Dependent Double Branching Annihilating Random Walk

Joint with Attila László Nagy

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University of Bristol

New Perspectives in Analysis and Probability University of Sussex 6th March, 2015.

Attractive and non-attractive models

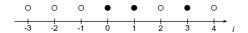
Totally asymmetric simple exclusion process $A \oplus \ominus 0$ model Totally asymmetric zero range process

On large scales Shocks Rarefaction waves

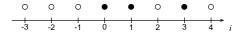
A mean field version

Positive recurrence

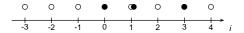
Two words on the proof



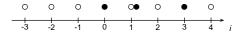
 $\omega_i(0) \sim \text{Bernoulli}(\varrho)$ product distribution.



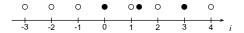
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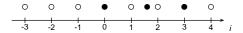
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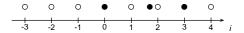
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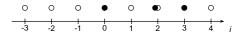
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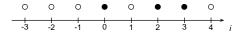
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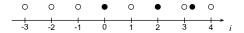
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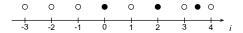
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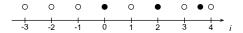
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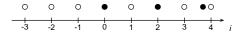
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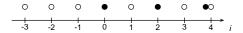
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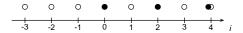
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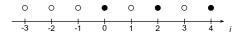
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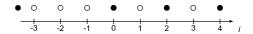
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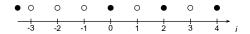
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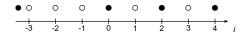
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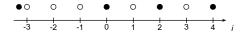
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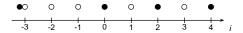
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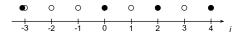
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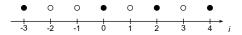
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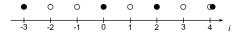
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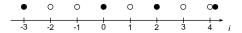
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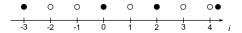
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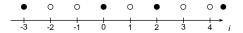
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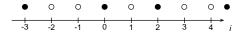
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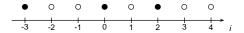
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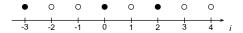
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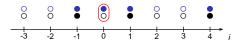
Particles step to the right with rate 1, unless the destination site is occupied.

The Bernoulli(ϱ) product distribution is stationary (and non-reversible) for all $0 \le \varrho \le 1$: $\omega_i(t) \sim \text{Bernoulli}(\varrho)$.

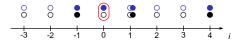
These are the important (= ergodic) stationary distributions.



Stochastic coupling: evolution as close as possible



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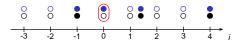
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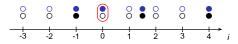
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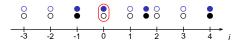
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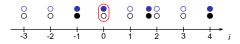
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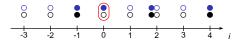
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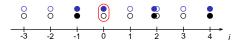
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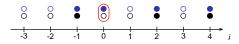
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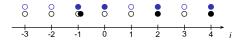
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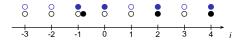
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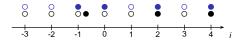
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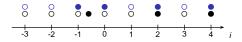
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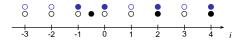
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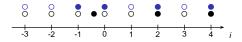
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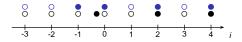
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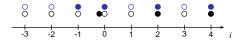
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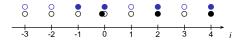
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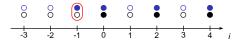
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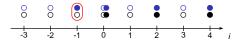
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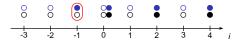
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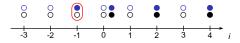
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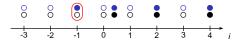
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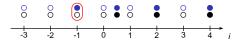
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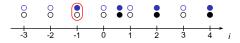
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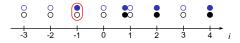
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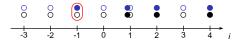
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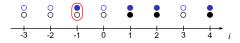
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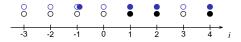
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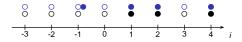
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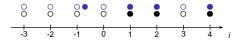
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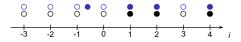
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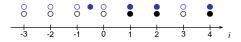
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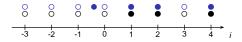
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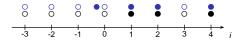
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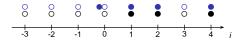
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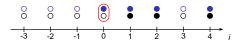
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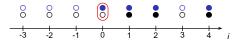


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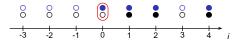
Second class particle. Its position at time t: Q(t).



Stochastic coupling: evolution as close as possible

Second class particle. Its position at time t: Q(t).

$$\begin{aligned} \mathbf{Cov}(\omega_i(t), \, \omega_0(0)) &= \mathbf{E}[\omega_i(t)\omega_0(0) \,|\, \omega_0(0) = 0] \cdot (1-\varrho) \\ &+ \mathbf{E}[\omega_i(t)\omega_0(0) \,|\, \omega_0(0) = 1] \cdot \varrho - \varrho^2 \\ &= \mathbf{E}[\omega_i(t)] \cdot \varrho - \varrho^2. \end{aligned}$$



Stochastic coupling: evolution as close as possible

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 $\mathsf{P}\{\mathsf{Q}(t)=i\}=\mathsf{E}[\omega_i(t)-\omega_i(t)]=\mathsf{E}[\omega_i(t)]-\mathsf{E}[\omega_i(t)].$



Stochastic coupling: evolution as close as possible

Second class particle. Its position at time t: Q(t).

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 $\mathsf{P}\{\mathsf{Q}(t)=i\}=\mathsf{E}[\omega_i(t)-\omega_i(t)]=\mathsf{E}[\omega_i(t)]-\mathsf{E}[\omega_i(t)].$

$$\begin{split} \varrho &= \mathsf{E}[\omega_i(t)] = \mathsf{E}[\omega_i(t) \,|\, \omega_0(0) = 0] \cdot (1 - \varrho) \\ &+ \mathsf{E}[\omega_i(t) \,|\, \omega_0(0) = 1] \cdot \varrho \\ &= \mathsf{E}[\omega_i(t)] \cdot (1 - \varrho) + \mathsf{E}[\omega_i(t)] \cdot \varrho \end{split}$$

$$\mathbf{Cov}(\omega_i(t), \, \omega_0(0)) = \mathbf{E}[\omega_i(t)] \cdot \varrho - \varrho^2. \tag{1}$$

$$\mathsf{P}\{\mathsf{Q}(t)=i\}=\mathsf{E}[\omega_i(t)]-\mathsf{E}[\omega_i(t)]. \tag{2}$$

$$\varrho = \mathbf{E}[\omega_i(t)] \cdot (1 - \varrho) + \mathbf{E}[\omega_i(t)] \cdot \varrho.$$
 (3)

$$\mathbf{Cov}(\omega_i(t),\,\omega_0(0)) = \mathbf{E}[\omega_i(t)] \cdot \varrho - \varrho^2. \tag{1}$$

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$$\varrho = \mathbf{E}[\omega_i(t)] \cdot (1 - \varrho) + \mathbf{E}[\omega_i(t)] \cdot \varrho.$$
 (3)

So,

 $\mathbf{Cov}(\omega_i(t), \omega_0(0)) \stackrel{(1)}{=} \varrho \cdot (\mathbf{E}[\omega_i(t)] - \varrho)$

$$\mathbf{Cov}(\omega_i(t), \, \omega_0(0)) = \mathbf{E}[\omega_i(t)] \cdot \varrho - \varrho^2. \tag{1}$$

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$$\varrho = \mathsf{E}[\omega_i(t)] \cdot (1-\varrho) + \mathsf{E}[\omega_i(t)] \cdot \varrho.$$
 (3)

So,

$$\begin{aligned} \mathbf{Cov}(\omega_i(t), \, \omega_0(0)) \stackrel{(1)}{=} \varrho \cdot \left(\mathbf{E}[\omega_i(t)] - \varrho \right) \\ \stackrel{(3)}{=} \varrho(1 - \varrho) \cdot \left(\mathbf{E}[\omega_i(t)] - \mathbf{E}[\omega_i(t)] \right) \end{aligned}$$

$$\mathbf{Cov}(\omega_i(t), \, \omega_0(0)) = \mathbf{E}[\omega_i(t)] \cdot \varrho - \varrho^2. \tag{1}$$

$$\mathsf{P}\{\mathsf{Q}(t)=i\}=\mathsf{E}[\omega_i(t)]-\mathsf{E}[\omega_i(t)]. \tag{2}$$

$$\varrho = \mathbf{E}[\omega_i(t)] \cdot (1 - \varrho) + \mathbf{E}[\omega_i(t)] \cdot \varrho.$$
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So,

$$\begin{aligned} \mathbf{Cov}(\omega_i(t), \, \omega_0(0)) &\stackrel{(1)}{=} \varrho \cdot \left(\mathbf{E}[\omega_i(t)] - \varrho \right) \\ &\stackrel{(3)}{=} \varrho(1 - \varrho) \cdot \left(\mathbf{E}[\omega_i(t)] - \mathbf{E}[\omega_i(t)] \right) \\ &\stackrel{(2)}{=} \varrho(1 - \varrho) \cdot \mathbf{P}\{\mathbf{Q}(t) = i\}. \end{aligned}$$

$$\mathbf{Cov}(\omega_i(t), \, \omega_0(0)) = \mathbf{E}[\omega_i(t)] \cdot \varrho - \varrho^2. \tag{1}$$

$$\mathsf{P}\{\mathsf{Q}(t)=i\}=\mathsf{E}[\omega_i(t)]-\mathsf{E}[\omega_i(t)]. \tag{2}$$

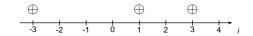
$$\varrho = \mathsf{E}[\omega_i(t)] \cdot (1 - \varrho) + \mathsf{E}[\omega_i(t)] \cdot \varrho.$$
 (3)

So,

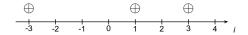
$$\begin{aligned} \mathbf{Cov}(\omega_i(t), \, \omega_0(0)) \stackrel{(1)}{=} \varrho \cdot \left(\mathbf{E}[\omega_i(t)] - \varrho \right) \\ \stackrel{(3)}{=} \varrho(1 - \varrho) \cdot \left(\mathbf{E}[\omega_i(t)] - \mathbf{E}[\omega_i(t)] \right) \\ \stackrel{(2)}{=} \varrho(1 - \varrho) \cdot \mathbf{P}\{Q(t) = i\}. \end{aligned}$$

The second class particle traces information propagation.

