Order of current variance in the simple exclusion process

Márton Balázs

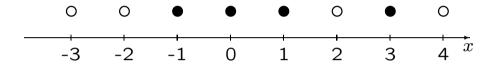
(University of Wisconsin - Madison)
(Budapest University of Technology and Economics)

Joint work with

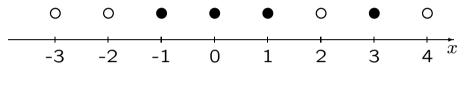
Timo Seppäläinen (University of Wisconsin - Madison)

Prague, December 4, 2006.

- 1. ASEP: Interacting particles
- 2. ASEP: Surface growth
 - 3. Growth fluctuations
 - 4. The second class particle
 - 5. The upper bound
 - 6. The lower bound
 - 7. Open questions



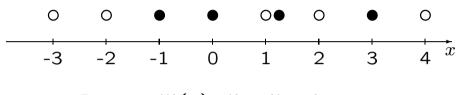
Bernoulli(ϱ) distribution



Bernoulli(ϱ) distribution

Particles try to jump

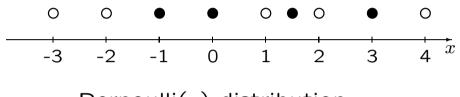
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

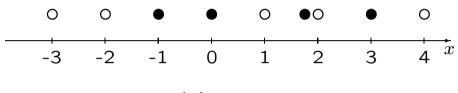
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

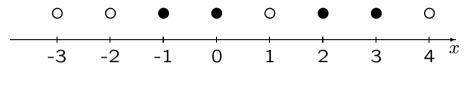
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

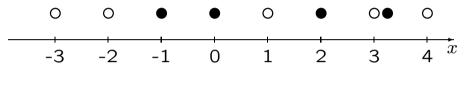
to the right with rate p, to the left with rate q = 1 - p < p.



 $\mathsf{Bernoulli}(\underline{\varrho}) \ \mathsf{distribution}$

Particles try to jump

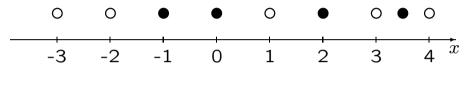
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

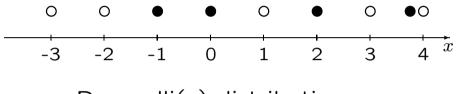
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

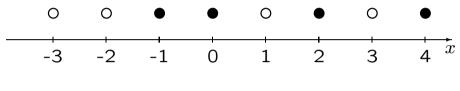
to the right with rate p, to the left with rate q = 1 - p < p.



 $\mathsf{Bernoulli}(\underline{\varrho}) \ \mathsf{distribution}$

Particles try to jump

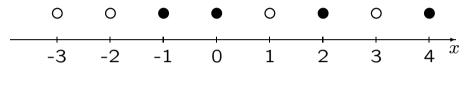
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

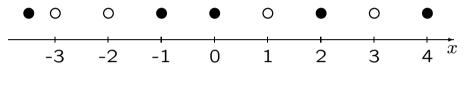
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

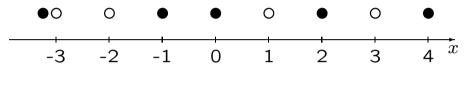
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

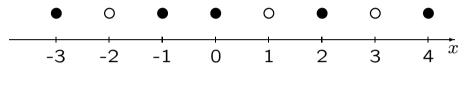
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

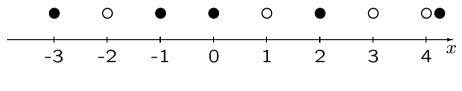
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

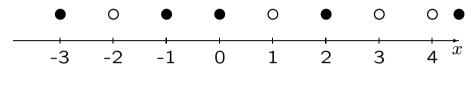
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

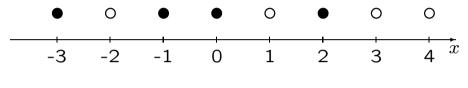
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

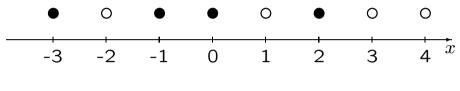
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

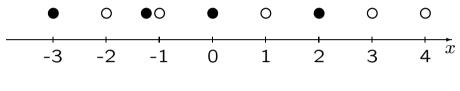
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

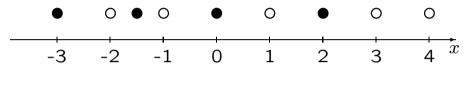
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

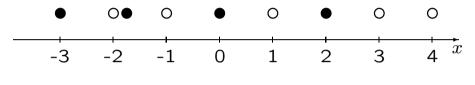
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

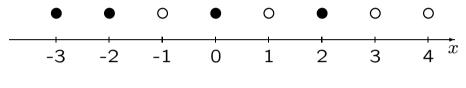
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

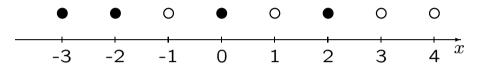
to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

to the right with rate p, to the left with rate q = 1 - p < p.



Bernoulli(ϱ) distribution

Particles try to jump

to the right with rate p, to the left with rate q = 1 - p < p.

The jump is suppressed if the destination site is occupied by another particle.

The Bernoulli(ϱ) distribution is time-stationary for any $(0 \le \varrho \le 1)$. Any translation-invariant stationary distribution is a mixture of Bernoullis.

Let T and X be some large-scale time and space parameters. \longrightarrow Set initially $\varrho = \varrho(T=0,X)$ to be the density at position $x=X/\varepsilon$. (Changes on the large scale.)

- \leadsto Set initially $\varrho = \varrho(T=0,X)$ to be the density at position $x=X/\varepsilon$. (Changes on the large scale.)
- $\rightsquigarrow \varrho(T,X)$ is the density of particles after a long time $t=T/\varepsilon$ at position $x=X/\varepsilon$.

Let T and X be some large-scale time and space parameters.

 \leadsto Set initially $\varrho = \varrho(T=0,X)$ to be the density at position $x=X/\varepsilon$. (Changes on the large scale.)

 $\sim \varrho(T,X)$ is the density of particles after a long time $t=T/\varepsilon$ at position $x=X/\varepsilon$. It satisfies, with a:=p-q,

$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1 - \varrho)] = 0 \quad \text{(inviscid Burgers)}$$

Let T and X be some large-scale time and space parameters.

 \leadsto Set initially $\varrho = \varrho(T=0,X)$ to be the density at position $x=X/\varepsilon$. (Changes on the large scale.)

 $\sim \varrho(T,X)$ is the density of particles after a long time $t=T/\varepsilon$ at position $x=X/\varepsilon$. It satisfies, with a:=p-q,

$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1 - \varrho)] = 0 \quad \text{(inviscid Burgers)}$$

$$\frac{\partial}{\partial T} \varrho + a[1 - 2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$

- \leadsto Set initially $\varrho = \varrho(T=0,X)$ to be the density at position $x=X/\varepsilon$. (Changes on the large scale.)
- $\sim \varrho(T,X)$ is the density of particles after a long time $t=T/\varepsilon$ at position $x=X/\varepsilon$. It satisfies, with a:=p-q,

$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1 - \varrho)] = 0 \quad \text{(inviscid Burgers)}$$

$$\frac{\partial}{\partial T} \varrho + a[1 - 2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$

$$\frac{\mathrm{d}}{\mathrm{d}T} \varrho(T, X(T)) = 0$$

- \leadsto Set initially $\varrho = \varrho(T=0,X)$ to be the density at position $x=X/\varepsilon$. (Changes on the large scale.)
- $\sim \varrho(T,X)$ is the density of particles after a long time $t=T/\varepsilon$ at position $x=X/\varepsilon$. It satisfies, with a:=p-q,

$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1 - \varrho)] = 0 \quad \text{(inviscid Burgers)}$$

$$\frac{\partial}{\partial T} \varrho + a[1 - 2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$

$$\frac{\partial}{\partial T} \varrho + \frac{dX(T)}{dT} \cdot \frac{\partial}{\partial X} \varrho = \frac{d}{dT} \varrho(T, X(T)) = 0$$

- \leadsto Set initially $\varrho = \varrho(T=0,X)$ to be the density at position $x=X/\varepsilon$. (Changes on the large scale.)
- $\sim \varrho(T,X)$ is the density of particles after a long time $t=T/\varepsilon$ at position $x=X/\varepsilon$. It satisfies, with a:=p-q,

$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1 - \varrho)] = 0 \quad \text{(inviscid Burgers)}$$

$$\frac{\partial}{\partial T} \varrho + a[1 - 2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$

$$\frac{\partial}{\partial T} \varrho + \frac{dX(T)}{dT} \cdot \frac{\partial}{\partial X} \varrho = \frac{d}{dT} \varrho(T, X(T)) = 0$$

Let T and X be some large-scale time and space parameters.

