

# The Growth Model: Busemann Functions, Shape, Geodesics, and Other Stories

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May 1, 2015

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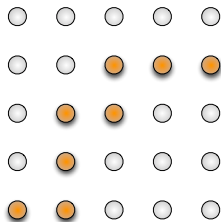
Lightning captured at 7,207 images per second  
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# Last Passage Percolation

**Random potential**  $\omega = (\omega_x)_{x \in \mathbb{Z}^2} \in \Omega$ ,  $\mathbb{R}$ -valued i.i.d.,  $2 + \varepsilon$  moments.

**Up-right paths**  $x_{0,n} = (x_0, \dots, x_n)$   
take steps  $e_1 = (1, 0)$  or  $e_2 = (0, 1)$ .

Passage time of path  $x_{0,n}$  is  $\sum_{i=0}^{n-1} \omega_{x_i}$ .

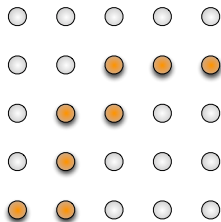


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**Point-to-point last passage time:**  $G_{x,y}(\omega) = \max_{x_0=x, x_n=y} \sum_{k=0}^{n-1} \omega_{x_k}$ .

Connections to: Totally Asymmetric Simple Exclusion, Queuing Theory, Corner Growth Model, etc.

# Shape Theorem

LLN says sum of i.i.d. grows linearly.

$G_{0,x}$  is not quite a sum of i.i.d.

It is however superadditive:  $G_{x,y} + G_{y,z} \leq G_{x,z}$ .

Then: outside one null set, for all  $\xi \in \mathbb{R}_+^2$  and all  $x_n/n \rightarrow \xi$  simultaneously

$g_{pp}(\xi) = \lim_{n \rightarrow \infty} n^{-1} G_{0,x_n}$  exists, is deterministic, concave, homogenous  
( $g_{pp}(c\xi) = cg_{pp}(\xi)$ ) and continuous all the way to the boundary.

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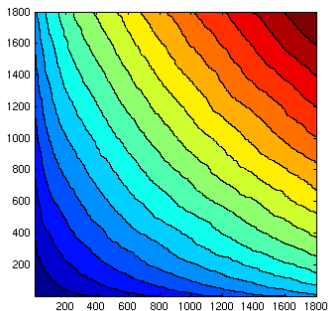
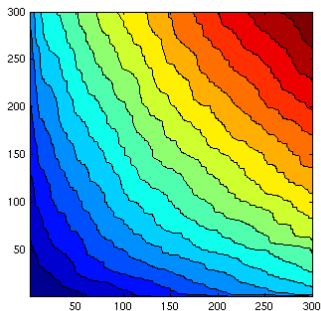
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$$\{x : G_{0,x} \leq t\}$$

Path  $x_{0,n}$  that maximizes  $G_{x,y}$  is called a **geodesic**.

$x_{0,\infty}$  is a  $\xi$ -geodesic if  $\forall n$   $x_{0,n}$  is a geodesic and  $x_n/n \rightarrow \xi$ .

Given  $\xi \in \mathbb{R}_+^2$ , is there an infinite  **$\xi$ -geodesic**?

Is it the limit of the geodesic from 0 to  $x_n$  as  $n \rightarrow \infty$  and  $x_n/n \rightarrow \xi$ ?

If  $\omega_0$  is continuous then finite geodesics are unique.

Is the infinite  $\xi$ -geodesic unique?

Do  $\xi$ -geodesics out of  $x$  and  $y$  coalesce (i.e. eventually merge)?

Licea and Newman '96: answers are in the positive for standard first passage percolation (nearest-neighbor paths minimizing the passage time) if  $g_{pp}(\xi)$  satisfies a global curvature assumption.

Problem: the curvature assumption has not been proved. Though conjectured to hold.

Damron and Hanson '14: Existence holds under just strict convexity or differentiability of  $g_{pp}$  (which presumably should be “easier” to prove).

Ferrari and Pimentel '05: answers are in the positive also for the last passage percolation model we are considering, but with  $\omega_0$  exponential.

The exponential model is one of the **solvable** models for which explicit computations are possible. In particular, an explicit formula is available for the shape  $g_{pp}(\xi)$ .

Would like to allow more general weight distributions.



# Understanding the shape

Consider a finite subset  $V \subset \mathbb{Z}^2$  containing 0 (e.g.  $\{u : |u| \leq L\}$ ).

$\{G_{0,z_n-u} - G_{0,z_n} : u \in V\}$  describes the **microscopic** shape around  $z_n$ .

Expect this random vector to converge in distribution as  $z_n/n \rightarrow \xi$ .