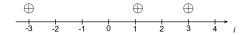
$A \oplus \ominus 0$ model



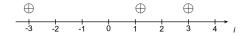
\oplus to the right: rate 1



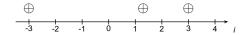
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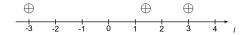
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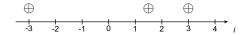
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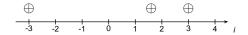
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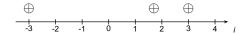
\oplus to the right: rate 1



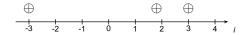
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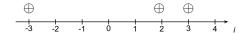
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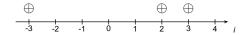
\oplus to the right: rate 1



\oplus to the right: rate 1

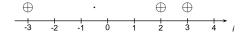


\oplus to the right: rate 1

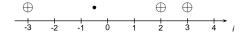


$A \oplus \ominus 0$ model

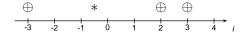
pair creation from vacuum: rate c



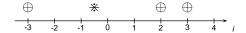
pair creation from vacuum: rate c



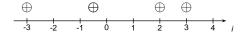
pair creation from vacuum: rate c



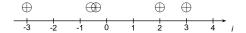
pair creation from vacuum: rate c



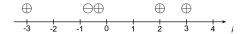
pair creation from vacuum: rate c



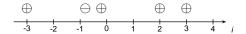
pair creation from vacuum: rate c



pair creation from vacuum: rate c



pair creation from vacuum: rate c



pair creation from vacuum: rate c



pair creation from vacuum: rate c



\ominus to the left: rate 1



\ominus to the left: rate 1



\ominus to the left: rate 1



\ominus to the left: rate 1



\ominus to the left: rate 1



\ominus to the left: rate 1



\ominus to the left: rate 1



\ominus to the left: rate 1



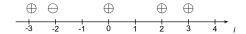
\ominus to the left: rate 1



\ominus to the left: rate 1



annihilation: rate 2



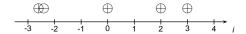
annihilation: rate 2



annihilation: rate 2



annihilation: rate 2



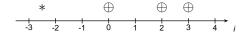
annihilation: rate 2



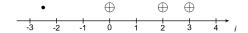
annihilation: rate 2



annihilation: rate 2





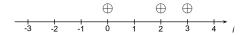


annihilation: rate 2



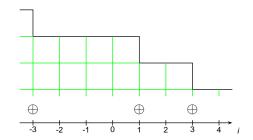
annihilation: rate 2





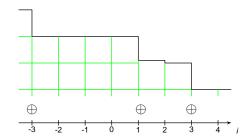
 $\omega_i = -1, 0, 1$: a family of product initial distribution.

Those product distributions are stationary (and non-reversible).



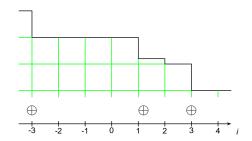
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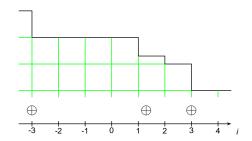
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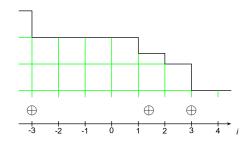
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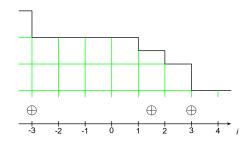
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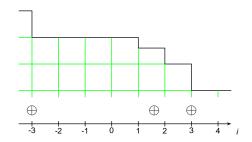
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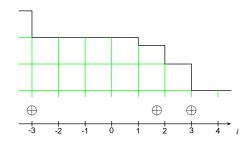
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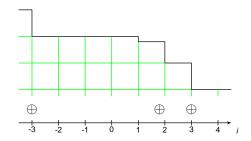
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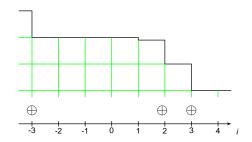
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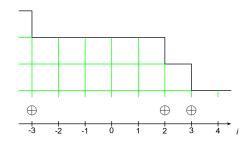
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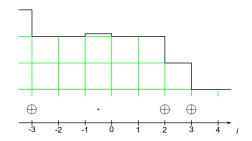
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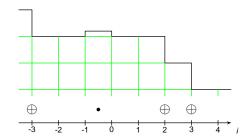
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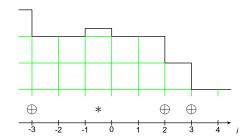
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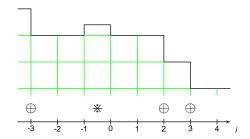
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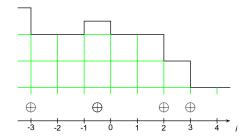
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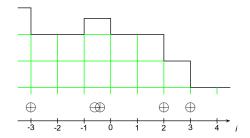
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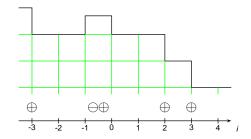
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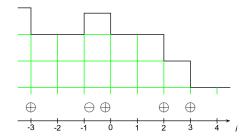
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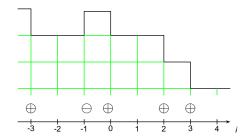
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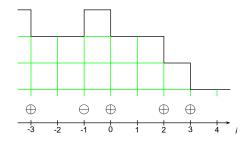
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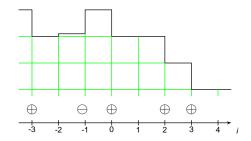
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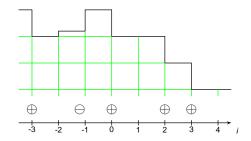
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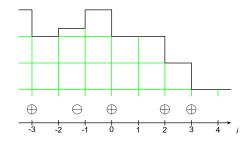
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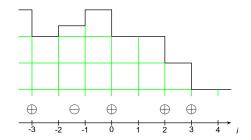
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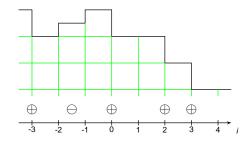
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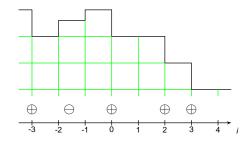
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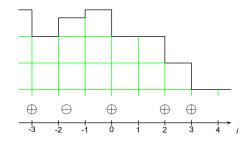
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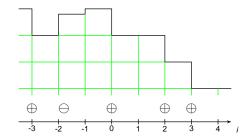
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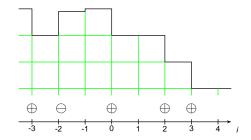
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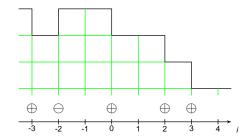
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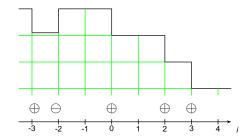
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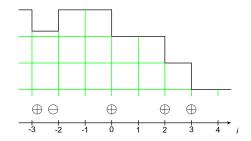
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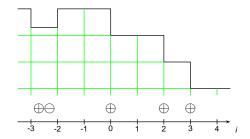
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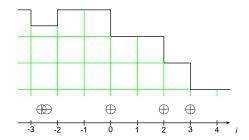
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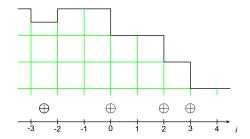
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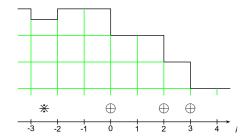
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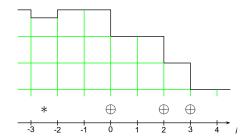
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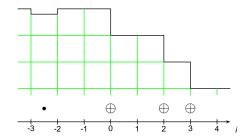
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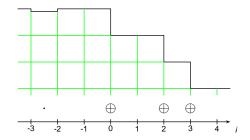
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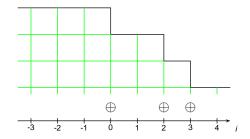
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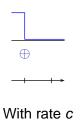


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The second class particle

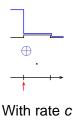
... works much like in TASEP for $c \le 1$. The interesting case:

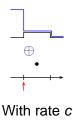


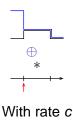
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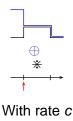


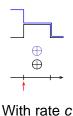
With rate c

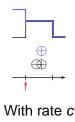


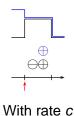


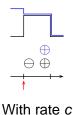


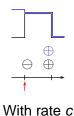








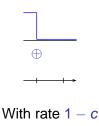




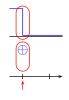
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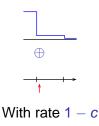
With rate c

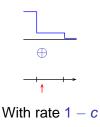


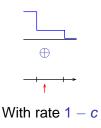
... works much like in TASEP for $c \le 1$. The interesting case:

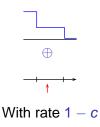


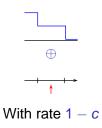
With rate 1 - c

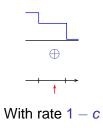


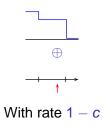


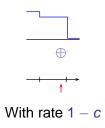


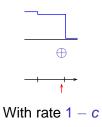




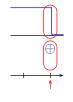






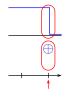


... works much like in TASEP for $c \leq 1$. The interesting case:



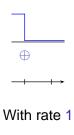
With rate 1 - c

... works much like in TASEP for $c \leq 1$. The interesting case:

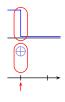


With rate 1 - c

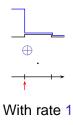
Attractivity

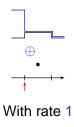


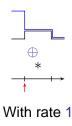
But, for c > 1:

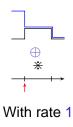


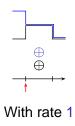
With rate 1

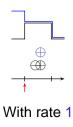


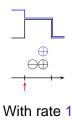


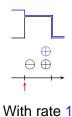


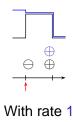










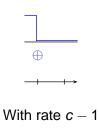


But, for c > 1:

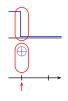


With rate 1

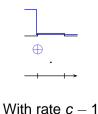
But, for c > 1:



But, for c > 1:



But, for c > 1:



But, for c > 1:



But, for c > 1:



But, for c > 1:



But, for c > 1:



But, for c > 1:



But, for *c* > 1:



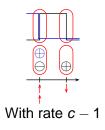
But, for *c* > 1:



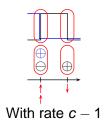
But, for c > 1:



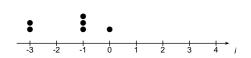
But, for *c* > 1:

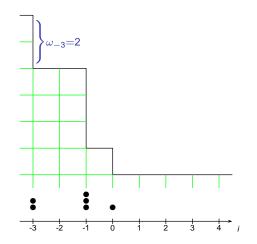


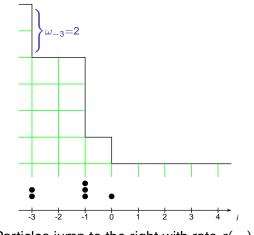
But, for *c* > 1:



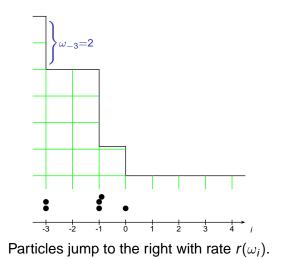
Non-attractivity

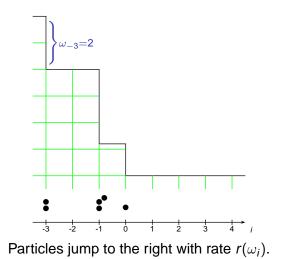


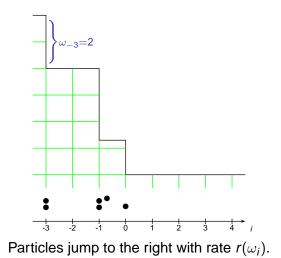


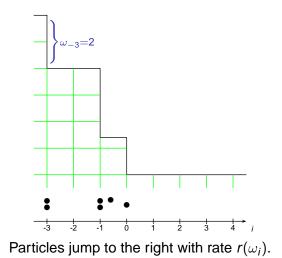


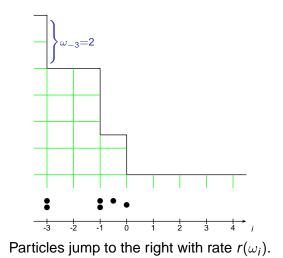
Particles jump to the right with rate $r(\omega_i)$.