- \leadsto Set initially $\varrho = \varrho(T=0,X)$ to be the density at position $x=X/\varepsilon$. (Changes on the large scale.)
- $\sim \varrho(T,X)$ is the density of particles after a long time $t=T/\varepsilon$ at position $x=X/\varepsilon$. It satisfies, with a:=p-q,

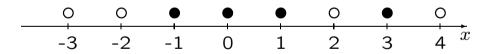
$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1 - \varrho)] = 0 \quad \text{(inviscid Burgers)}$$

$$\frac{\partial}{\partial T} \varrho + a[1 - 2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$

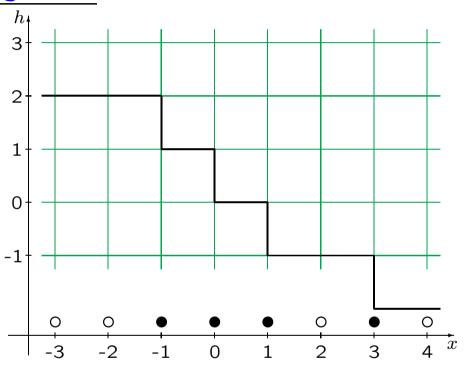
$$\frac{\partial}{\partial T} \varrho + \frac{dX(T)}{dT} \cdot \frac{\partial}{\partial X} \varrho = \frac{d}{dT} \varrho(T, X(T)) = 0$$

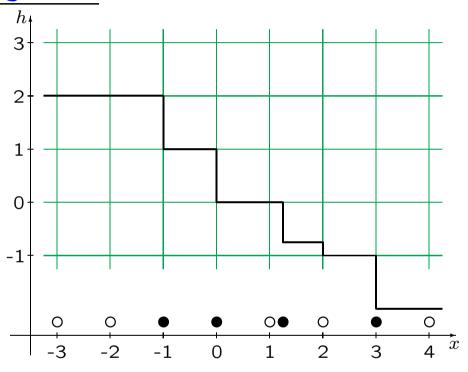
 \leadsto The characteristic speed $C(\varrho) := a[1 - 2\varrho]$. (ϱ is constant along $\dot{X}(T) = C(\varrho)$.)

