Order of current variance in the simple exclusion process

Márton Balázs

(University of Wisconsin - Madison) (Budapest University of Technology and Economics)

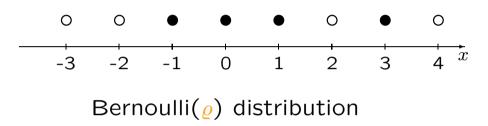
Joint work with

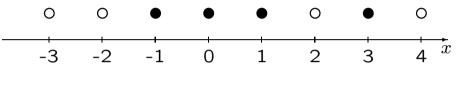
Timo Seppäläinen (University of Wisconsin - Madison)

Toronto, March 12, 2007.

- 1. ASEP: Interacting particles
- 2. ASEP: Surface growth
 - 3. Growth fluctuations
 - 4. The second class particle
 - 5. The upper bound
 - 6. The lower bound
 - 7. Open questions



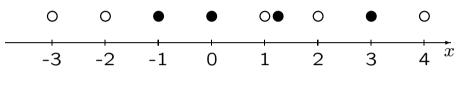




 $Bernoulli(\varrho)$ distribution

Particles try to jump

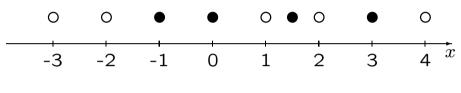
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

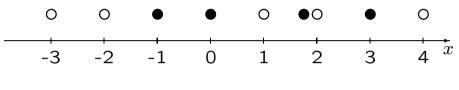
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

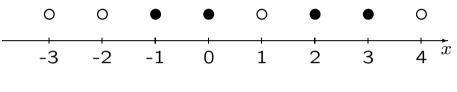
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

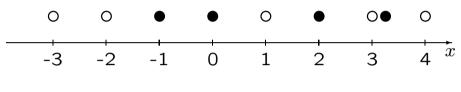
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

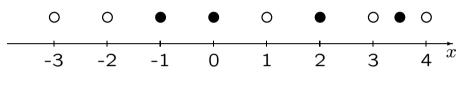
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

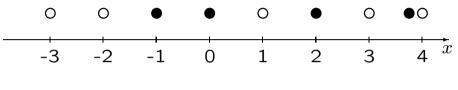
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

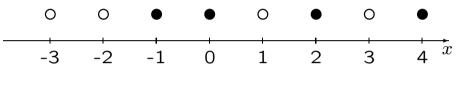
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

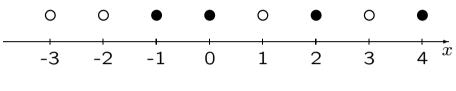
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

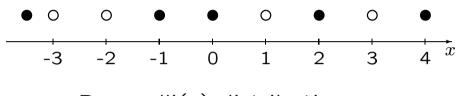
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

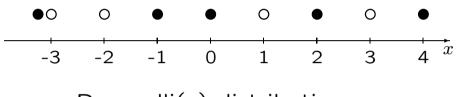
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

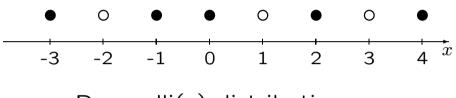
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

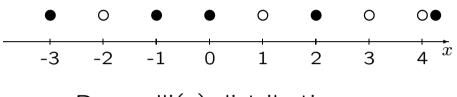
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

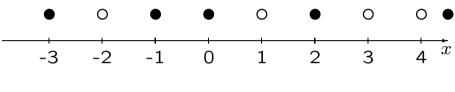
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

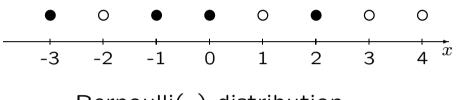
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

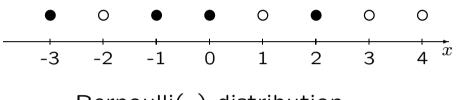
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

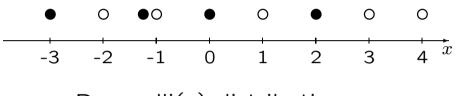
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

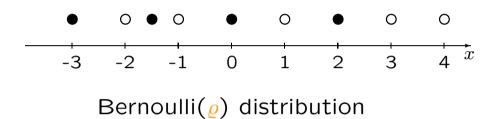
to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

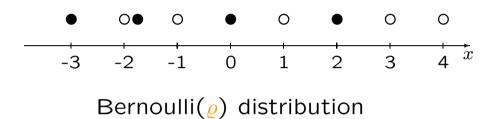
Particles try to jump

to the right with rate p, to the left with rate q = 1 - p < p.



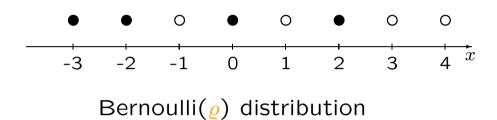
Particles try to jump

to the right with rate p, to the left with rate q = 1 - p < p.