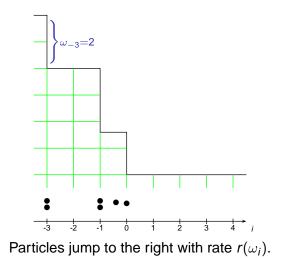


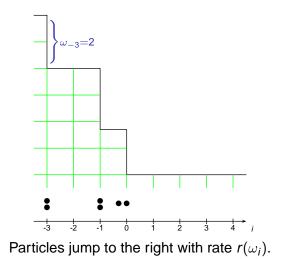


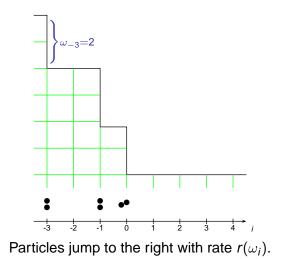


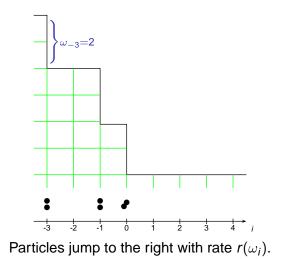


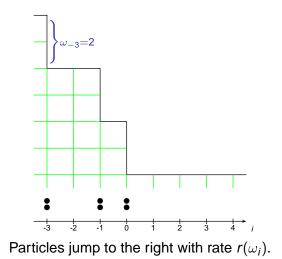


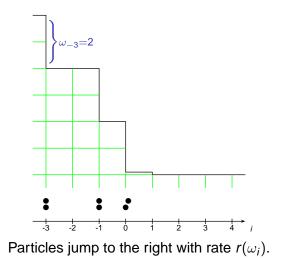


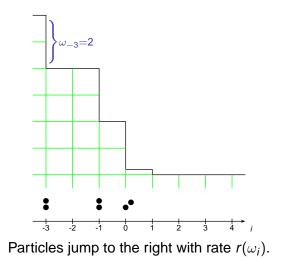


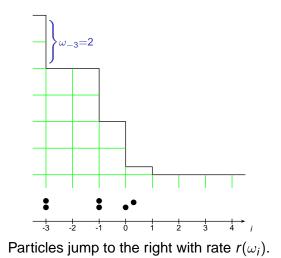


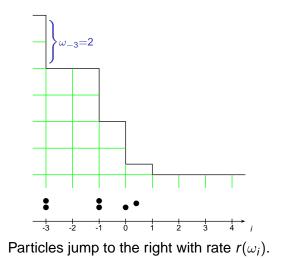


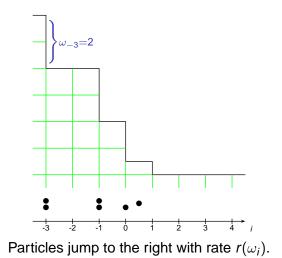


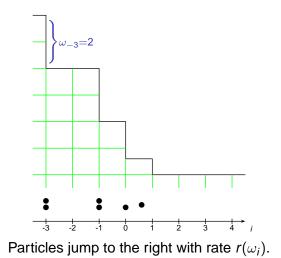


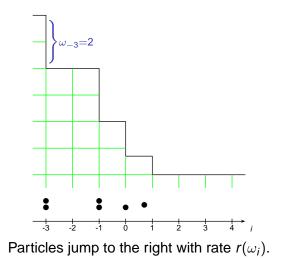


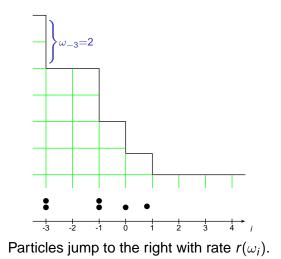


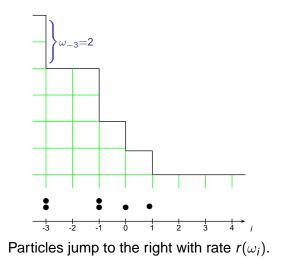


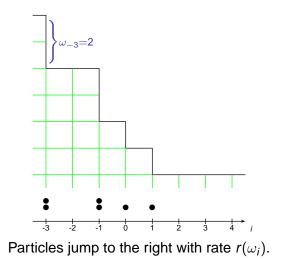


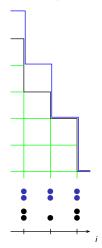


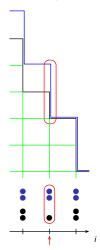


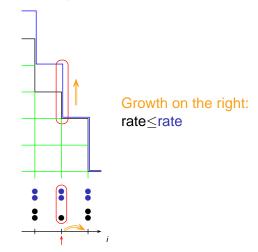


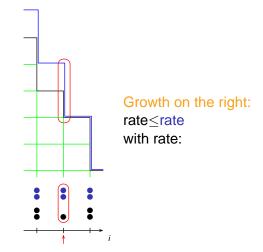


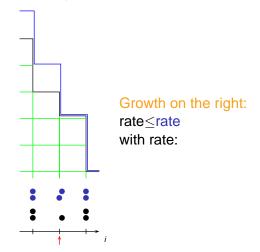


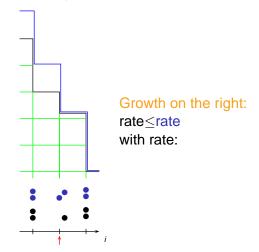


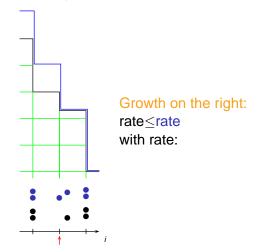


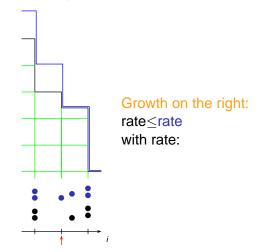


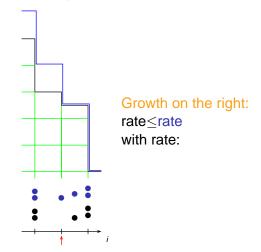


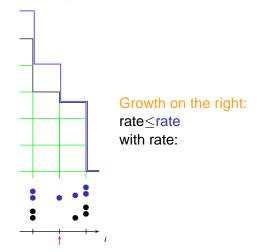


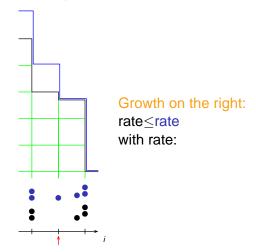


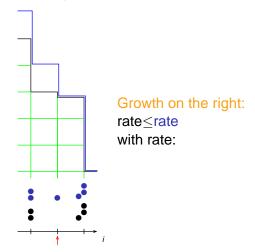


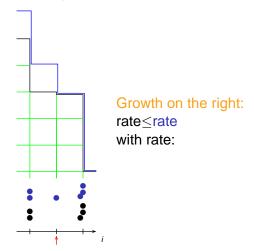


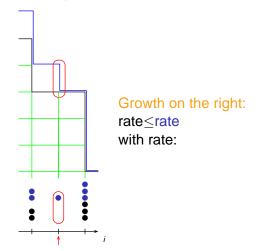


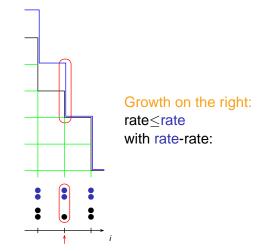


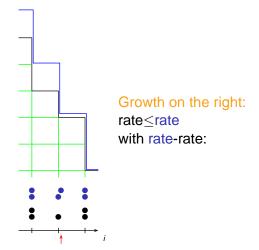


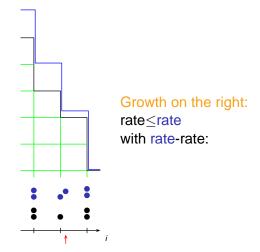


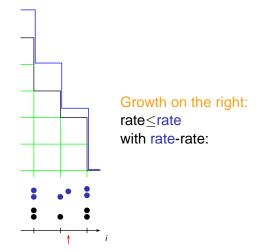


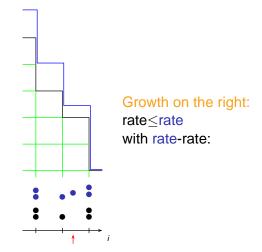


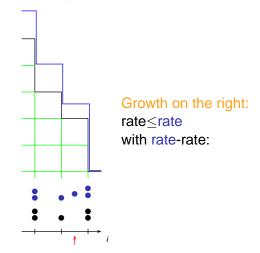


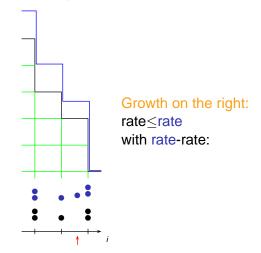


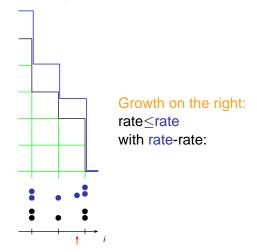


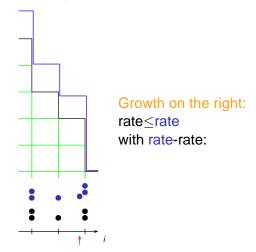


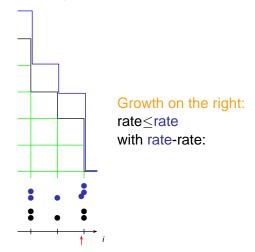


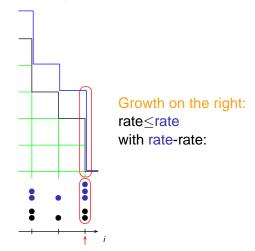


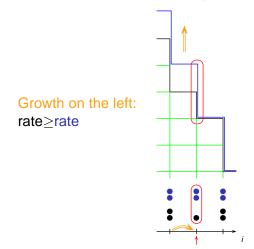


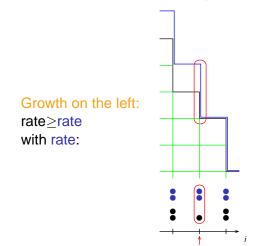


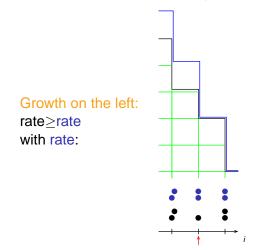


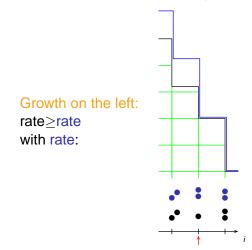


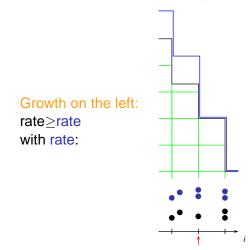


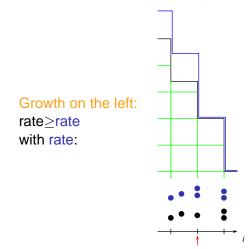


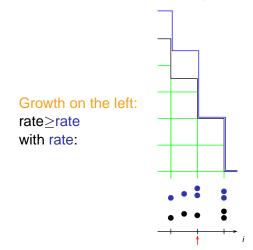


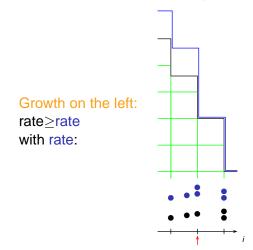


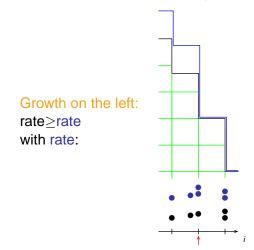


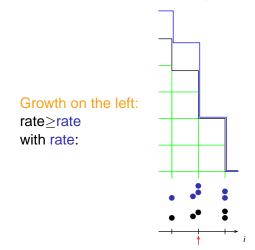


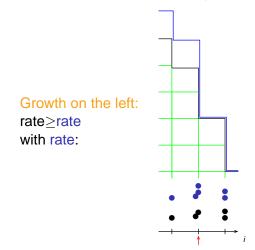


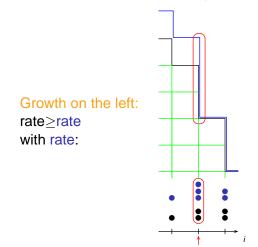


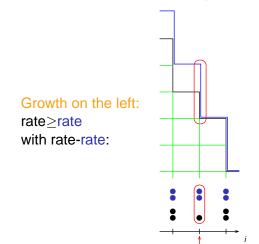


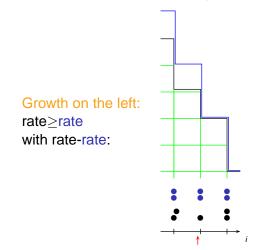


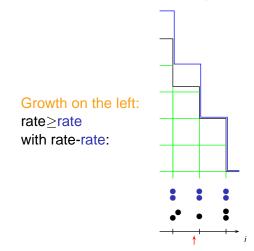


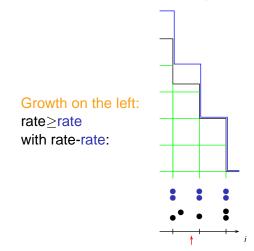


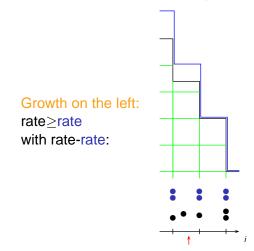


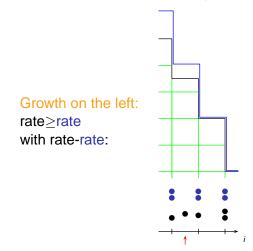


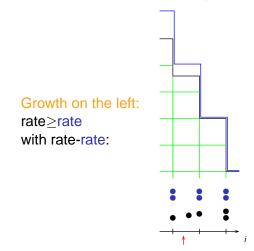


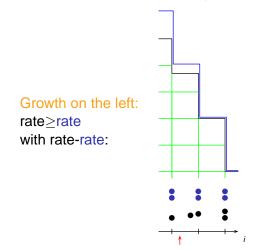


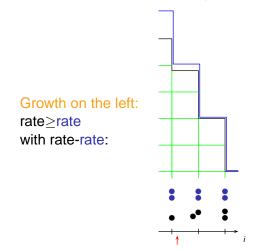


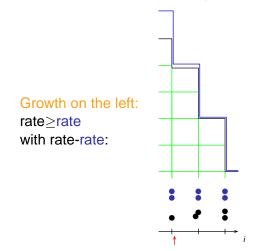


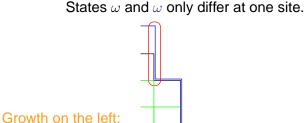




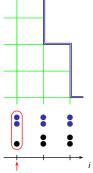


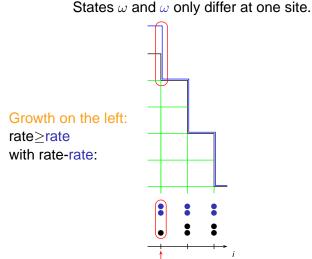




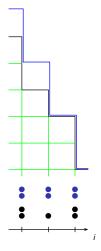


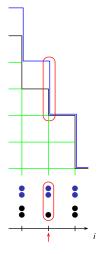
rate≥rate with rate-rate:

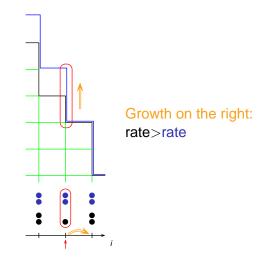


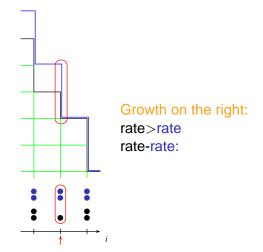


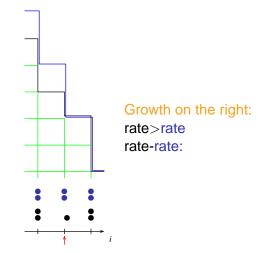
A single discrepancy t, the second class particle, is conserved.

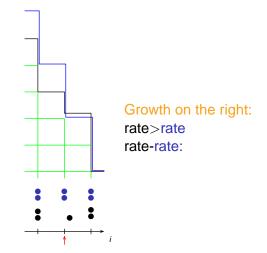


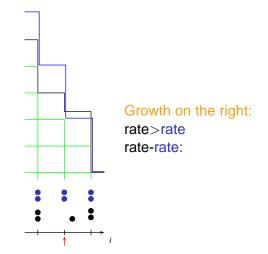


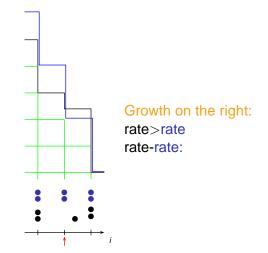


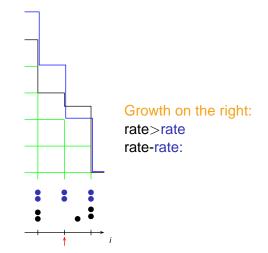


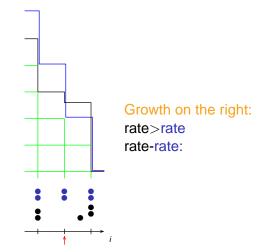


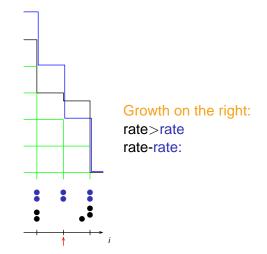


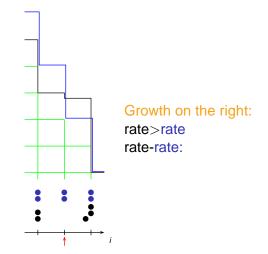


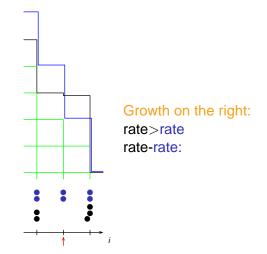


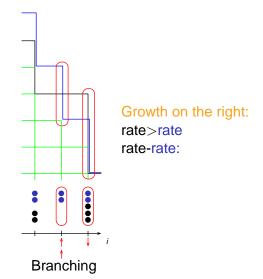


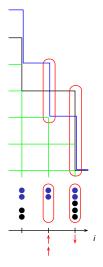


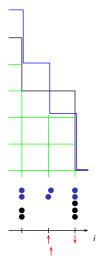


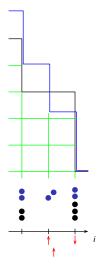


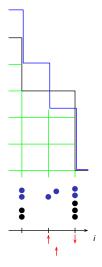


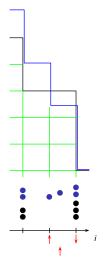


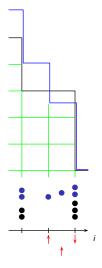


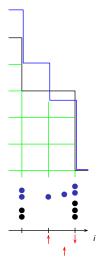


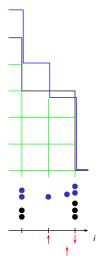


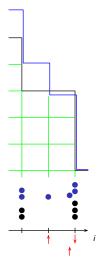


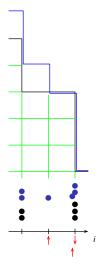


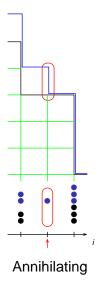


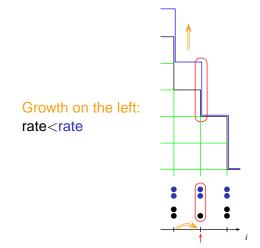


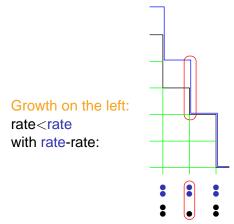


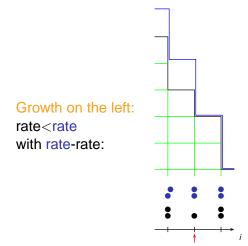


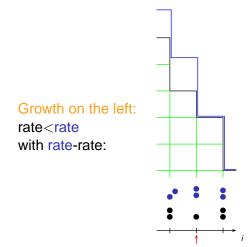


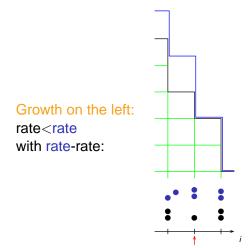


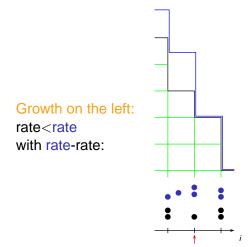


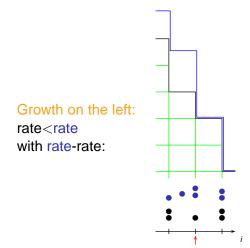


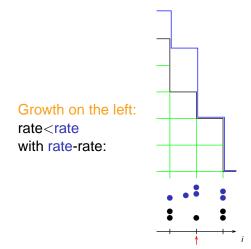


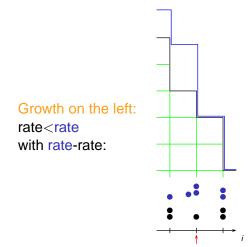


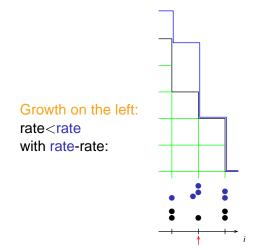


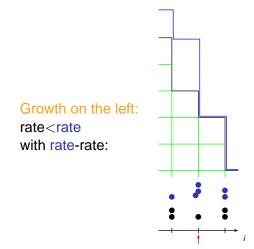


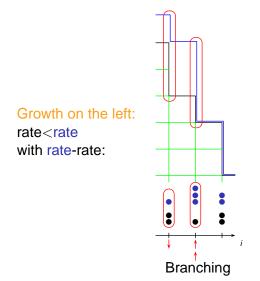


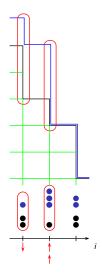


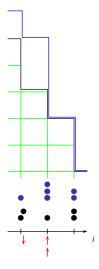


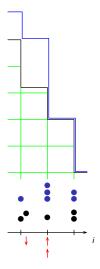


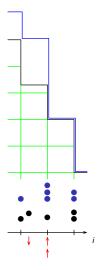


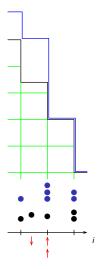


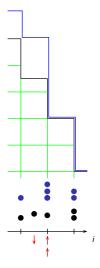


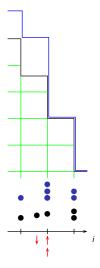


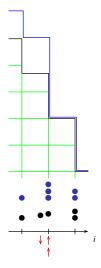


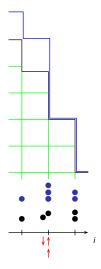


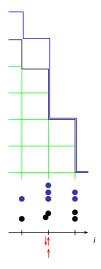


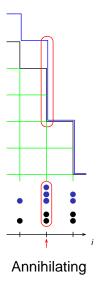












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$$= \lim_{t \to 0} \frac{\mathbf{P}\{\omega_{i-1} = 1, \ \omega_i = 0\}t - \mathbf{P}\{\omega_i = 1, \ \omega_{i+1} = 0\}t + \mathfrak{o}(t)}{t}$$

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$$= \varrho_{i-1}[1 - \varrho_{i}] - \varrho_{i}[1 - \varrho_{i+1}].$$

Let us now allow the density to change slowly in space. The change of density at position *i*:

$$\frac{\mathrm{d}}{\mathrm{d}t}\varrho_i = \frac{\mathrm{d}}{\mathrm{d}t}\mathbf{E}\omega_i$$

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$$\frac{\partial}{\partial T}\hat{\varrho} + \frac{\partial}{\partial X}[\hat{\varrho}(1-\hat{\varrho})] = 0$$
 (Burgers eq.).