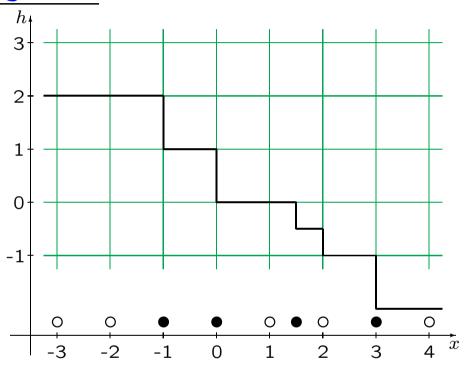
2. ASEP: Surface growth

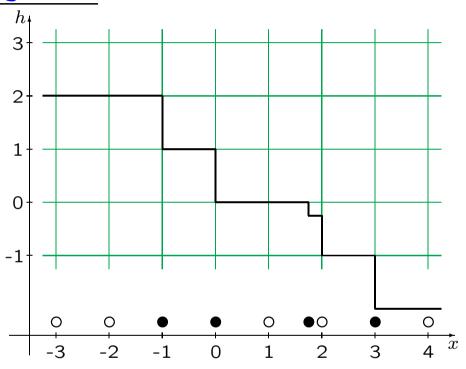


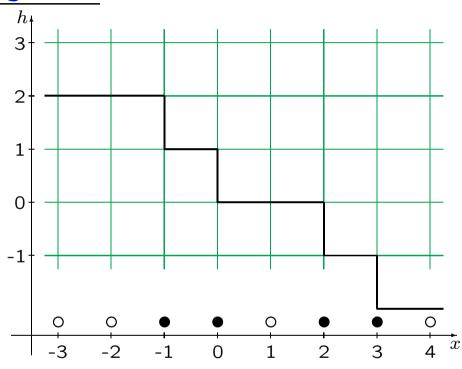
2. ASEP: Surface growth

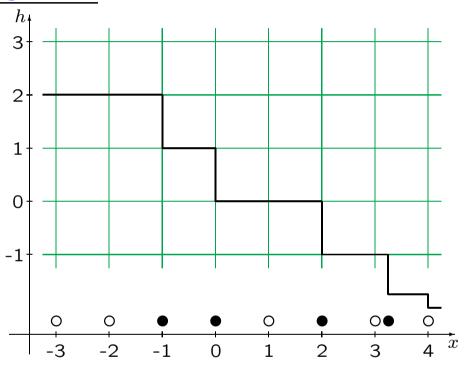


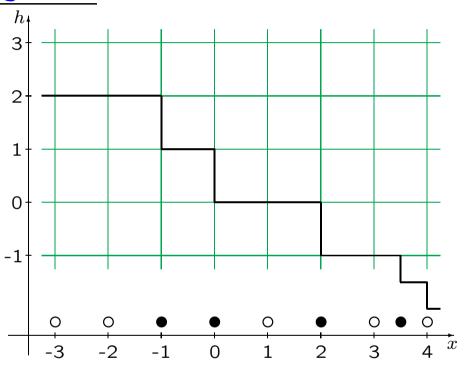


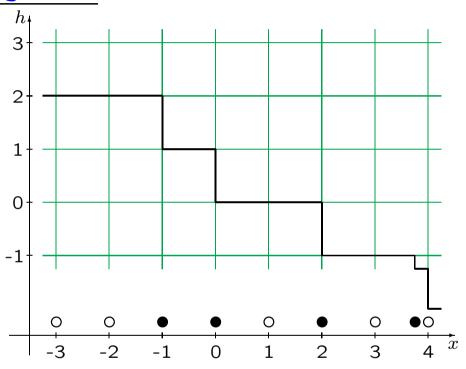


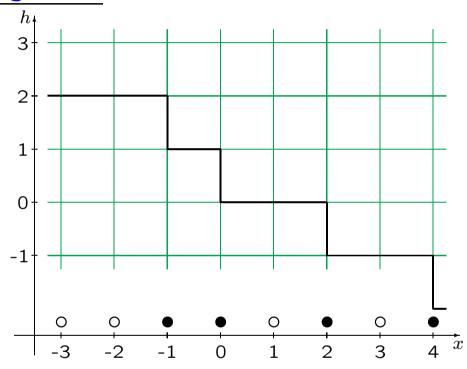


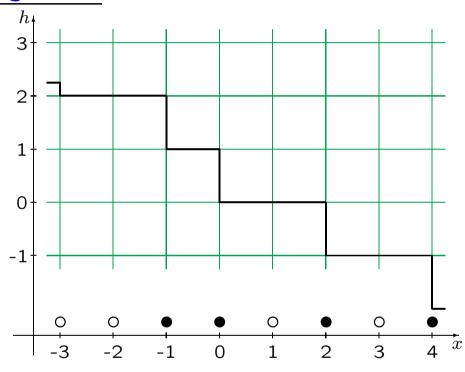


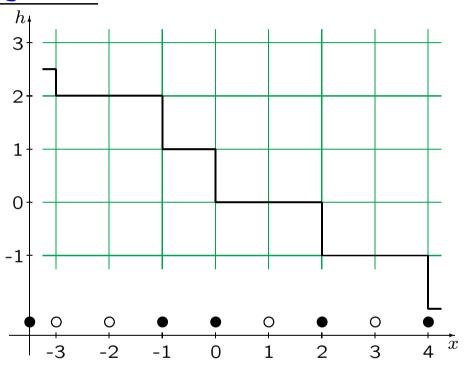


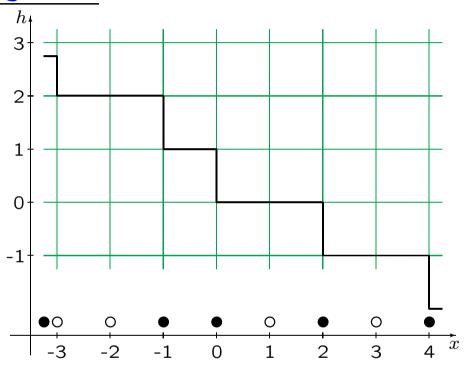


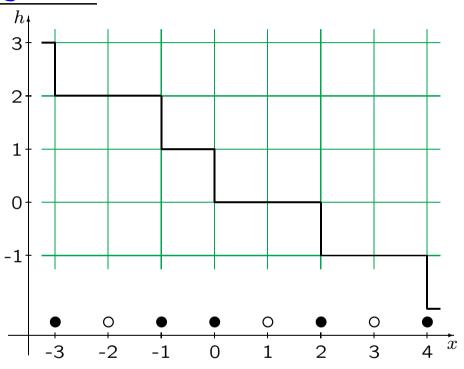


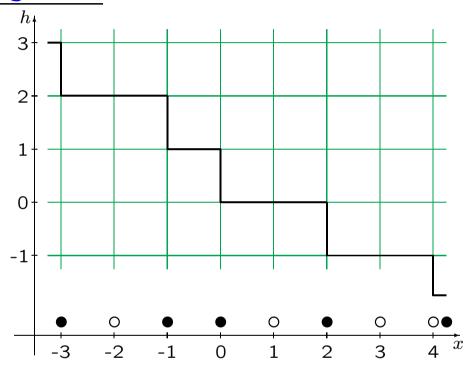


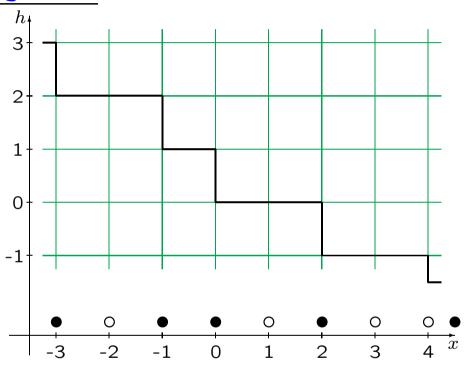


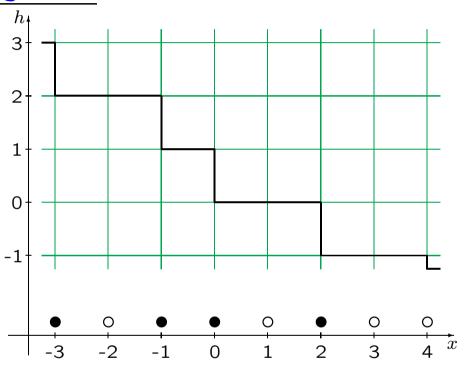


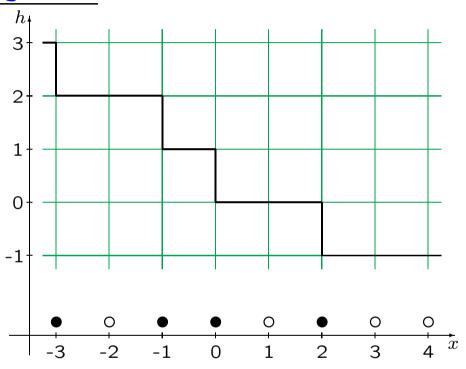


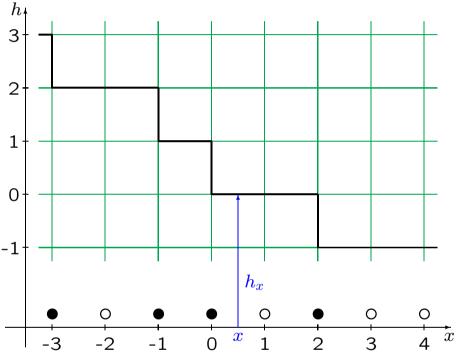




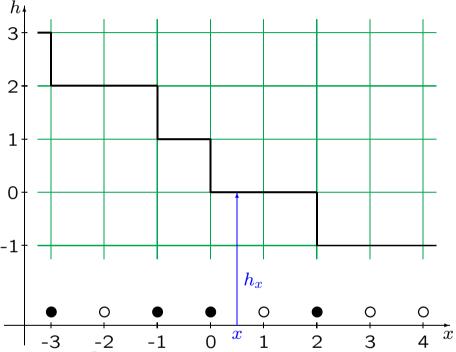




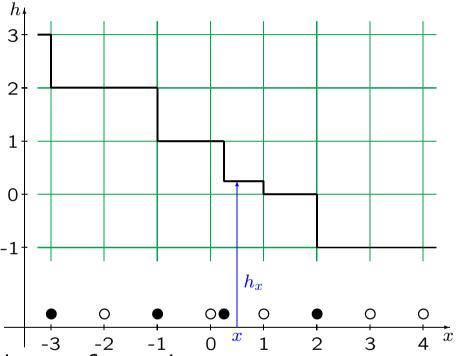




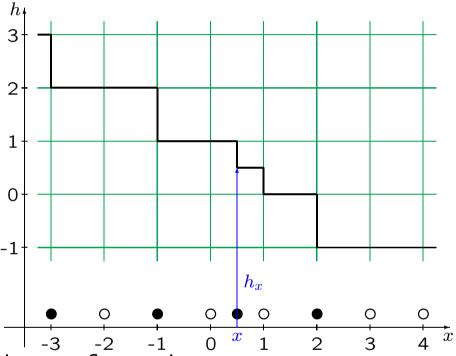
 $h_x(t) = \text{height of the surface above } x.$



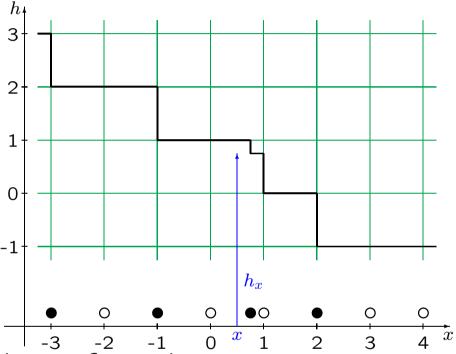
 $h_x(t)$ = height of the surface above x.



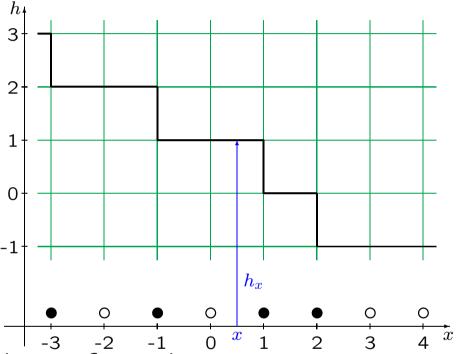
 $h_x(t)$ = height of the surface above x.



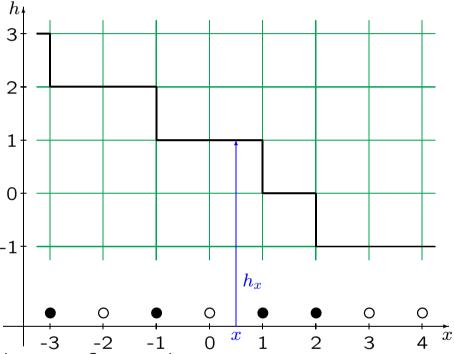
 $h_x(t)$ = height of the surface above x.



 $h_x(t)$ = height of the surface above x.



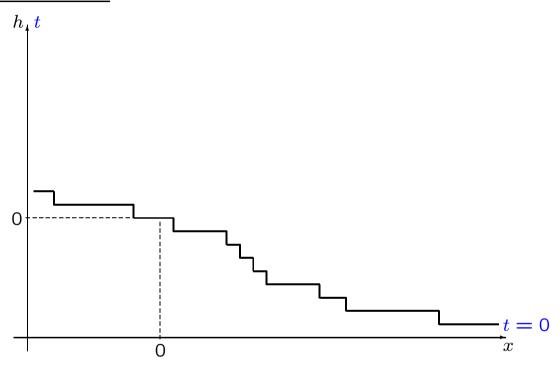
 $h_x(t)$ = height of the surface above x.

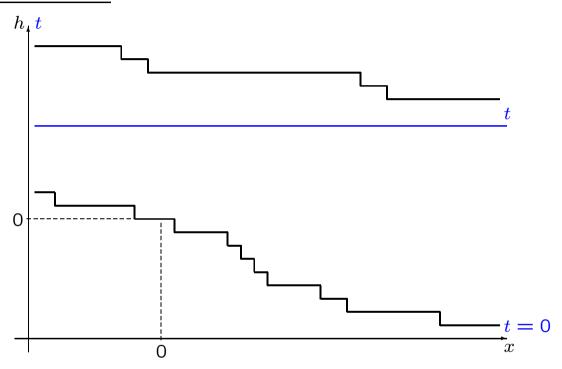


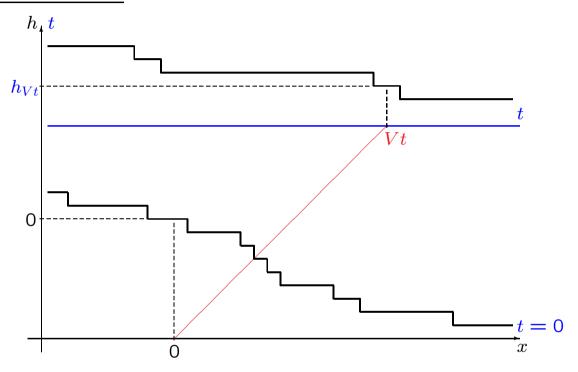
 $h_x(t)$ = height of the surface above x.

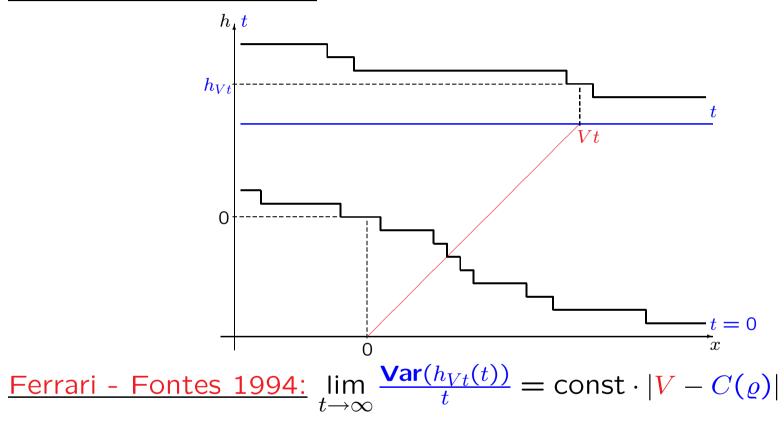
 $h_x(t) - h_x(0)$ = net number of particles passed above x.

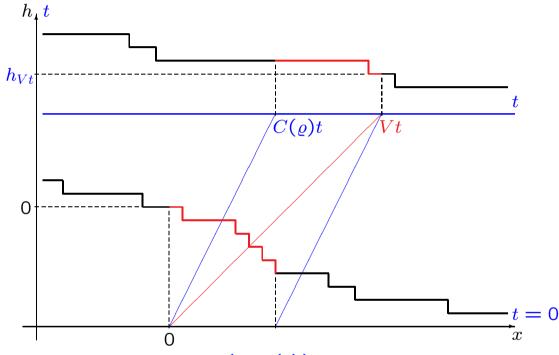
 $h_{Vt}(t)$ = net number of particles passed through the moving window at Vt $(V \in \mathbb{R})$.





