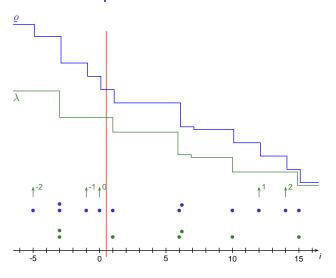
There are limit distribution results for TASEP by Johansson 2000, Prähofer and Spohn 2001, Ferrari and Spohn 2006.

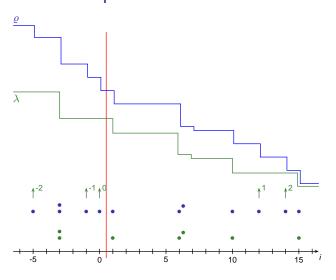
Their methods give limit distributions as well, but are very model-dependent: they rewrite the model as a determinantal process, and perform asymptotic analysis of the determinants.

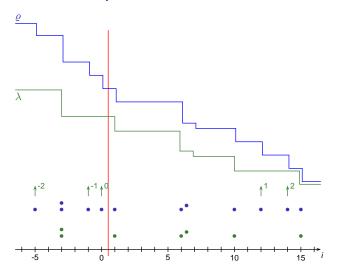


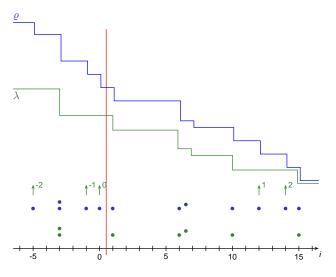


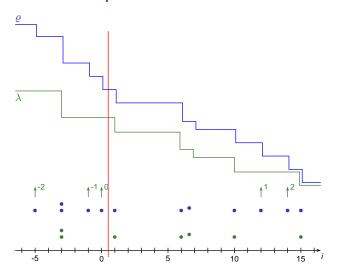


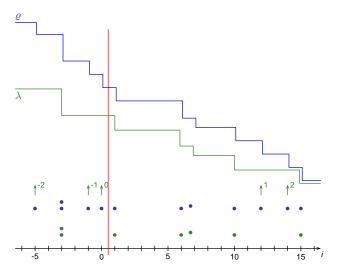


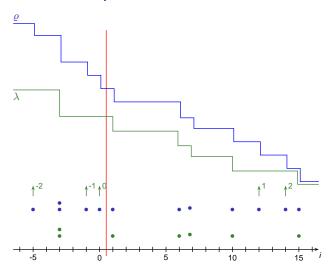