Burgers eq.: characteristics

 $\frac{\partial}{\partial T}\hat{\varrho} + \frac{\partial}{\partial X}[\hat{\varrho}(1-\hat{\varrho})] = 0 \qquad \text{Burgers eq.: nonlinear PDE.}$

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Characteristics: find a path X(T) where $\hat{\varrho}(T, X(T))$ is a constant:

$$\frac{\mathrm{d}}{\mathrm{d}T}\hat{\varrho}\big(T,\,X(T)\big)=0$$

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The characteristic velocity: $\dot{X}(T) = 1 - 2\hat{\varrho}$.

Burgers eq.: characteristics

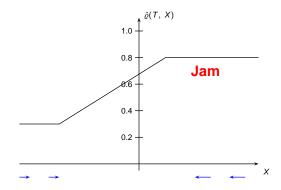
 $\frac{\partial}{\partial T}\hat{\varrho} + \frac{\partial}{\partial X}[\hat{\varrho}(1-\hat{\varrho})] = 0 \qquad \text{Burgers eq.: nonlinear PDE.}$

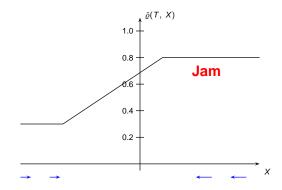
Characteristics: find a path X(T) where $\hat{\varrho}(T, X(T))$ is a constant:

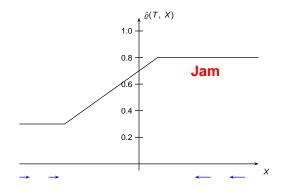
$$\frac{\mathrm{d}}{\mathrm{d}T}\hat{\varrho}(T, X(T)) = 0$$
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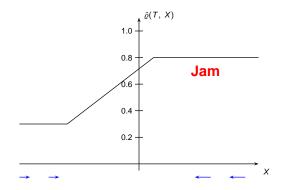
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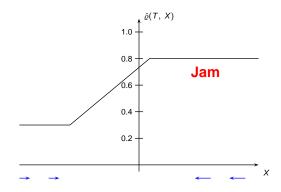
Second class particles are known to follow the characteristics.