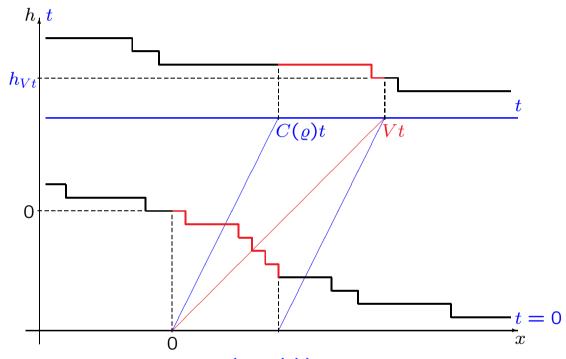






Ferrari - Fontes 1994: $\lim_{t \to \infty} \frac{\operatorname{Var}(h_{Vt}(t))}{t} = \operatorname{const} \cdot |V - C(\varrho)|$

→ Initial fluctuations are transported along the characteristics.



Ferrari - Fontes 1994: $\lim_{t \to \infty} \frac{\operatorname{Var}(h_{Vt}(t))}{t} = \operatorname{const} \cdot |V - C(\varrho)|$

- → Initial fluctuations are transported along the characteristics.
- \rightsquigarrow How about $V = C(\varrho)$?

Conjecture:

$$\lim_{t\to\infty}\frac{\mathrm{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}=\text{[sg. non trivial]}.$$

Conjecture:

$$\lim_{t\to\infty}\frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}=[\operatorname{sg. non trivial}].$$

Theorem (B., Seppäläinen): For any $0 < \varrho < 1$, and any q < p,

$$\begin{aligned} &0 < \liminf_{t \to \infty} \frac{ \frac{ \operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} \\ &\leq \limsup_{t \to \infty} \frac{ \operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} < \infty. \end{aligned}$$

Conjecture:

$$\lim_{t\to\infty}\frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}=[\operatorname{sg. non trivial}].$$

Theorem (B., Seppäläinen): For any $0 < \varrho < 1$, and any q < p,

$$\begin{aligned} &0 < \liminf_{t \to \infty} \frac{ \operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} \\ &\leq \limsup_{t \to \infty} \frac{ \operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} < \infty. \end{aligned}$$

Corollary: The corresponding scaling of the diffusivity is also proved.

Limit distributions (not yet controlling the second moment) in terms of the Tracy-Widom distribution (GUE random matrices) were found by Baik, Deift and Johansson 1999, Johansson 2000, and Ferrari and Spohn 2006 for the *totally* asymmetric exclusion (TASEP: p = 1, q = 0).

Limit distributions (not yet controlling the second moment) in terms of the Tracy-Widom distribution (GUE random matrices) were found by Baik, Deift and Johansson 1999, Johansson 2000, and Ferrari and Spohn 2006 for the *totally* asymmetric exclusion (TASEP: p = 1, q = 0).

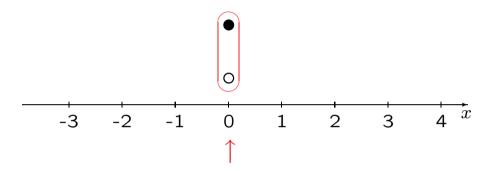
Method was: Last passage percolation, heavy combinatorics and asymptotic analysis.

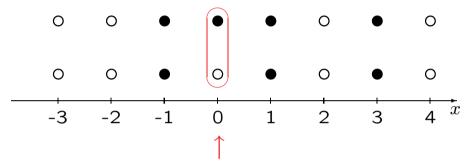
Limit distributions (not yet controlling the second moment) in terms of the Tracy-Widom distribution (GUE random matrices) were found by Baik, Deift and Johansson 1999, Johansson 2000, and Ferrari and Spohn 2006 for the *totally* asymmetric exclusion (TASEP: p = 1, q = 0).

Method was: Last passage percolation, heavy combinatorics and asymptotic analysis.

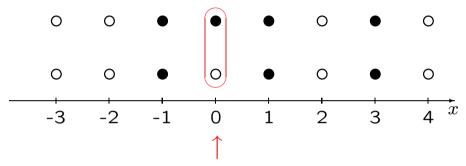
We needed to get rid of these tools. Premises: Cator and Groene-boom 2006 (Hammersley's process), B., Cator and Seppäläinen 2006 (TASEP, last passage).

4. The second class particle

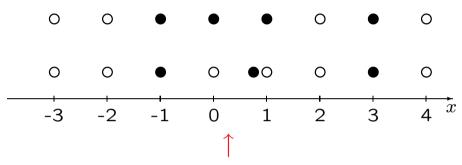




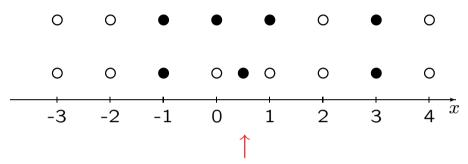
Bernoulli(ϱ) distribution except for 0



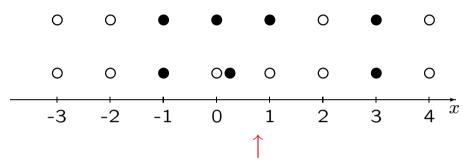
Bernoulli(ϱ) distribution except for 0



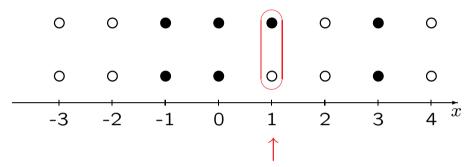
Bernoulli(ϱ) distribution except for 0



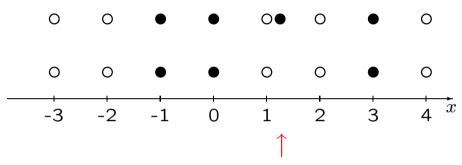
Bernoulli(ϱ) distribution except for 0



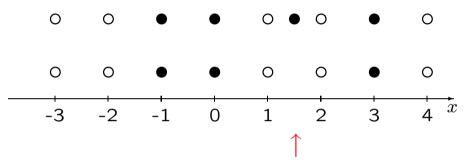
Bernoulli(ϱ) distribution except for 0



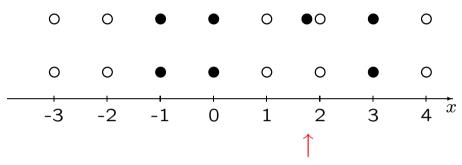
Bernoulli(ϱ) distribution except for 0



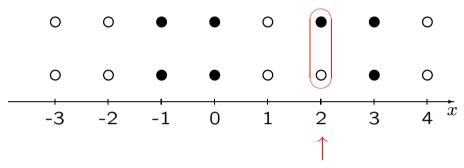
Bernoulli(ϱ) distribution except for 0



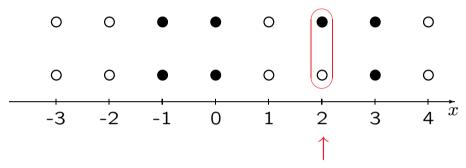
Bernoulli(ϱ) distribution except for 0



Bernoulli(ϱ) distribution except for 0

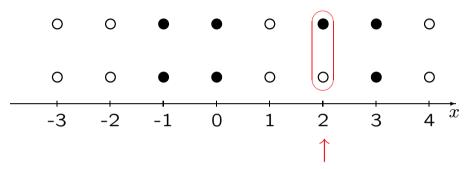


Bernoulli(ϱ) distribution except for 0



Bernoulli(o) distribution except for 0

Coupling: A single discrepancy is always conserved = the second class particle. Its location at time t is Q(t).



Bernoulli(ϱ) distribution except for 0

Coupling: A single discrepancy is always conserved = the second class particle. Its location at time t is Q(t).

The second class particle is a highly nontrivial object. For example, the Bernoulli(ϱ) distribution is *not* stationary as seen by the second class particle.

Theorem:

$$E(Q(t)) = C(\varrho)t$$

(characteristic speed),

Theorem:

$$E(Q(t)) = C(\varrho)t$$

(characteristic speed), and

$$Var(h_{Vt}(t)) = const \cdot E|Vt - Q(t)|.$$

Theorem:

$$E(Q(t)) = C(\varrho)t$$

(characteristic speed), and

$$Var(h_{Vt}(t)) = const \cdot E|Vt - Q(t)|.$$

Method of proof: martingale arguments, time-reversal, and conservation of particles.

Theorem:

$$E(Q(t)) = C(\varrho)t$$

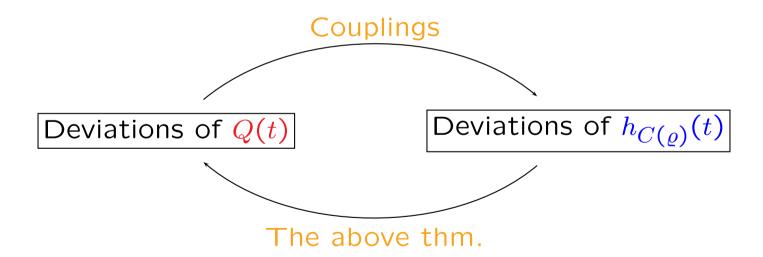
(characteristic speed), and

$$Var(h_{Vt}(t)) = const \cdot E|Vt - Q(t)|.$$

Method of proof: martingale arguments, time-reversal, and conservation of particles.