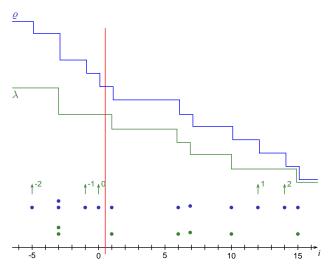


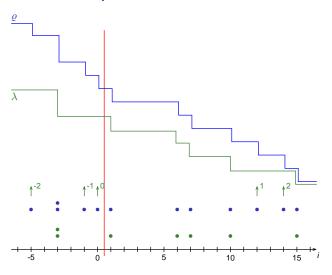


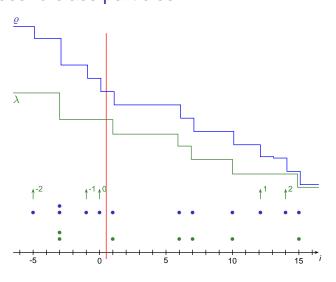


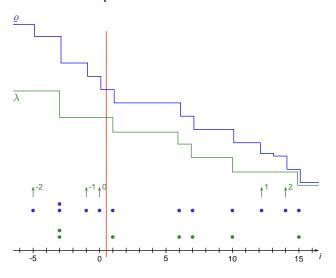


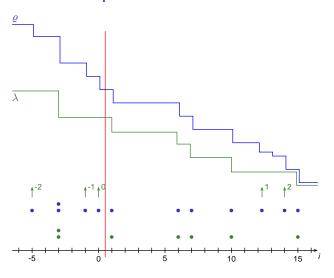


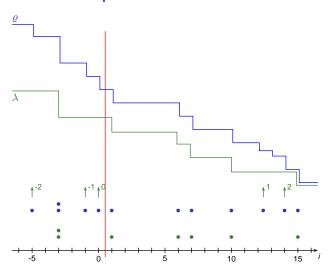


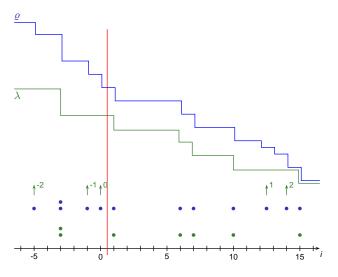


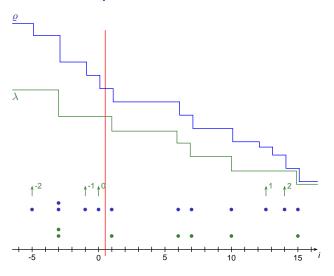


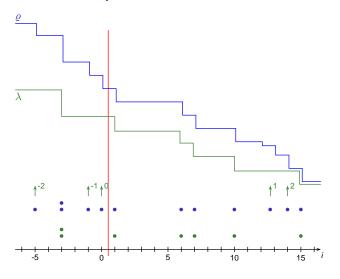


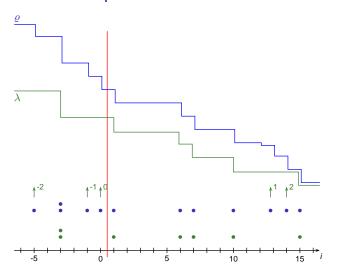


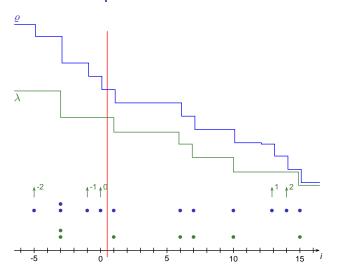


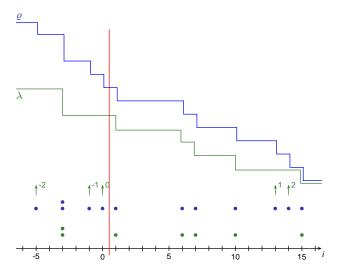


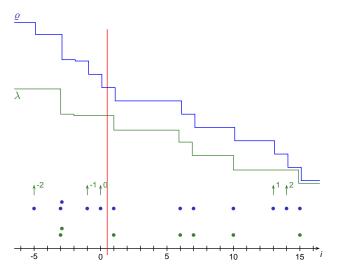


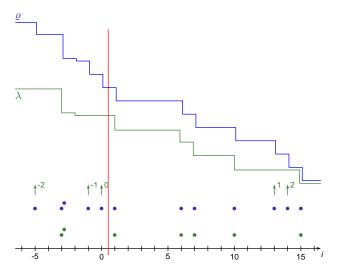