Shifting by  $-z_n$  and reflecting  $\omega_x \mapsto \omega_{-x}$  turns the above into

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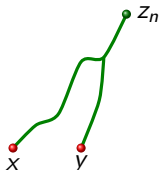
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**Busemann functions:**  $B^\xi(x, y; \omega) = \lim_{z_n/n \rightarrow \xi} (G_{x, z_n} - G_{y, z_n})$ .

Limit exists if e.g. geodesics coalesce.



# Understanding geodesics

Note that  $G_{x,z_n} = \omega_x + \max(G_{x+e_1,z_n}, G_{x+e_2,z_n})$ .

So  $(G_{x,z_n} - G_{x+e_1,z_n}) \wedge (G_{x,z_n} - G_{x+e_2,z_n}) = \omega_x$  almost surely.

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$n \rightarrow \infty$  gives  $B^\xi(x, x + e_1) \wedge B^\xi(x, x + e_2) = \omega_x$  almost surely.

The above suggests that  $\xi$ -geodesic out of  $x$  should follow the smallest  $B^\xi(x, x + z)$ ,  $z \in \{e_1, e_2\}$ .

## Busemann functions

Consider  $g_{pp}$  as a concave function on  $\mathcal{U} = \{(t, 1 - t) : t \in (0, 1)\}$ .

Given  $\xi \in \mathcal{U}$  let  $[\underline{\xi}, \bar{\xi}] \subset \mathcal{U}$  be the maximal (possibly degenerate) interval containing  $\xi$  on which  $g_{pp}$  is linear.

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**Standing assumptions:**  $\mathbb{P}\{\omega_0 \geq c\} = 1$ ,  $\omega_x$  i.i.d. with  $2 + \varepsilon$  moments,  $\underline{\xi}$  and  $\bar{\xi}$  are points of differentiability.

**Theorem.**  $B^\xi(x, y; \omega) = \lim_{z_n/n \rightarrow \xi} (G_{x, z_n} - G_{y, z_n})$  exists **a.S.**

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**Remark.**  $\omega_0 \geq c$  only because we use results from queuing where service times were assumed nonnegative. All the queuing results seem to go through without this assumption.

# Properties

$L^1$ :  $\mathbb{E}[|B^\xi(x, y)|] < \infty$ . (Comes from construction.)

**Stationary:**  $B^\xi(x, y; T_z\omega) = B^\xi(x + z, y + z; \omega)$  ( $(T_z\omega)_x = \omega_{x+z}$ )

**Cocycle:**  $B^\xi(x, y) + B^\xi(y, z) = B^\xi(x, z)$ .

The space of  $L^1$  stationary cocycles:  $\mathcal{C}$ .

**Potential recovery:**  $B^\xi(0, e_1) \wedge B^\xi(0, e_2) = \omega_0$  almost surely.



# Geodesics

If  $B \in \mathcal{C}$  then a  $B$ -geodesic is a path that follows the minimal  $B(x, x + z)$ ,  $z \in \{e_1, e_2\}$ . (In case of ties, OK to go either way.)

**Theorem.** If  $B$  recovers potential  $\omega$  ( $B(0, e_1) \wedge B(0, e_2) = \omega_0$  a.s.) then a  $B$ -geodesic is a geodesic: every finite piece of it is a geodesic.

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Given  $\xi \in \mathcal{U}$ , recall the maximal linear segment  $[\underline{\xi}, \bar{\xi}]$ .

A geodesic  $x_{0,\infty}$  is directed in  $[\underline{\xi}, \bar{\xi}]$  if all limit points of  $x_n/n$  belong to this interval.

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**Theorem.**

- Any  $B^\xi$ -geodesic is directed in  $[\underline{\xi}, \bar{\xi}]$ .
- Any geodesic directed in  $[\underline{\xi}, \bar{\xi}]$  is a  $B^\xi$ -geodesic.
- The  $B^\xi$ -geodesic with  $e_2$ -tie breaks is the topmost of all geodesics directed in  $[\underline{\xi}, \bar{\xi}]$ . Similarly for rightmost.

**Corollary.** If  $g_{pp}$  is differentiable everywhere, then every geodesic is directed in  $[\underline{\xi}, \bar{\xi}]$  for some  $\xi$ .

**Remark.** Can also handle corners, but will omit.

Thus can show: If  $g_{pp}$  is strictly concave, then every geodesic has a direction  $\xi$ , i.e.  $\lim x_n/n$  exists almost surely.

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Thus can show: If  $g_{pp}$  is strictly concave, then every geodesic has a direction  $\xi$ , i.e.  $\lim x_n/n$  exists almost surely.

**Theorem.** Assume also  $\mathbb{P}\{\omega_0 \leq r\}$  is continuous in  $r$ . Then  $\mathbb{P}\{B^\xi(0, e_1) = B^\xi(0, e_2)\} = 0$ .

**Corollary.** If  $\omega_0$  is continuous, then there exists a **unique** geodesic directed in  $[\underline{\xi}, \bar{\xi}]$  out of every point  $x \in \mathbb{Z}^2$ .