The proof is based on ideas of Bálint Tóth, he said these ideas were standard.

Main idea for prooving $t^{1/3}$ scaling:



The coupling measure

Let $\lambda < \varrho$, and

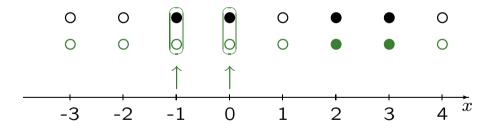
$$\mu({}^{\circ}_{\circ}) = 1 - {}_{\varrho}, \quad \mu({}^{\bullet}_{\circ}) = {}_{\varrho} - \lambda, \quad \mu({}^{\bullet}_{\bullet}) = \lambda.$$

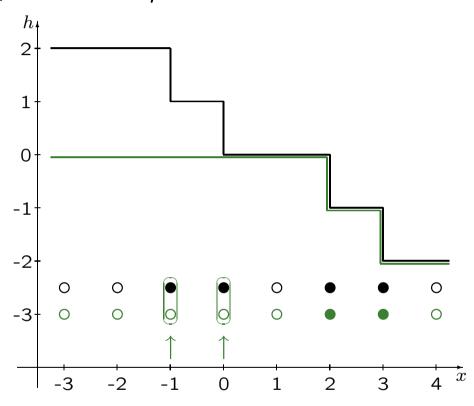
The coupling measure

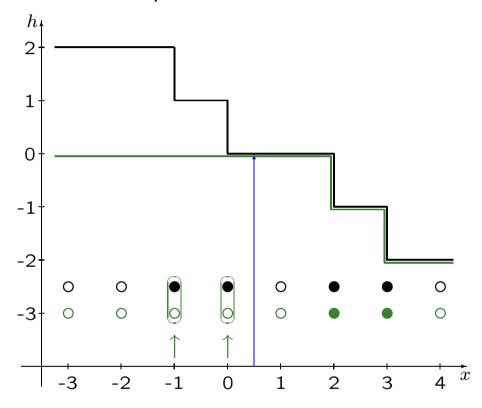
Let $\lambda < \varrho$, and

$$\mu({}^{\circ}_{\circ}) = 1 - {}_{\varrho}, \quad \mu({}^{\bullet}_{\circ}) = {}_{\varrho} - \lambda, \quad \mu({}^{\bullet}_{\bullet}) = \lambda.$$

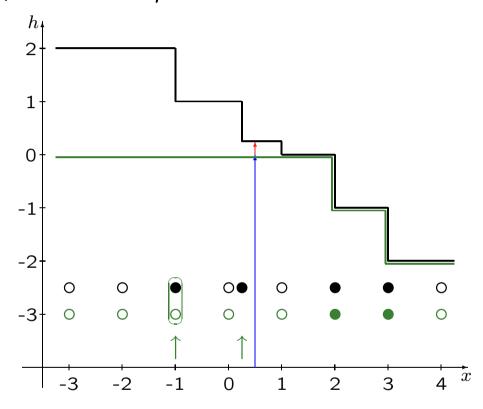
Then the "upper" marginal is Bernoulli(ϱ), and the "lower" marginal is Bernoulli(λ).



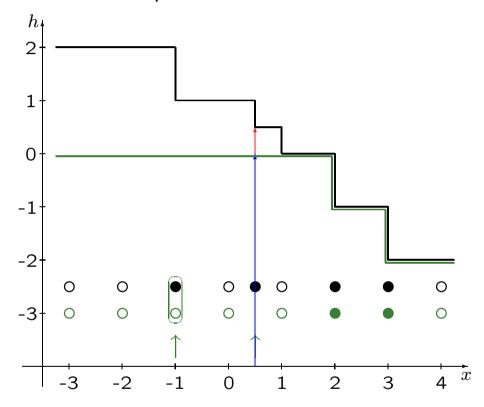




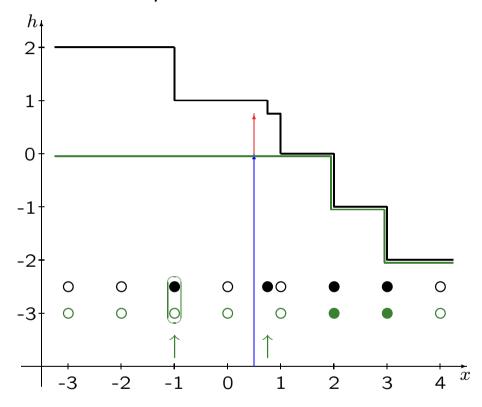
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



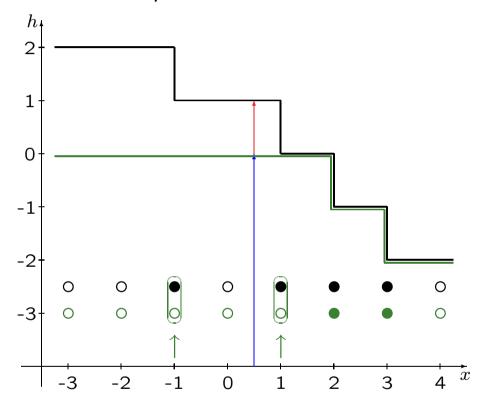
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



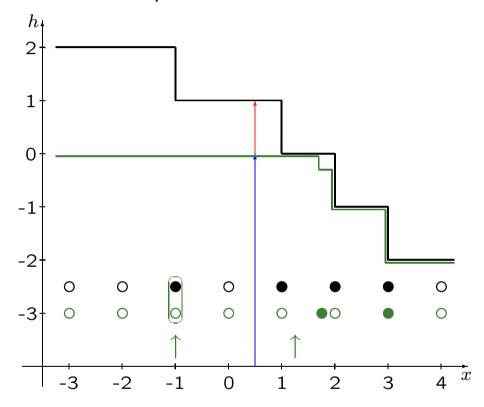
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



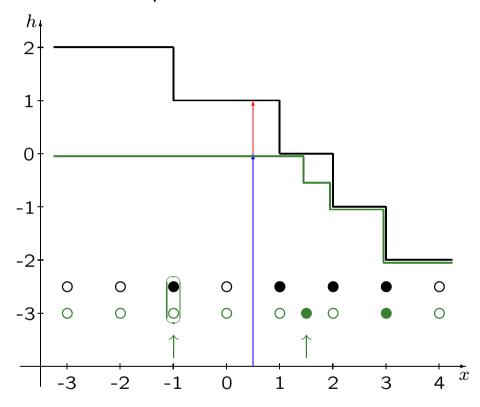
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



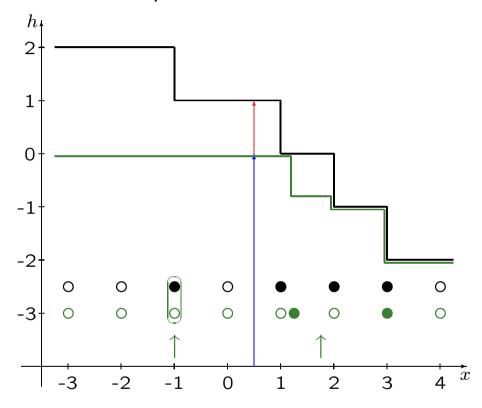
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



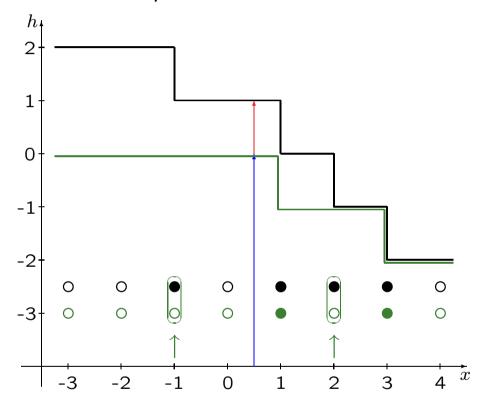
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



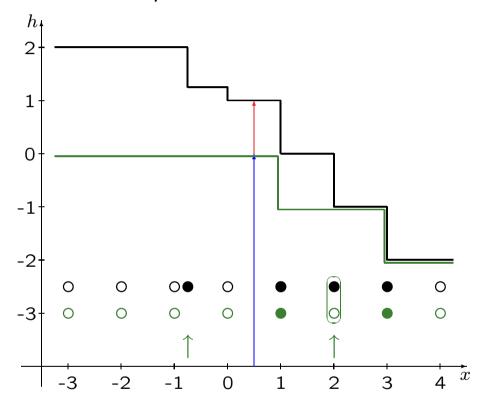
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



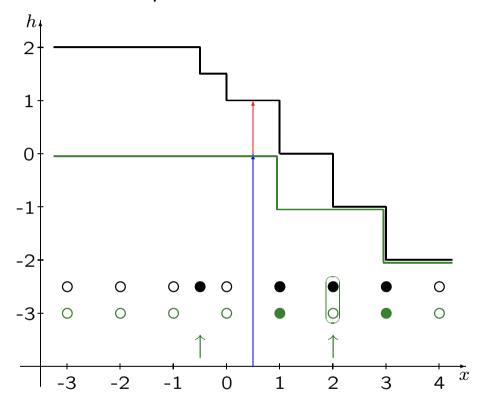
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