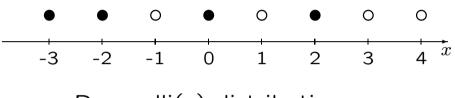
Particles try to jump

to the right with rate p, to the left with rate q = 1 - p < p.



Particles try to jump

to the right with rate p, to the left with rate q = 1 - p < p.



 $Bernoulli(\varrho)$ distribution

Particles try to jump

to the right with rate p, to the left with rate q = 1 - p < p.

The jump is suppressed if the destination site is occupied by another particle.

The Bernoulli(ϱ) distribution is time-stationary for any ($0 \le \varrho \le 1$). Any translation-invariant stationary distribution is a mixture of Bernoullis.

Let T and X be some large-scale time and space parameters.

Let T and X be some large-scale time and space parameters. \rightsquigarrow Set initially $\varrho = \varrho(T = 0, X)$ to be the density at position $x = X/\varepsilon$. (Changes on the large scale.)

Let T and X be some large-scale time and space parameters. \rightarrow Set initially $\varrho = \varrho(T = 0, X)$ to be the density at position $x = X/\varepsilon$. (Changes on the large scale.) $\rightarrow \varrho(T, X)$ is the density of particles after a long time $t = T/\varepsilon$ at position $x = X/\varepsilon$.

Let T and X be some large-scale time and space parameters. \rightarrow Set initially $\varrho = \varrho(T = 0, X)$ to be the density at position $x = X/\varepsilon$. (Changes on the large scale.)

$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1-\varrho)] = 0 \quad \text{(inviscid Burgers)}$$

Let T and X be some large-scale time and space parameters. \rightarrow Set initially $\varrho = \varrho(T = 0, X)$ to be the density at position $x = X/\varepsilon$. (Changes on the large scale.)

$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1-\varrho)] = 0 \quad \text{(inviscid Burgers})$$
$$\frac{\partial}{\partial T} \varrho + a[1-2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$

Let T and X be some large-scale time and space parameters. \rightarrow Set initially $\varrho = \varrho(T = 0, X)$ to be the density at position $x = X/\varepsilon$. (Changes on the large scale.)

$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1-\varrho)] = 0 \quad \text{(inviscid Burgers)}$$
$$\frac{\partial}{\partial T} \varrho + a[1-2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$
$$\frac{d}{dT} \varrho(T, X(T)) = 0$$

Let T and X be some large-scale time and space parameters. \rightarrow Set initially $\varrho = \varrho(T = 0, X)$ to be the density at position $x = X/\varepsilon$. (Changes on the large scale.)

$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1-\varrho)] = 0 \quad \text{(inviscid Burgers)}$$
$$\frac{\partial}{\partial T} \varrho + a[1-2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$
$$\frac{\partial}{\partial T} \varrho + \frac{\mathrm{d}X(T)}{\mathrm{d}T} \cdot \frac{\partial}{\partial X} \varrho = \frac{\mathrm{d}}{\mathrm{d}T} \varrho(T, X(T)) = 0$$

Let T and X be some large-scale time and space parameters. \rightarrow Set initially $\varrho = \varrho(T = 0, X)$ to be the density at position $x = X/\varepsilon$. (Changes on the large scale.)

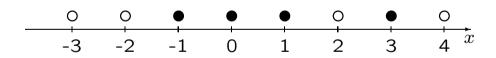
$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1-\varrho)] = 0 \quad \text{(inviscid Burgers)}$$
$$\frac{\partial}{\partial T} \varrho + a[1-2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$
$$\frac{\partial}{\partial T} \varrho + \frac{\mathrm{d}X(T)}{\mathrm{d}T} \cdot \frac{\partial}{\partial X} \varrho = \frac{\mathrm{d}}{\mathrm{d}T} \varrho(T, X(T)) = 0$$

Let T and X be some large-scale time and space parameters. \rightarrow Set initially $\varrho = \varrho(T = 0, X)$ to be the density at position $x = X/\varepsilon$. (Changes on the large scale.)

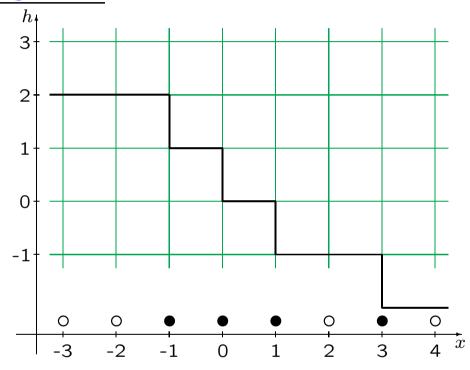
 $\rightsquigarrow \varrho(T, X)$ is the density of particles after a long time $t = T/\varepsilon$ at position $x = X/\varepsilon$. It satisfies, with a := p - q,