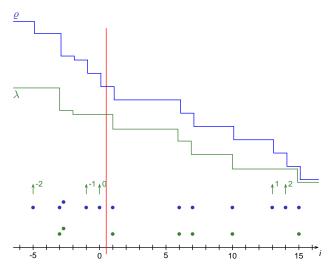


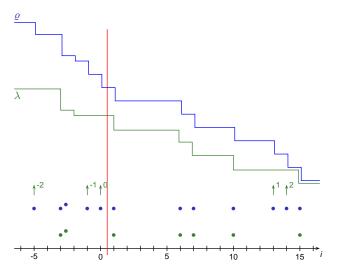


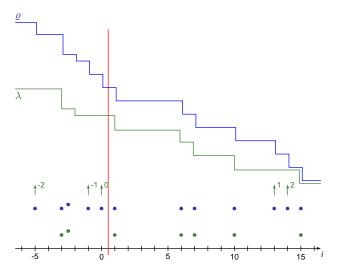


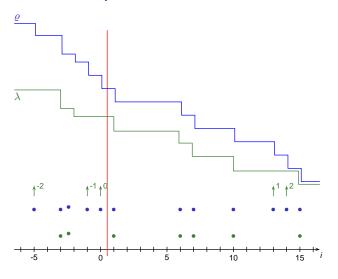


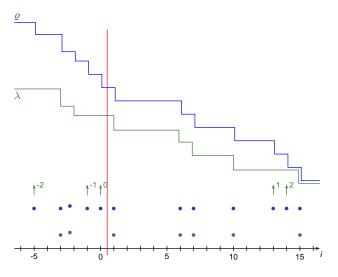


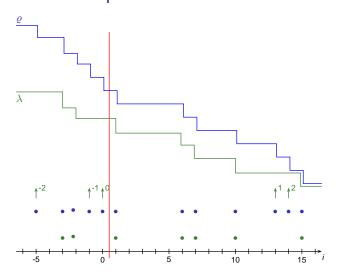


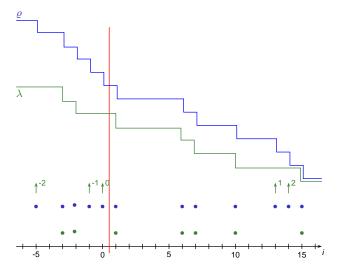


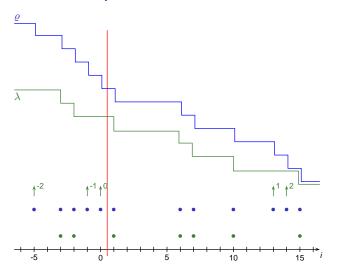


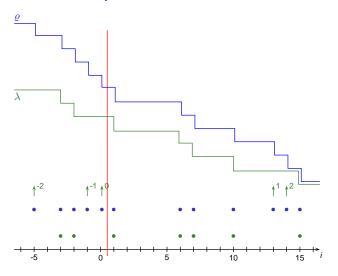


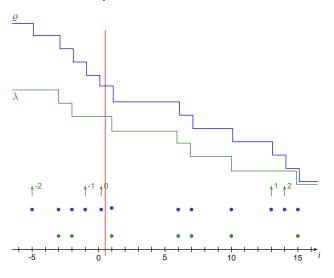


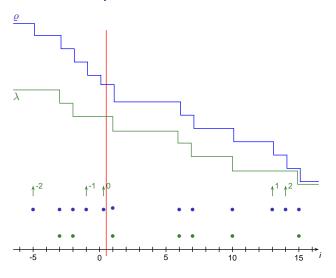


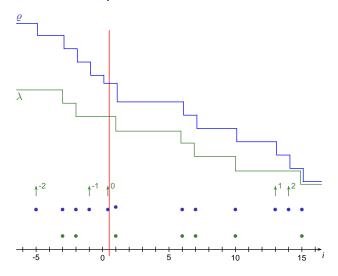


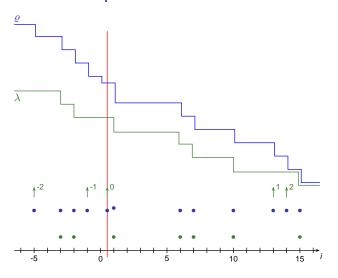


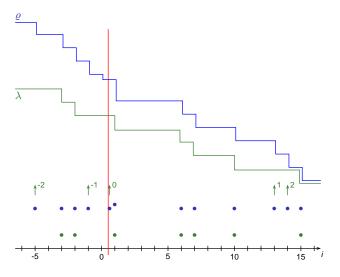