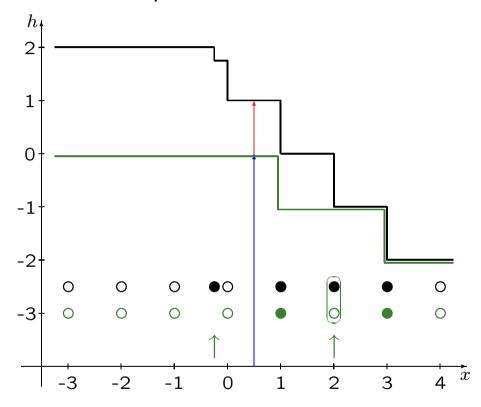
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



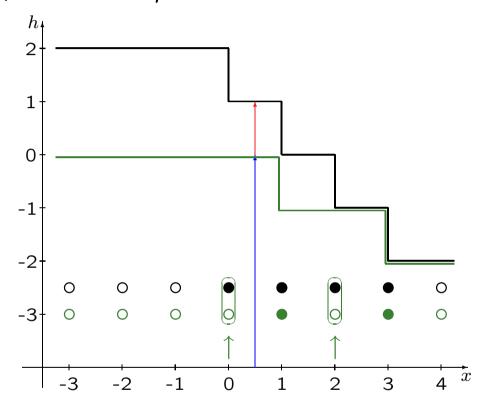
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



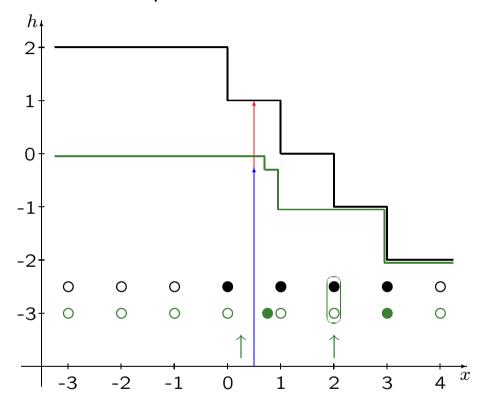
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



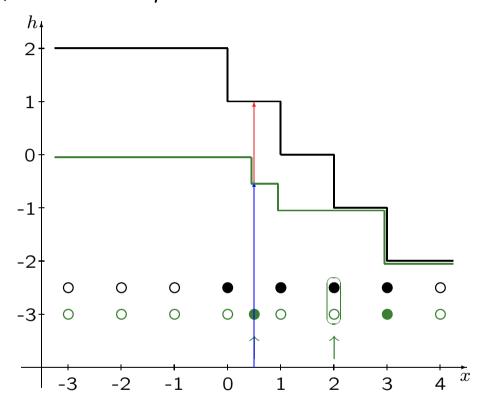
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



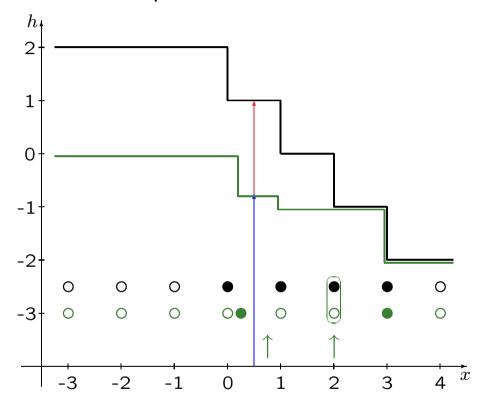
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



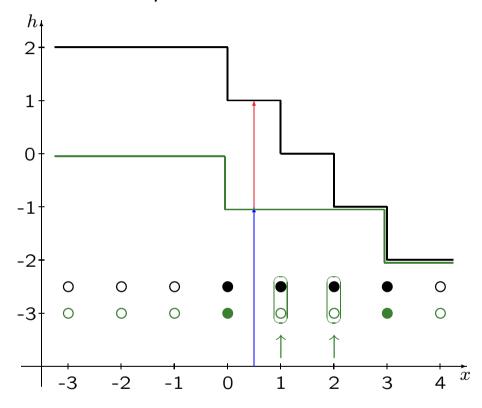
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



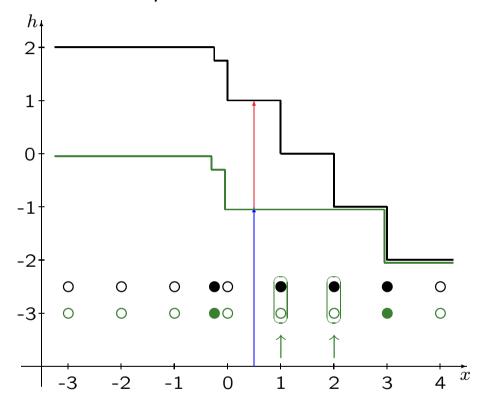
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



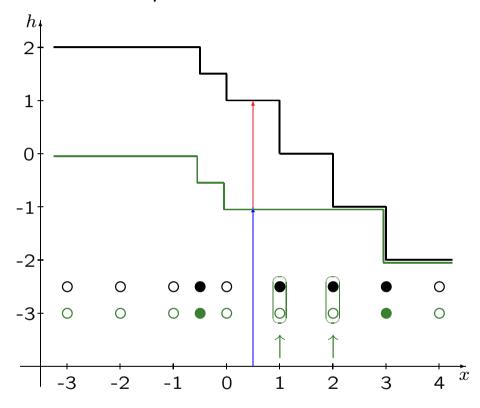
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



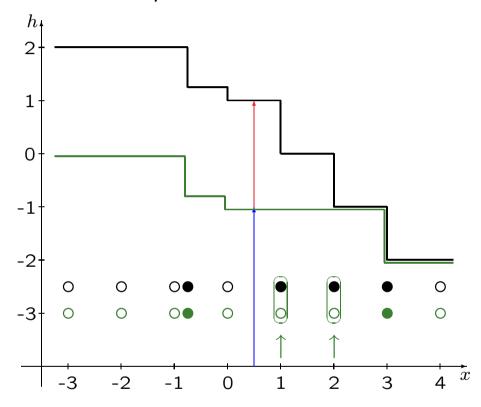
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



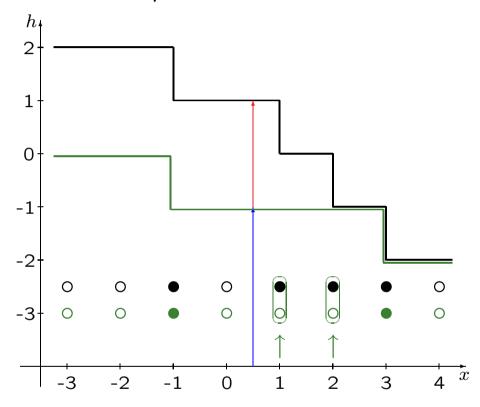
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.



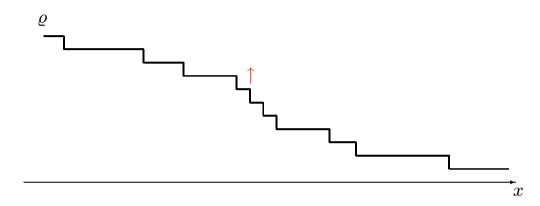
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.

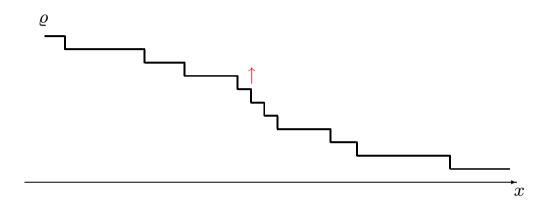


 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.

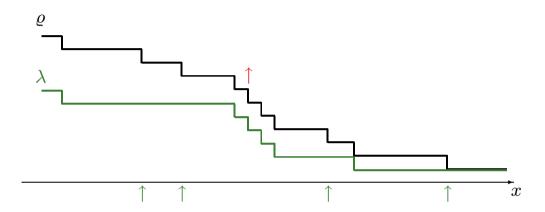


 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt $(V \in \mathbb{R})$.

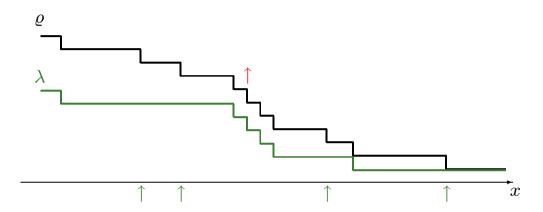




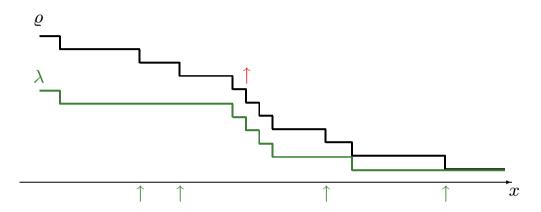
Connect Q(t)



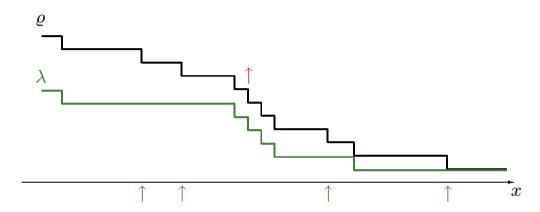
Connect Q(t) with the \uparrow 's



Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

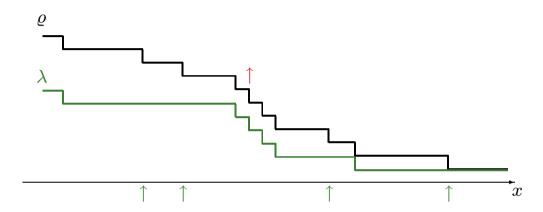


Connect Q(t) with the \uparrow 's (this needs nontrivial couplings): $\mathbf{P}\{Q(t) \text{ is too large}\}$