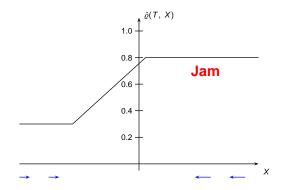


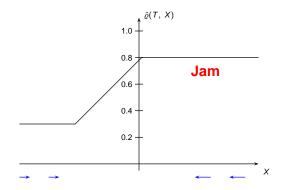


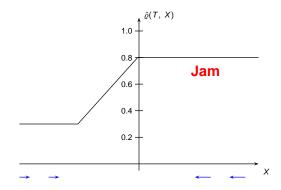


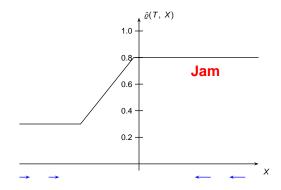


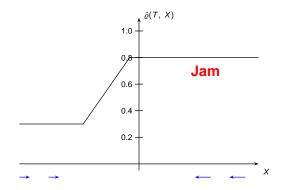


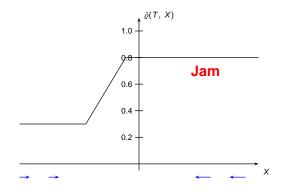


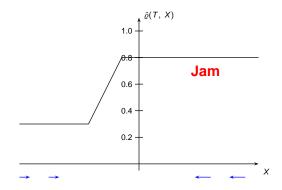


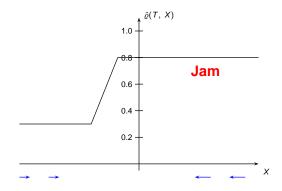


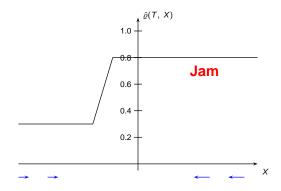


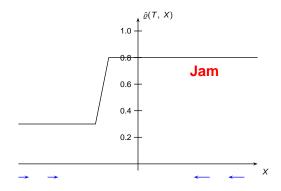


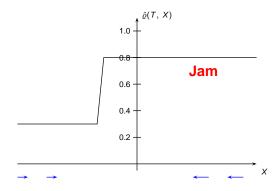


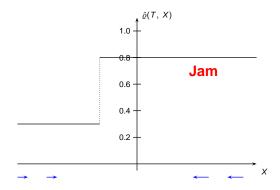


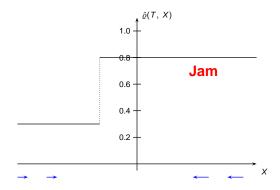


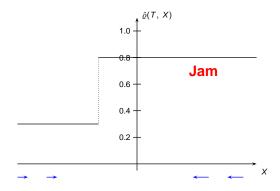


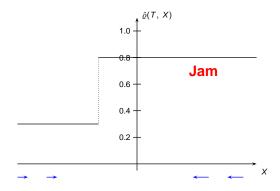


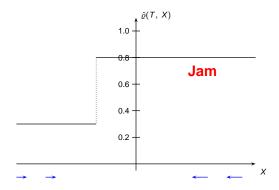


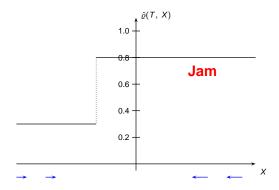


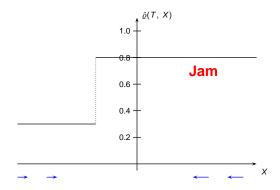


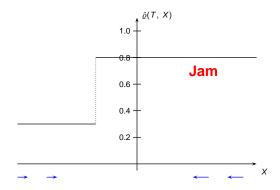


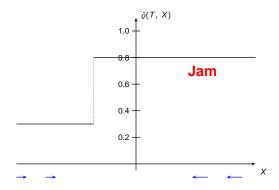


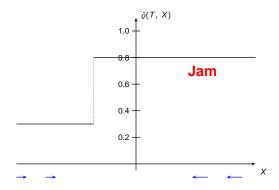


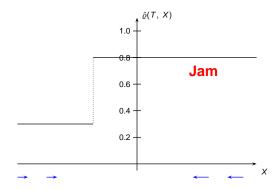


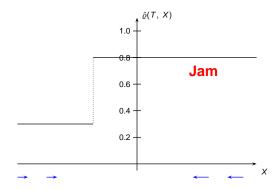


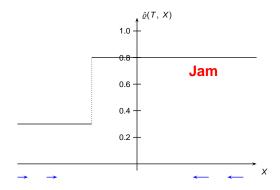


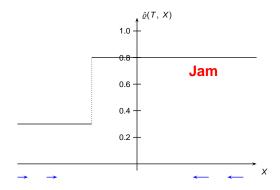


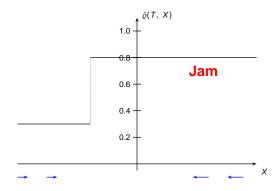


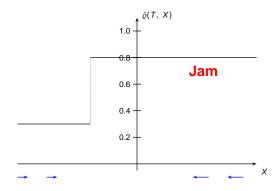


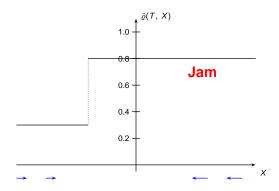


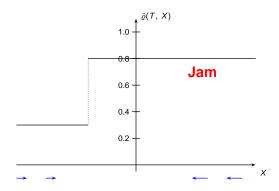


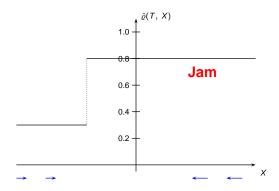


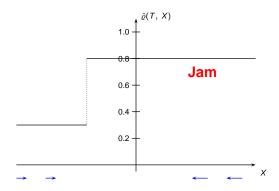


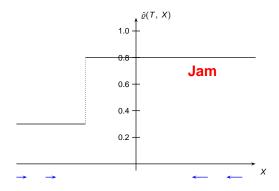


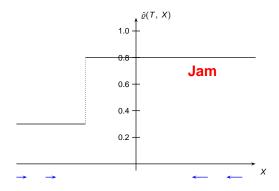


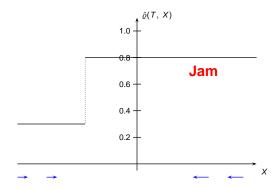


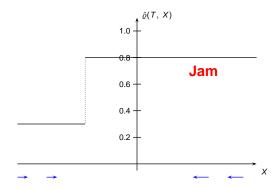


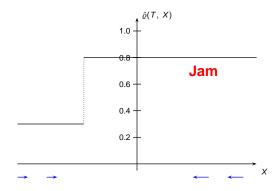


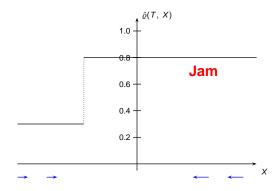


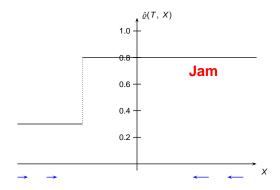


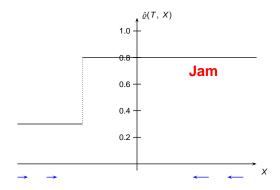


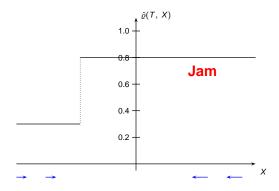


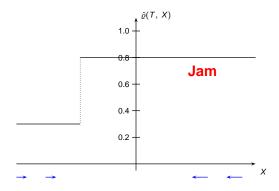


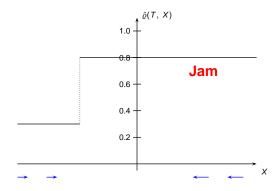


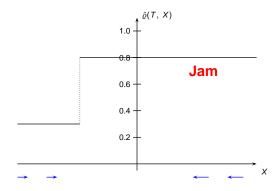


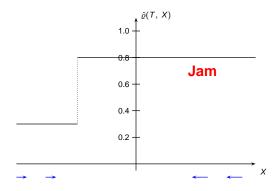


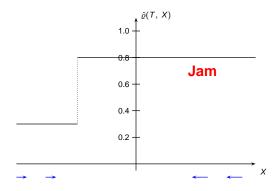


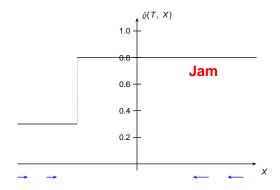


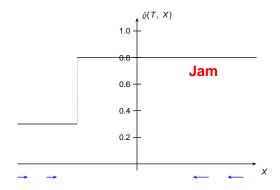


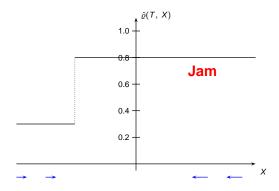


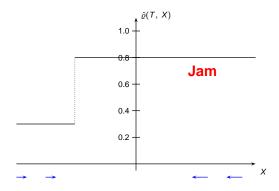


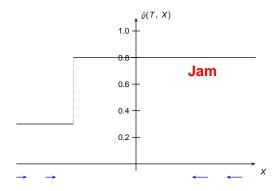


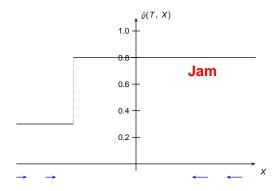


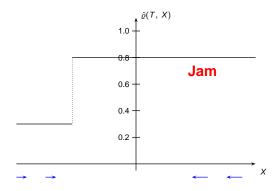


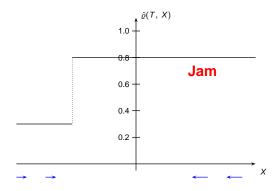


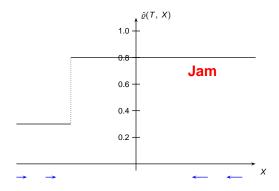


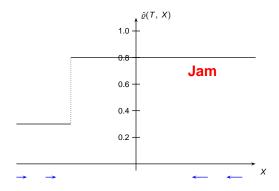


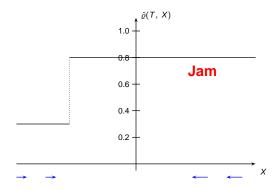


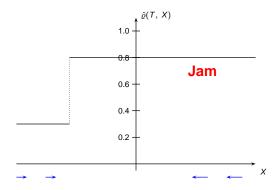






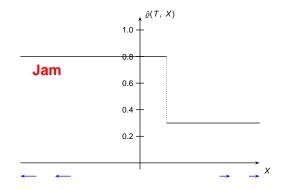


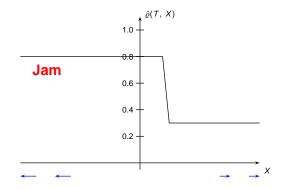


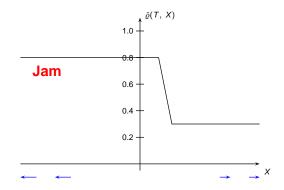


 $\dot{X}(T) = 1 - 2\hat{\varrho}$

Shock

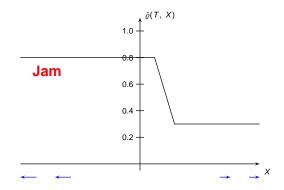


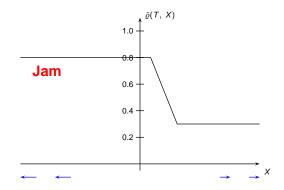


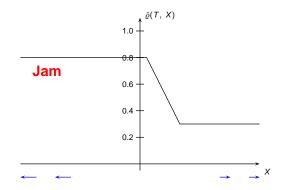


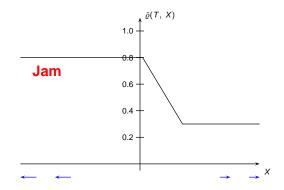
ocks Rarefaction waves

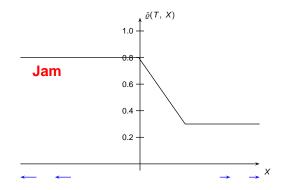
On large scales

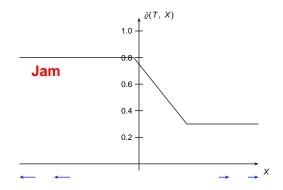


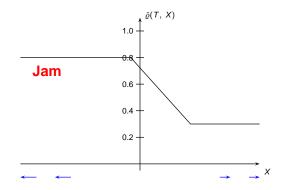


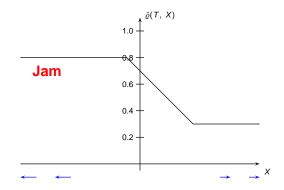


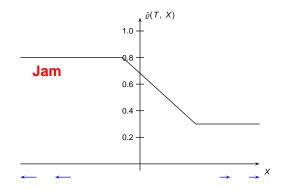


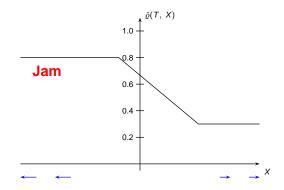


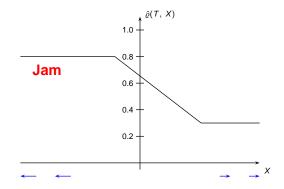


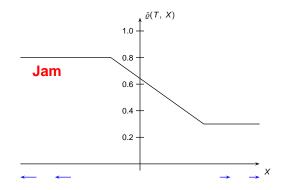


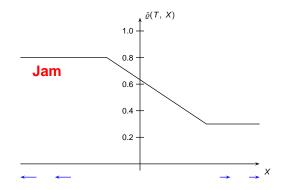


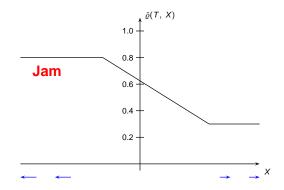


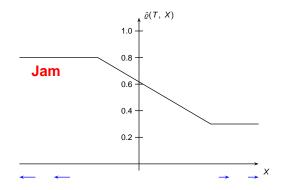


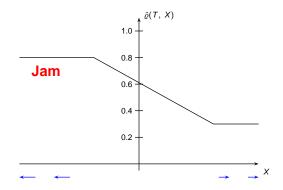


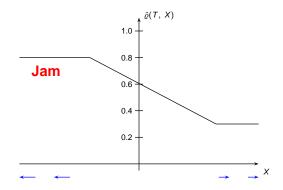


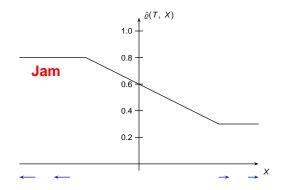


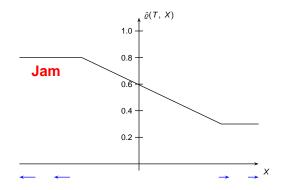


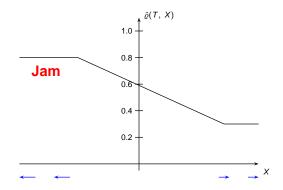


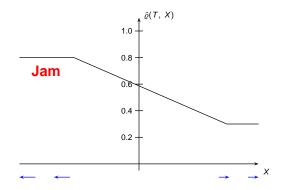


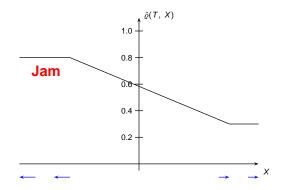


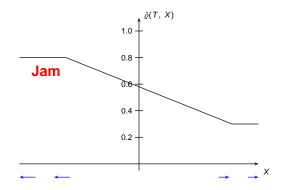


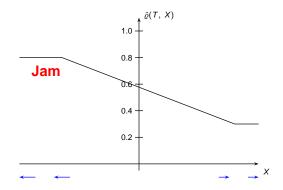


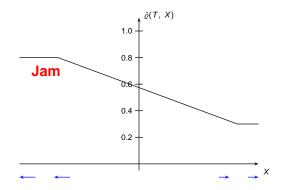


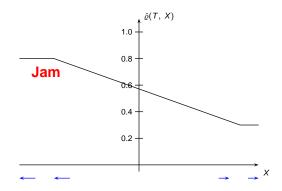


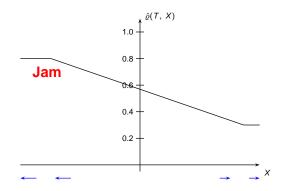


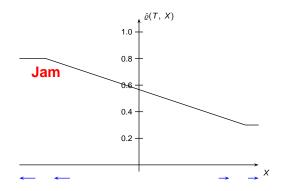


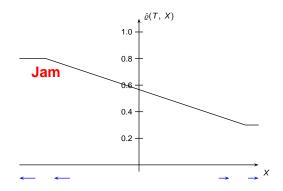












 $\dot{X}(T) = 1 - 2\hat{\varrho}$

Rarefaction wave

The second class particle: non-attractive case

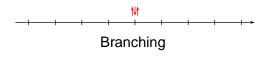
We are facing a

- nearest neighbour
- parity conserving
- branching
- annihilating process
- on the dynamic background of first class particles.