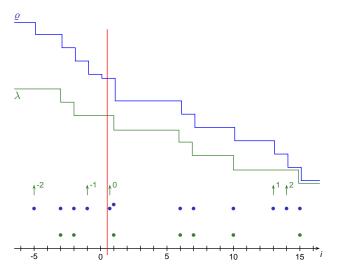


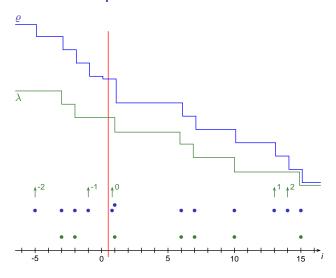


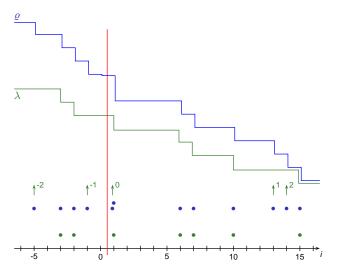


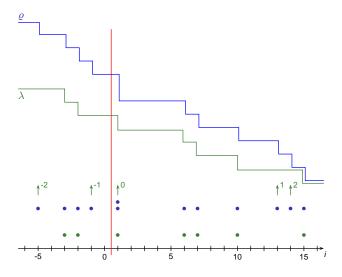


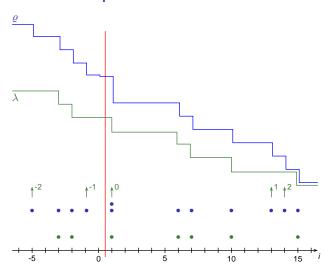


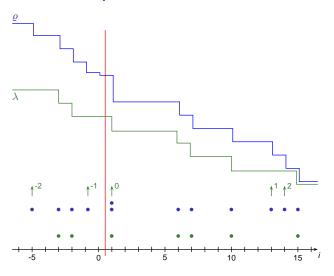


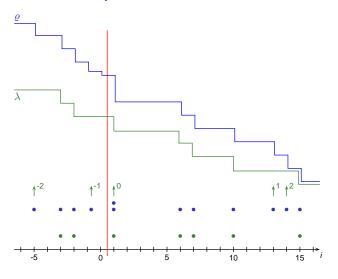


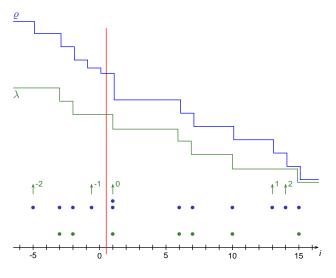


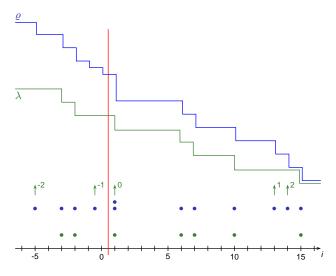


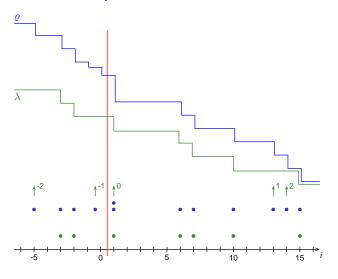


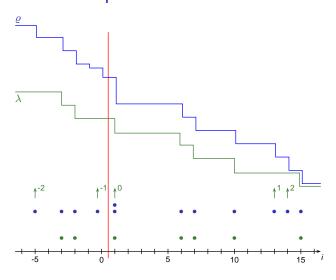


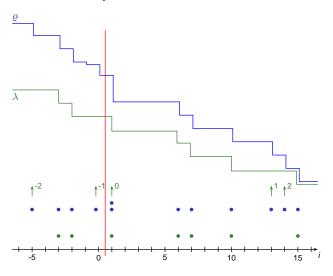


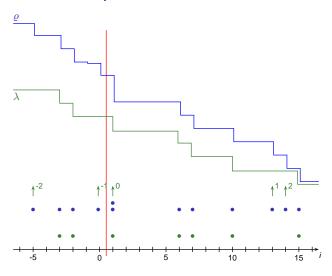


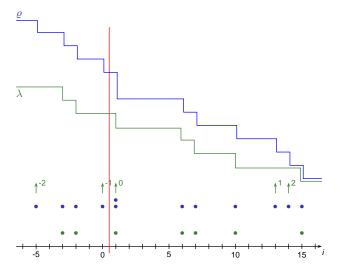