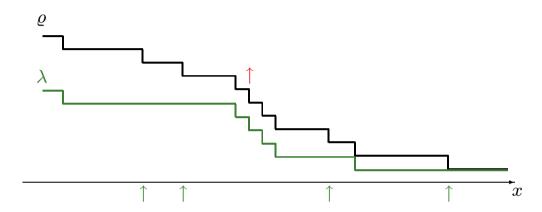
Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

 $\mathbf{P}\{Q(t) \text{ is too large}\} \leq \mathbf{P}\{\text{too many }\uparrow\text{'s have crossed } C(\varrho)t\}$



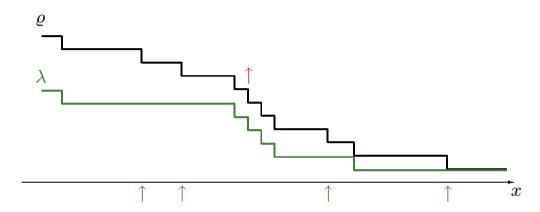
Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

$$\begin{split} \mathbf{P}\{Q(t) \text{ is too large}\} &\leq \mathbf{P}\{\text{too many } \uparrow \text{'s have crossed } C(\varrho)t\} \\ &\leq \mathbf{P}\{h_{C(\varrho)t}(t) - h_{C(\varrho)t}(t) \text{ is too large}\}. \end{split}$$



Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

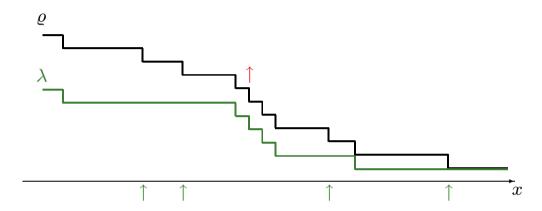
$$\begin{split} \mathbf{P}\{Q(t) \text{ is too large}\} &\leq \mathbf{P}\{\text{too many }\uparrow\text{'s have crossed } C(\varrho)t\} \\ &\leq \mathbf{P}\{h_{C(\varrho)t}(t) - h_{C(\varrho)t}(t) \text{ is too large}(\lambda)\}. \end{split}$$



Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

$$\begin{split} \mathbf{P}\{Q(t) \text{ is too large}\} &\leq \mathbf{P}\{\text{too many }\uparrow\text{'s have crossed } C(\varrho)t\} \\ &\leq \mathbf{P}\{h_{C(\varrho)t}(t) - h_{C(\varrho)t}(t) \text{ is too large}(\lambda)\}. \end{split}$$

Optimize "too large(λ)" in λ ,



Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

$$\begin{split} \mathbf{P}\{Q(t) \text{ is too large}\} &\leq \mathbf{P}\{\text{too many }\uparrow\text{'s have crossed } C(\varrho)t\} \\ &\leq \mathbf{P}\{h_{C(\varrho)t}(t) - h_{C(\varrho)t}(t) \text{ is too large}(\lambda)\}. \end{split}$$

Optimize "too large(λ)" in λ , use Chebyshev's inequality and relate $Var(h_{C(\rho)t}(t))$ to $Var(h_{C(\rho)t}(t))$.

The computations result in

$$\mathbf{P}\{Q(t) - C(\varrho)t \ge u\} \le c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{C(\varrho)t}(t))$$

$$\mathbf{P}\{Q(t) - C(\varrho)t \ge u\} \le c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{C(\varrho)t}(t))$$

$$\mathbf{P}\{Q(t) - C(\varrho)t \ge u\} \le c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{C(\varrho)t}(t))$$

$$\stackrel{\mathsf{Thm}}{=} c \cdot \frac{t^2}{u^4} \cdot \mathbf{E}|Q(t) - C(\varrho)t|.$$

$$\mathbf{P}\{\mathbf{Q}(t) - C(\varrho)t \ge u\} \le c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{C(\varrho)t}(t))$$

$$\stackrel{\mathsf{Thm}}{=} c \cdot \frac{t^2}{u^4} \cdot \mathbf{E}|\mathbf{Q}(t) - C(\varrho)t|.$$

With

$$\widetilde{Q}(t) := Q(t) - C(\varrho)t$$
 and $E := \mathbf{E}|\widetilde{Q}(t)|,$

we have (with a similar lower deviation bound)

$$\mathbf{P}\{|\widetilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E.$$

$$\mathbf{P}\{\mathbf{Q}(t) - C(\varrho)t \ge u\} \le c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{C(\varrho)t}(t))$$

$$\stackrel{\mathsf{Thm}}{=} c \cdot \frac{t^2}{u^4} \cdot \mathbf{E}|\mathbf{Q}(t) - C(\varrho)t|.$$

With

$$\widetilde{Q}(t) := Q(t) - C(\varrho)t$$
 and $E := \mathbf{E}|\widetilde{Q}(t)|,$

we have (with a similar lower deviation bound)

$$\mathbf{P}\{|\widetilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E.$$

Claim: this already implies the $t^{2/3}$ upper bound:

We had
$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.

We had
$$\mathbf{P}\{|\tilde{Q}(t)|>u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.
$$E = \mathbf{E}|\tilde{Q}(t)| = \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)|>u\} \ \mathrm{d}u$$

We had
$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.
$$E = \mathbf{E}|\tilde{Q}(t)| = \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > u\} \ \mathrm{d}u$$
$$= E \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \ \mathrm{d}v$$

We had
$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.
$$E = \mathbf{E}|\tilde{Q}(t)| = \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > u\} \ \mathrm{d}u$$
$$= E \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \ \mathrm{d}v$$
$$\le E \int_{1/2}^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \ \mathrm{d}v + \frac{1}{2}E$$

We had
$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.
$$E = \mathbf{E}|\tilde{Q}(t)| = \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > u\} \, \mathrm{d}u$$
$$= E \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \, \mathrm{d}v$$
$$\le E \int_{1/2}^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \, \mathrm{d}v + \frac{1}{2}E$$
$$\le c \cdot \frac{t^2}{E^2} + \frac{1}{2}E,$$

that is, $E^3 \leq c \cdot t^2$.

We had
$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.

$$E = \mathbf{E}|\tilde{Q}(t)| = \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > u\} \, du$$

$$= E \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \, dv$$

$$\leq E \int_{1/2}^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \, dv + \frac{1}{2}E$$

$$\leq c \cdot \frac{t^2}{E^2} + \frac{1}{2}E,$$

that is, $E^3 \leq c \cdot t^2$.

$$\operatorname{Var}(h_{C(\varrho)t}(t)) \stackrel{\mathsf{Thm}}{=} \operatorname{const.} \cdot \operatorname{E}|Q(t) - C(\varrho)t|$$

We had
$$\mathbf{P}\{|\widetilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.

$$E = \mathbf{E}|\tilde{Q}(t)| = \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > u\} \, du$$

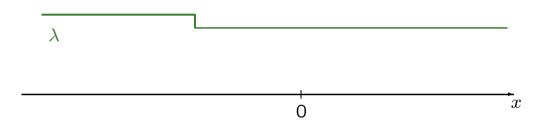
$$= E \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \, dv$$

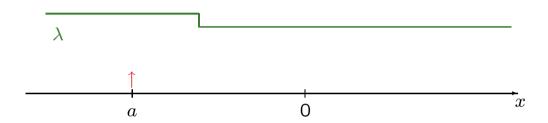
$$\leq E \int_{1/2}^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \, dv + \frac{1}{2}E$$

$$\leq c \cdot \frac{t^2}{E^2} + \frac{1}{2}E,$$

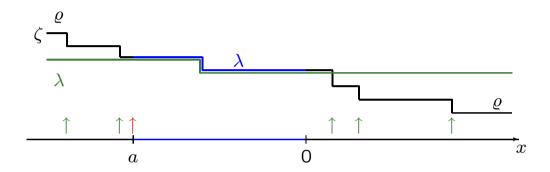
that is, $E^3 \leq c \cdot t^2$.

$$\mathbf{Var}(h_{C(\varrho)t}(t)) \stackrel{\mathsf{Thm}}{=} \operatorname{const.} \cdot \mathbf{E}|Q(t) - C(\varrho)t|$$
$$= \operatorname{const.} \cdot E \le c \cdot t^{2/3}.$$

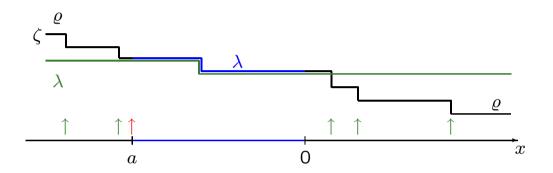




Let $Q^a(0) = a < 0$.

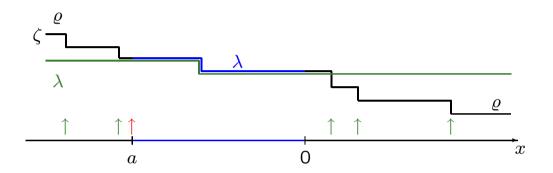


Let $Q^a(0) = a < 0$.