**Theorem.** Topmost  $[\underline{\xi}, \bar{\xi}]$ -directed geodesics coalesce, rightmost  $[\underline{\xi}, \bar{\xi}]$ -geodesics coalesce, and when  $\omega_0$  is continuous,  $[\underline{\xi}, \bar{\xi}]$ -geodesics coalesce.

# Variational formula

Until recently, the only description of  $g_{pp}(\xi)$  was from superadditivity:

$$g_{pp}(\xi) = \sup_n n^{-1} \mathbb{E}[G_{0,[n\xi]}] \text{ (e.g. if } \xi \in \mathbb{Z}_+^2 \text{)}.$$

Going through random polymer models:

**Theorem.**  $g_{pp}(\xi) = \inf_{B \in \mathcal{C}} \text{ess sup} \{ \omega_0 - B(0, e_1; \omega) \wedge B(0, e_2; \omega) + \bar{B} \cdot \xi \}.$

( $\bar{B} = (\mathbb{E}[B(0, e_1)], \mathbb{E}[B(0, e_2)])$  and  $\mathcal{C}$  is class of  $L^1$  stationary cocycles.)

Such formulas are important in statistical mechanics: their solutions are expected to describe the infinite-volume system (i.e. geodesics and shape as  $n \rightarrow \infty$ ).

**Theorem.** Under the standing assumptions,  $B^\xi$  solves the variational formula for  $g_{pp}(\xi)$ . In fact, the essential supremum is not needed and we have almost surely

$$g_{pp}(\xi) = \omega_0 - B^\xi(0, e_1, \omega) \wedge B^\xi(0, e_2, \omega) + \overline{B^\xi} \cdot \xi.$$

**Corollary.** Due to potential recovery, we have  $g_{pp}(\xi) = \overline{B^\xi} \cdot \xi$ .

Using some calculus one then gets that  $\overline{B^\xi} = \nabla g_{pp}(\xi)$ .

Nice interpretation: average **microscopic** gradient is **macroscopic** gradient.

## Solvable models

When  $\omega_0$  are exponential or geometric we in fact can calculate explicitly the distributions of  $B^\xi(0, e_1)$  and  $B^\xi(0, e_2)$  for all  $\xi \in \mathcal{U}$ .

For example, if  $\omega_0$  is exponential with rate 1, then  $B^\xi(0, e_1)$  is exponential with rate  $\alpha$  and  $B^\xi(0, e_2)$  is exponential with rate  $1 - \alpha$ , where

$$\alpha = \frac{\sqrt{\xi_1}}{\sqrt{\xi \cdot e_1} + \sqrt{\xi \cdot e_2}}.$$

Then  $\overline{B^\xi} = (\mathbb{E}[B^\xi(0, e_1)], \mathbb{E}[B^\xi(0, e_2)]) = (\frac{1}{\alpha}, \frac{1}{1-\alpha})$

and  $g_{pp}(\xi) = \overline{B^\xi} \cdot \xi = (\sqrt{\xi \cdot e_1} + \sqrt{\xi \cdot e_2})^2$ .

This is the known formula derived by Rost '81.



# Fluctuations

CLT says that if  $X_{0,n}$  has increments  $e_1$  or  $e_2$  equally likely, then it fluctuates from its average (straight line from 0 to  $(n/2, n/2)$ ) by  $n^{1/2}$ .

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Say  $\omega_0$  is continuous.

What are the fluctuations of the geodesic from 0 to  $[n\xi]$ ?

Conjecture: with enough moments on  $\omega_0$  geodesic has fluctuations of order  $n^{2/3}$ .

Superdiffusivity is because the path goes “out of its way” looking for high values of the potential.

On the other hand,  $G_{0,[n\xi]}$  should have  $n^{1/3}$  fluctuations.

Limit distributions related to **Tracy-Widom** from random matrices.

Models with these fluctuation exponents are said to belong to the **Kardar-Parisi-Zhang (KPZ)** universality class.

Johansson '00 proved LPP with exponential weights is in the KPZ class.

Again: solvability of the model was key.

# Fluctuations

When  $\omega_0$  is exponential or geometric,  $B^\xi(ne_1, (n+1)e_1)$  are i.i.d. and so are  $B^\xi(ne_2, (n+1)e_2)$ .

Balázs, Cator, and Seppäläinen '06 used this to prove the  $n^{2/3}$  fluctuations of the geodesic and  $n^{1/3}$  fluctuations of the last passage time, in the exponential weights case, with less technology than Johansson's proof of the Tracy-Widom limit.

More generally, CLT exponents for fluctuations of  $B^\xi(0, ne_1)$  and  $B^\xi(0, ne_2)$  imply information about fluctuation exponents of last passage quantities. (The above BCS result is one way to achieve this.)

Now we have a promising route to proving universality of KPZ fluctuations for general weight distributions.

# Thank You