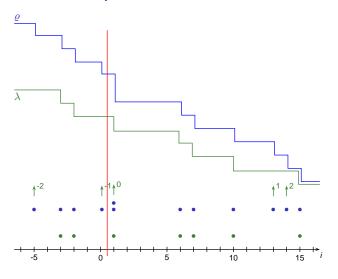


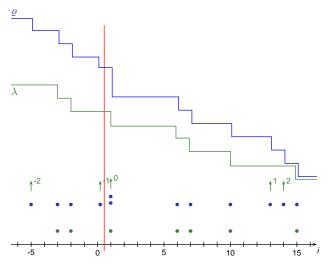


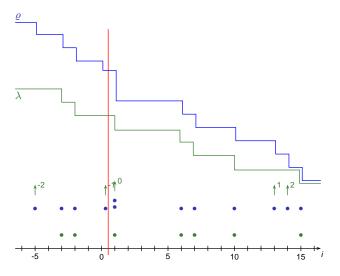


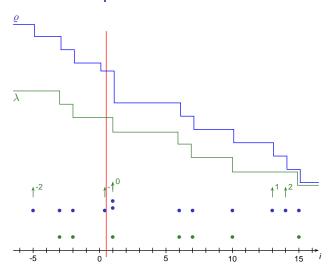


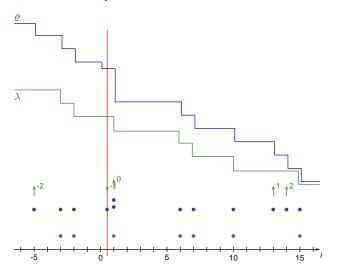


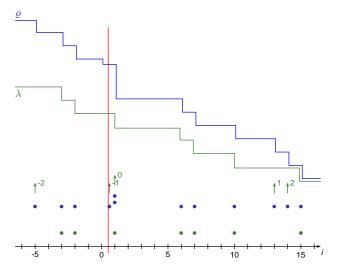


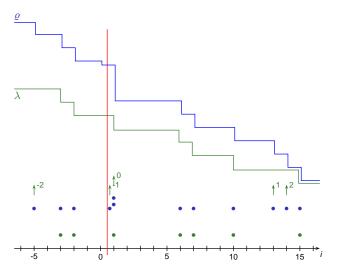


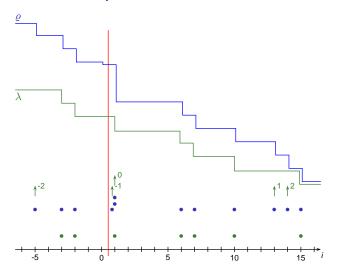


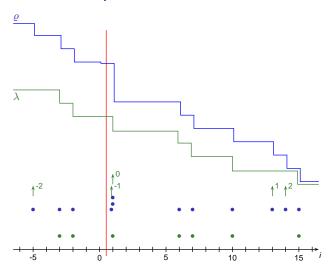


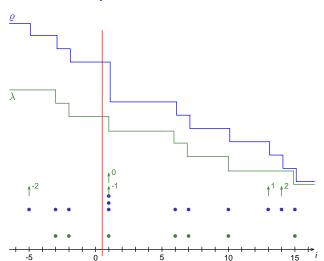


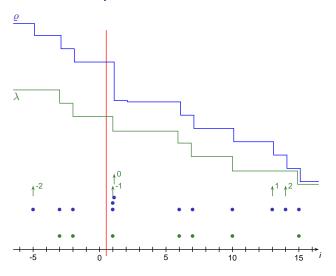


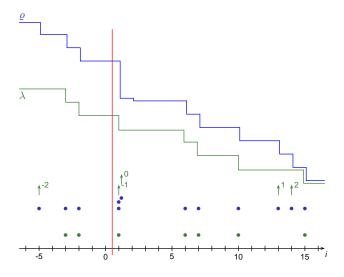


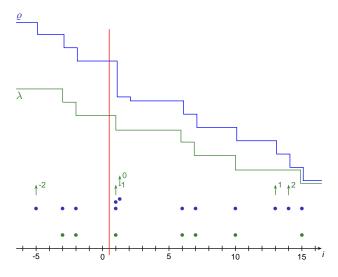


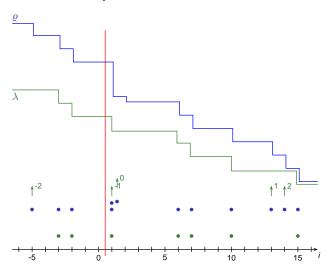