Let $Q^a(0) = a < 0$. If $Q^a(t) \le C(\varrho)t$, then the \uparrow 's have not crossed the path $C(\varrho)t$ from left to right:

$$\mathbf{P}\{Q^{a}(t) \leq C(\varrho)t\} \leq \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

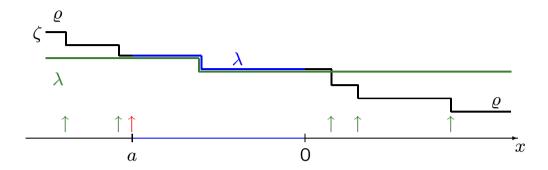


Let $Q^a(0) = a < 0$. If $Q^a(t) \le C(\varrho)t$, then the \uparrow 's have not crossed the path $C(\varrho)t$ from left to right:

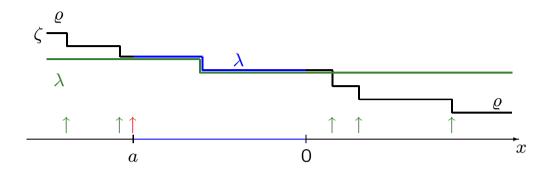
$$\mathbf{P}\{Q^{a}(t) \le C(\varrho)t\} \le \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

Therefore:

$$1 \le \mathbf{P}\{Q^a(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

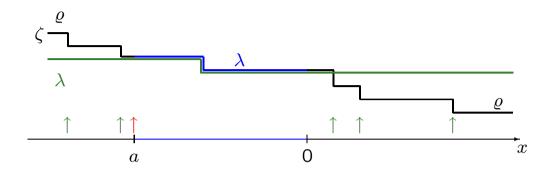


$$1 \le \mathbf{P}\{Q^a(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$$



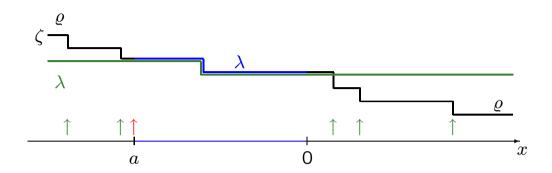
$$1 \le \mathbf{P}\{Q^a(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$$

 \rightsquigarrow Set a so that $\mathrm{E}(Q^a(t)) < C(\varrho)t$,



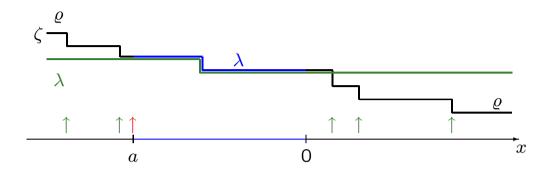
$$1 \le \mathbf{P}\{Q^a(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$$

- \leadsto Set a so that $\mathbf{E}(Q^a(t)) < C(\varrho)t$,
- $ightharpoonup \mathbf{E}(h_{C(\varrho)t}(t)) \mathbf{E}(h_{C(\varrho)t}(t)) \sim t(\varrho \lambda)^2 > 0$ would be the case, if ζ was Bernoulli(ϱ) distributed.



$$1 \le \mathbf{P}\{Q^a(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$$

- \leadsto Set a so that $\mathbf{E}(Q^a(t)) < C(\varrho)t$,
- $ightharpoonup \mathbf{E}(h_{C(\varrho)t}(t)) \mathbf{E}(h_{C(\varrho)t}(t)) \sim t(\varrho \lambda)^2 > 0$ would be the case, if ζ was Bernoulli(ϱ) distributed. Instead, $\mathbf{E}(h_{C(\varrho)t}(t))$ will have a harmless Radon-Nikodym factor.



$$1 \le \mathbf{P}\{Q^a(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$$

- \leadsto Set a so that $\mathbf{E}(Q^a(t)) < C(\varrho)t$,
- $ightharpoonup \mathbf{E}(h_{C(\varrho)t}(t)) \mathbf{E}(h_{C(\varrho)t}(t)) \sim t(\varrho-\lambda)^2 > 0$ would be the case, if ζ was Bernoulli(ϱ) distributed. Instead, $\mathbf{E}(h_{C(\varrho)t}(t))$ will have a harmless Radon-Nikodym factor.
- ⇒ Both probabilities are deviation probabilities.

$$1 \le \mathbf{P}\{Q^{a}(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

$$1 \le \mathbf{P}\{Q^a(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

Apply Markov's inequality on the first, Chebyshev's on the second probability (use again the connection between $Var(h_{C(\varrho)t}(t))$ and $Var(h_{C(\varrho)t}(t))$).

$$1 \le P\{Q^a(t) > C(\varrho)t\} + P\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

Apply Markov's inequality on the first, Chebyshev's on the second probability (use again the connection between $Var(h_{C(\rho)t}(t))$ and $Var(h_{C(\rho)t}(t))$).

The correct scaling of the parameters is: $\varrho - \lambda \sim t^{-1/3}, \quad a \sim -t^{2/3}.$

$$1 \le P\{Q^a(t) > C(\varrho)t\} + P\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

Apply Markov's inequality on the first, Chebyshev's on the second probability (use again the connection between $Var(h_{C(\varrho)t}(t))$ and $Var(h_{C(\varrho)t}(t))$).

The correct scaling of the parameters is: $\varrho - \lambda \sim t^{-1/3}, \quad a \sim -t^{2/3}.$ In this case

$$1 \leq c_1 \cdot \frac{\mathbf{E}(|\widetilde{Q^a}(t)|)}{t^{2/3}} + c_2 \cdot \frac{\mathbf{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}$$

$$1 \le P\{Q^a(t) > C(\varrho)t\} + P\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

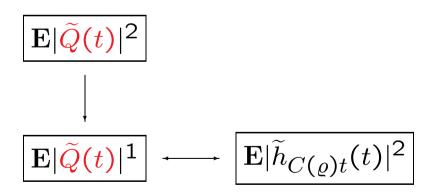
Apply Markov's inequality on the first, Chebyshev's on the second probability (use again the connection between $Var(h_{C(\varrho)t}(t))$ and $Var(h_{C(\varrho)t}(t))$).

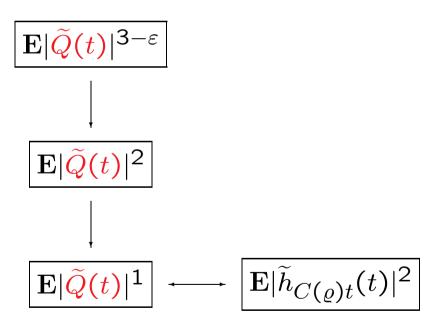
The correct scaling of the parameters is: $\varrho - \lambda \sim t^{-1/3}, \quad a \sim -t^{2/3}.$ In this case

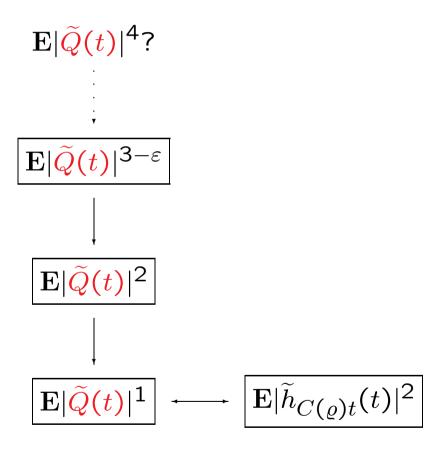
$$1 \leq c_1 \cdot \frac{\mathrm{E}(|Q^a(t)|)}{t^{2/3}} + c_2 \cdot \frac{\mathrm{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}$$

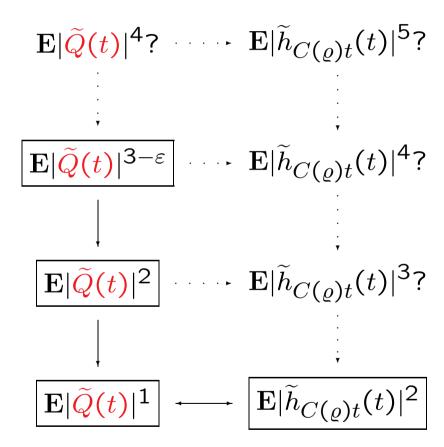
$$\stackrel{\mathsf{Thm}}{=} c \cdot \frac{\mathrm{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}.$$

$$\mathbf{E}|\widetilde{Q}(t)|^{1} \longrightarrow \mathbf{E}|\widetilde{h}_{C(\varrho)t}(t)|^{2}$$









 \rightarrow What is the limit $\lim_{t\rightarrow\infty}\frac{\mathbf{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}=$? What does it have to do with Gaussian random matrices? (Difficult.)

- ightharpoonup What is the limit $\lim_{t \to \infty} \frac{ {
 m Var}(h_{C(\varrho)t}(t))}{t^{2/3}} = ?$ What does it have to do with Gaussian random matrices? (Difficult.)
- → Other processes (zero range, *Bricklayers*', ...)?

- \rightarrow What is the limit $\lim_{t\to\infty} \frac{{\sf Var}(h_{C(\varrho)t}(t))}{t^{2/3}}=$? What does it have to do with Gaussian random matrices? (Difficult.)
- → Other processes (zero range, *Bricklayers'*, ...)?
- \rightarrow Some processes (e.g. symmetric simple exclusion, linear rate zero range) show $t^{1/4}$ scaling (with Gaussian limits), rather than $t^{1/3}$. Where is the borderline? Are there other scalings as well?

Thank you.