The aim is to control the number of \dagger and \downarrow 's. Idea from Bálint Tóth.

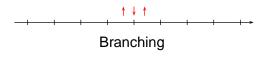
∽ homog2.avi







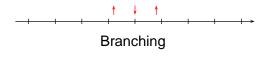














































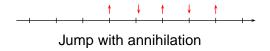


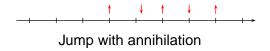


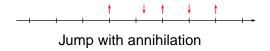


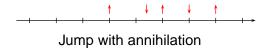


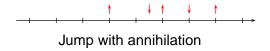


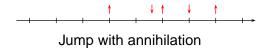


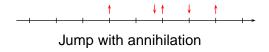


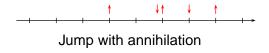


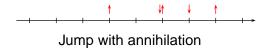


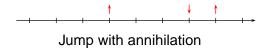


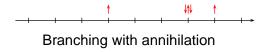


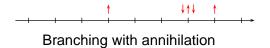


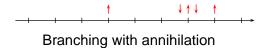




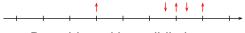








A model we can say something about:



Branching with annihilation

A model we can say something about:



Branching with annihilation

A model we can say something about:



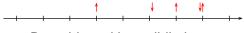
A model we can say something about:



A model we can say something about:

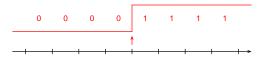


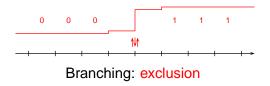
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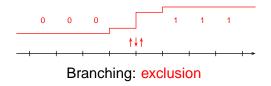


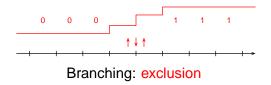
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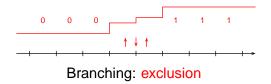


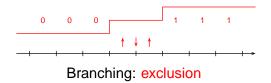


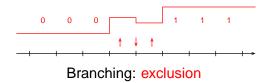


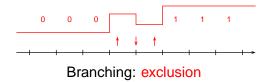


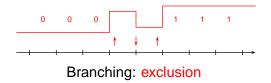


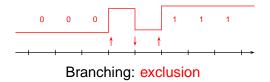


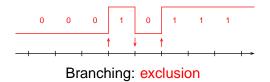


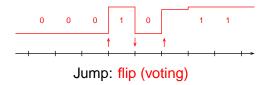


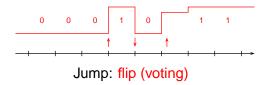


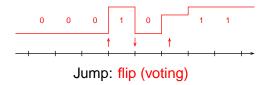


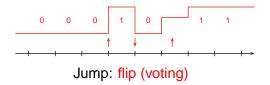


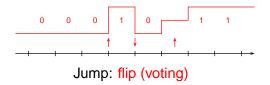


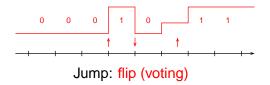


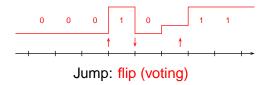


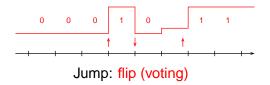


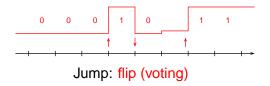


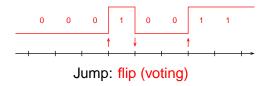


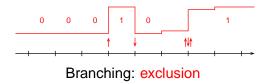


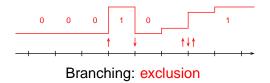


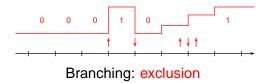


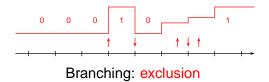


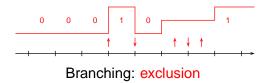


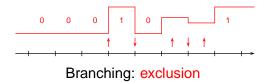


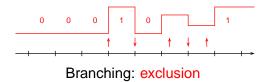


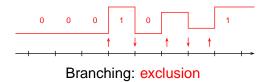




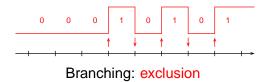




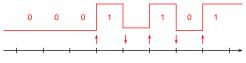


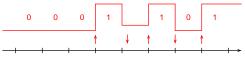






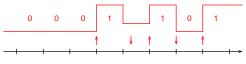
A model we can say something about:



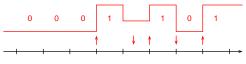


Jump with annihilation: flip (voting)

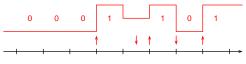
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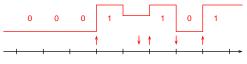
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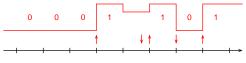
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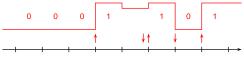
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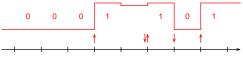
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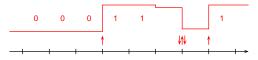


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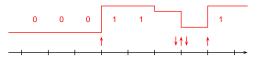




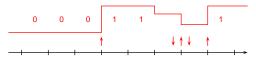
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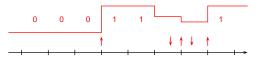
Branching with annihilation: exclusion



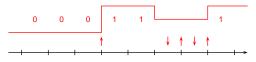
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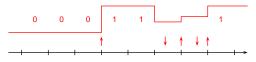
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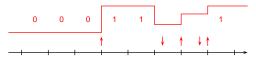
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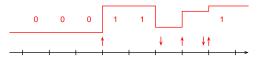
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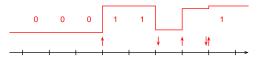
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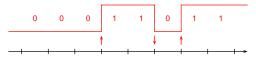
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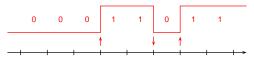


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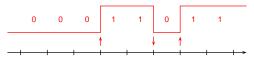
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Double branching-annihilating random walks (DBARW)

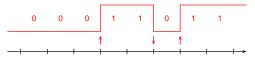
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Branching with annihilation: exclusion

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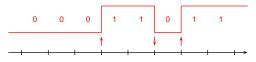
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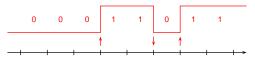
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Question: Is the process, as seen by the leftmost 1, recurrent?

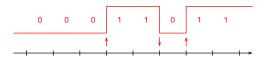


First instance of DBARW we could find in the literature: A. Sudbury '90. Positive recurrence: V. Belitsky, P.A. Ferrari, M.V. Menshikov and S.Y. Popov '01; A. Sturm and J.M. Swart '08. *Results are very sensitive to the details of branching.*



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But: true second class particles interact (*common background* of first class particles).

→ Repeat the Sturm-Swart proof with configuration dependent jump rates. Jump rates can depend on the whole configuration.

Conditions on the jumping and branching rates:

Translation invariance.

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- ▶ Bounds on the difference for branching rates of *t*'s and *t*'s.
- Weak dependence on particles far away.
- No repulsion in the jumping rates between particles. (A bit of repulsion locally is still OK.)

Theorem

Then, starting from a single t:

 The process takes finitely many steps in finite time (construction).

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- The process as seen from the leftmost t is positive recurrent.
- The stationary distribution sees a finite expected number of particles.
- (Extension of all this to non nearest neighbour symmetric branching.)

An example

- Branching rates: constant.
- Jump rate to the right:

$$\frac{1}{2} + \sum_{\text{particle on right}} \frac{1}{\text{distance}^{\alpha}},$$

jump rate to the left:



 $\alpha > 1.$

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Unfortunately we do not seem to be there yet... This is not covered at the moment. But a small modification that respects parity in a peculiar way works.

Another example

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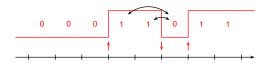
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This one is fine.

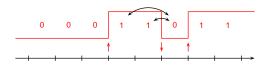
Two words on the proof



Main tool 1: the number of inversions, i.e., wrongly ordered 1-0 pairs.

If there are too many of them, the generator is negative on the number of these pairs.

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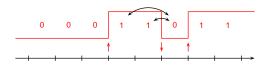
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Thank you.