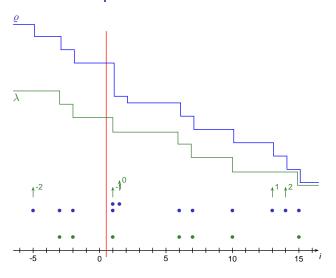


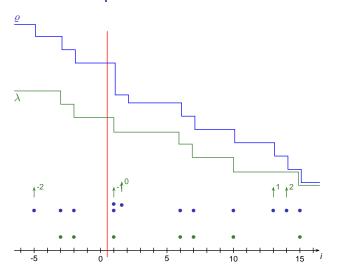


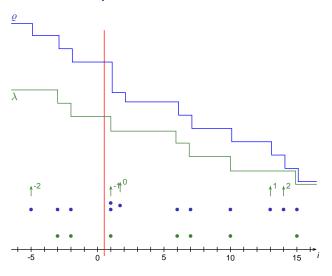


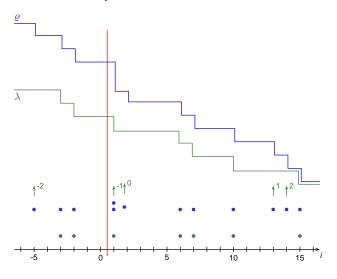


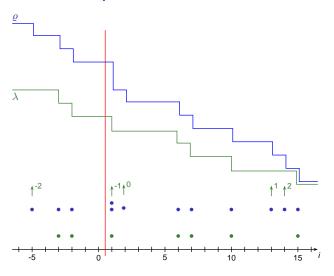


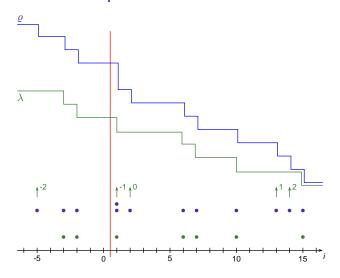












$P{Q(t) \text{ is too large}}$

 $P{Q(t) \text{ is too large}} \le P{X(t) \text{ is too large}}$

 $P{Q(t) \text{ is too large}} \le P{X(t) \text{ is too large}}$ $\leq \mathbf{P}\{\text{too many} \uparrow \text{'s have crossed } Ct\}$

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≤ P{too many \( \) 's have crossed \( Ct \) \)

                               < \mathbf{P}\{h_{Ct}(t) - h_{Ct}(t) \text{ is too large}\}.
```

Centering $h_{Ct}(t) - h_{Ct}(t)$ brings in a second-order Taylor-expansion of $H(\rho)$.

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Optimize "too large(λ)" in λ , use Chebyshev's inequality and relate $Var(h_{Ct}(t))$ to $Var(h_{Ct}(t))$.

$$\begin{split} \mathbf{P}\{\mathbf{Q}(t) \text{ is too large}\} &\leq \mathbf{P}\{X(t) \text{ is too large}\}\\ &\leq \mathbf{P}\{\text{too many} \uparrow \text{'s have crossed } Ct\}\\ &\leq \mathbf{P}\{h_{Ct}(t) - h_{Ct}(t) \text{ is too large}(\lambda)\}. \end{split}$$

Centering $h_{Ct}(t) - h_{Ct}(t)$ brings in a second-order Taylor-expansion of $H(\varrho)$. This is another point where concavity of the flux matters.

Optimize "too large(λ)" in λ , use Chebyshev's inequality and relate $Var(h_{Ct}(t))$ to $Var(h_{Ct}(t))$.

The computations result in (remember E(Q(t)) = Ct)

$$\mathbf{P}\{\mathbf{Q}(t)-Ct\geq u\}\leq c\cdot\frac{t^2}{u^4}\cdot\mathbf{Var}(h_{Ct}(t)).$$

Theorem (B. - Seppäläinen; also ideas from B. Tóth, H. Spohn and M. Prähofer)

Started from (almost) equilibrium,

$$\mathsf{Var}(h_{Ct}(t)) = c \cdot \mathsf{E}|\,\mathsf{Q}(t) - C \cdot t|$$

in the whole family of processes.

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Hence proceed with

$$\begin{aligned} \mathbf{P}\{\frac{\mathbf{Q}(t) - Ct \ge u\} \le c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{Ct}(t)) \\ &= c \cdot \frac{t^2}{u^4} \cdot \mathbf{E}|\mathbf{Q}(t) - C \cdot t|. \end{aligned}$$

With

$$\widetilde{Q}(t) := Q(t) - Ct$$
 and $E := E|\widetilde{Q}(t)|$,

we have (with a similar lower deviation bound)

$$\mathbf{P}\{|\widetilde{\mathbf{Q}}(t)|>u\}\leq c\cdot\frac{t^2}{u^4}\cdot E.$$

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Claim: this already implies the $t^{2/3}$ upper bound:

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that is, $E^3 < c \cdot t^2$.

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$$\leq c \cdot \frac{t^2}{E^2} + \frac{1}{2}E,$$

that is, $E^3 < c \cdot t^2$.

$$Var(h_{Ct}(t)) \stackrel{\text{1-hm}}{=} \text{const.} \cdot \mathbf{E}|\mathbf{Q}(t) - Ct|$$

$$= \text{const.} \cdot \mathbf{E} \le c \cdot t^{2/3}.$$

Lower bound

In the upper bound, the relevant orders were

$$u$$
 (deviation of $Q(t)$) $\sim t^{2/3}$, $\varrho - \lambda \sim t^{-1/3}$.

The lower bound works with similar arguments: compare models of which the densities differ by $t^{-1/3}$, and use connections between Q(t), X(t) and heights.

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The critical feature in both the upper bound and lower bound was Q(t) > X(t) (convex) or Q(t) < X(t) (concave).

Coupling results

$H(\varrho)$ is	Feature?	$t^{2/3}$ law
	$H(\varrho)$ is	H(ϱ) is Feature?

Model	$H(\varrho)$ is	Feature?	$t^{2/3}$ law
TASEP			

Model	$H(\varrho)$ is	Feature?	$t^{2/3}$ law
TASEP	concave		
	I	I	

Model	$H(\varrho)$ is	Feature?	$t^{2/3}$ law
TASEP	concave	$Q(t) \leq X(t)$	

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TASEP	concave	$Q(t) \leq X(t)$	proved (BS.)
ASEP	concave	$Q(t) \leq X(t) + Err$	
	I	I	l

Model	$H(\varrho)$ is	Feature?	$t^{2/3}$ law
TASEP	concave	$Q(t) \leq X(t)$	proved (BS.)
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	1		

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rate 1 TAZRP	concave	$Q(t) \leq X(t)$	

$t^{2/3}$ law Model $H(\varrho)$ is Feature? $Q(t) \leq X(t)$ proved (B.-S.) concave **ASEP** $Q(t) \leq X(t) + Err$ proved (B.-S.) concave $Q(t) \leq X(t)$ rate 1 TAZRP concave proved (B.-K.)

Model	$H(\varrho)$ is	Feature?	<i>t</i> ^{2/3} law
TASEP	concave	$Q(t) \leq X(t)$	proved (BS.)
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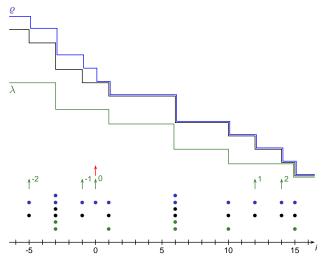
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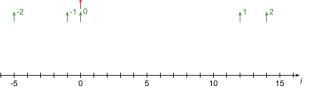
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concave exp rate TAZRP	concave	$Q(t) \le X(t) + Err$	seems ok (BKS.)
convex exp rate TABLP	convex	$Q(t) \ge X(t) - Err$	might work (BKS.)





$$m_{Q}(t) = [\text{the label of } \uparrow \text{ at } Q(t)] = 0$$

 $m_{Q}(t) \leq 0 \Rightarrow Q(t) \leq X(t).$





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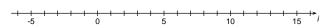




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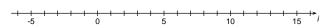




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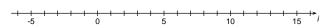




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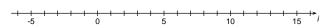




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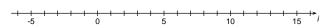




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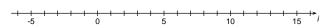




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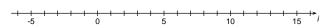




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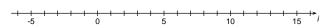




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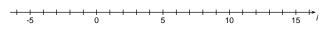




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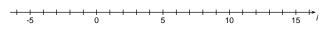




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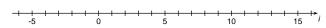
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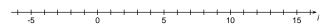




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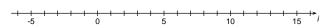




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$$m_{\mathbf{Q}}(t) = [\text{the label of } \uparrow \text{ at } \mathbf{Q}(t)] = 0$$

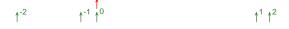
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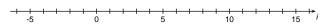




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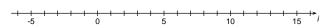




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$$m_{\mathsf{Q}}(t) = [\text{the label of} \uparrow \text{ at } \mathsf{Q}(t)] = 0 \downarrow$$

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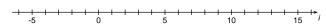




$$m_{\mathsf{Q}}(t) = [\text{the label of } \uparrow \text{ at } \mathsf{Q}(t)] = 0.1$$

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$$m_Q(t) =$$
[the label of † at $Q(t)$] = 0-1
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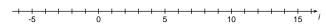




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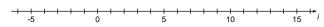




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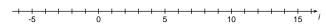




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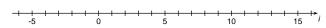
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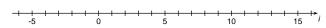




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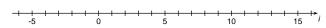
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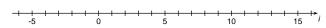




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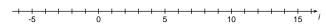
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 $m_{\mathcal{O}}(t) < 0 \Rightarrow Q(t) < X(t)$.





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Generalizing requires finding more models with nice $m_0(t)$ behavior, or handle less nice cases of $m_{\rm O}(t)$ processes. This is subject to future work.

Linear models

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In their cases, we have

$$\lim_{t\to\infty}\frac{\mathsf{Var}(h_{Ct}(t))}{t^{1/2}}=\ldots,$$

even convergence of the finite-dimensional distributions of the $h_{Ct}(t)$ process to Gaussian limits is known (Seppäläinen 2005, Ferrari and Fontes 1998, B., Rassoul-Agha and Seppäläinen 2006).

And there are attractive asymmetric models with nonlinear, nonconvex and nonconcave hydrodynamics:

▶ 2-jump exclusion: $\stackrel{\bullet}{\longleftarrow}$ $\stackrel{\bullet}{\vdash}$ $\stackrel{\bullet}{\vdash}$ $\stackrel{\bullet}{\vdash}$ $\stackrel{\bullet}{\vdash}$ is a cubic polynomial;

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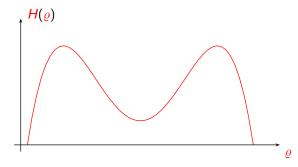
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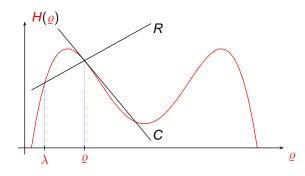
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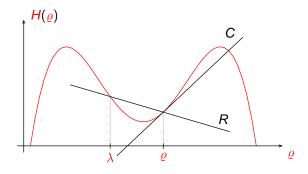
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- A three-state process with variable rates (B. Tóth I. Tóth).

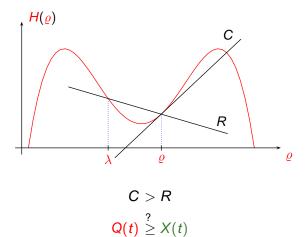




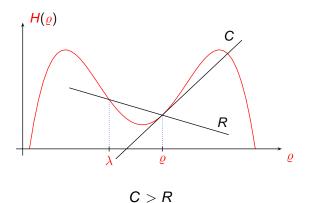
$$C < R$$
 $Q(t) \stackrel{?}{\leq} X(t)$



$$C > R$$
 $Q(t) \stackrel{?}{\geq} X(t)$



Inequality changes with the density...?



$$\frac{Q(t)}{Q(t)} \stackrel{?}{\geq} X(t)$$

Inequality changes with the density...? Any coupling must be very very tricky.

Thank you.