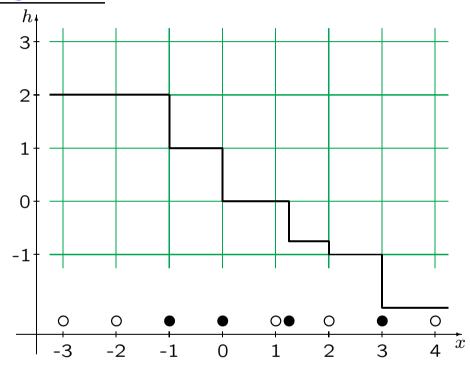
$$\frac{\partial}{\partial T} \varrho + \frac{\partial}{\partial X} a[\varrho(1-\varrho)] = 0 \quad \text{(inviscid Burgers)}$$
$$\frac{\partial}{\partial T} \varrho + a[1-2\varrho] \cdot \frac{\partial}{\partial X} \varrho = 0 \quad \text{(while smooth)}$$
$$\frac{\partial}{\partial T} \varrho + \frac{\mathrm{d}X(T)}{\mathrm{d}T} \cdot \frac{\partial}{\partial X} \varrho = \frac{\mathrm{d}}{\mathrm{d}T} \varrho(T, X(T)) = 0$$

→ The characteristic speed $C(\varrho) := a[1 - 2\varrho].$ (ϱ is constant along $\dot{X}(T) = C(\varrho).$) 2. ASEP: Surface growth

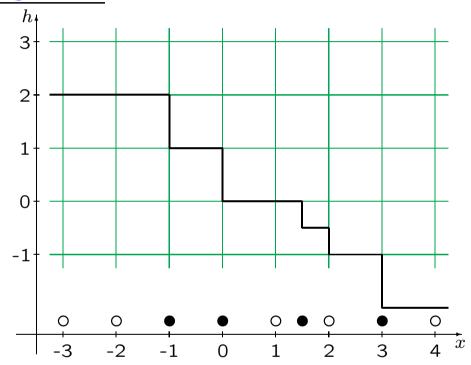


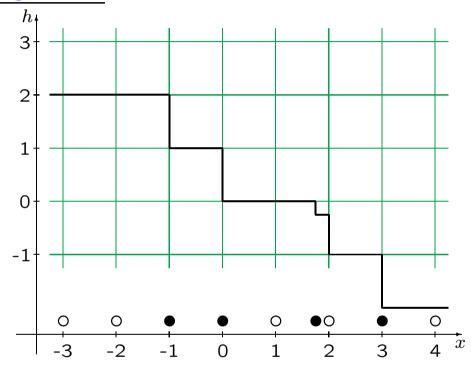
2. ASEP: Surface growth

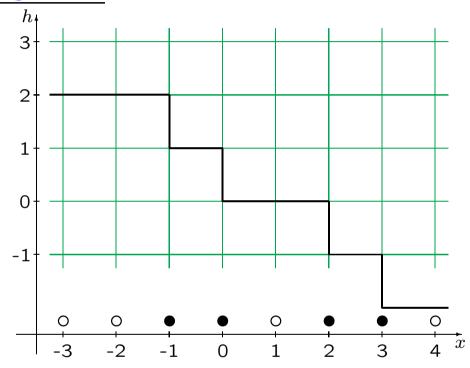


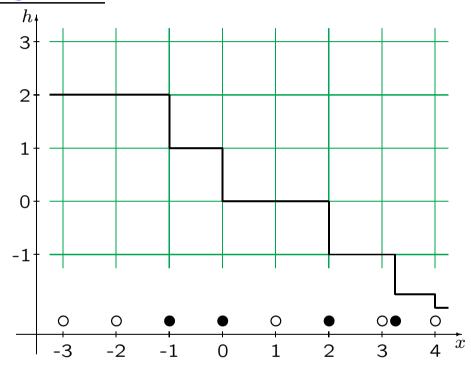


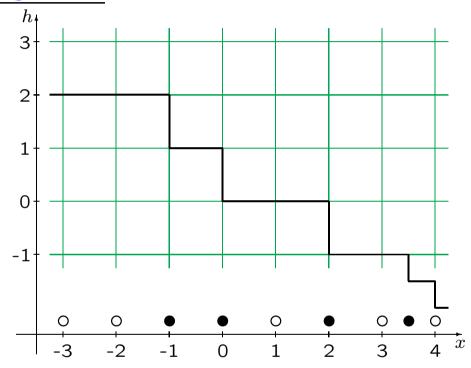
37

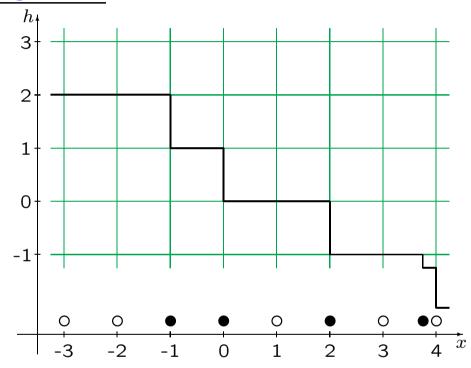


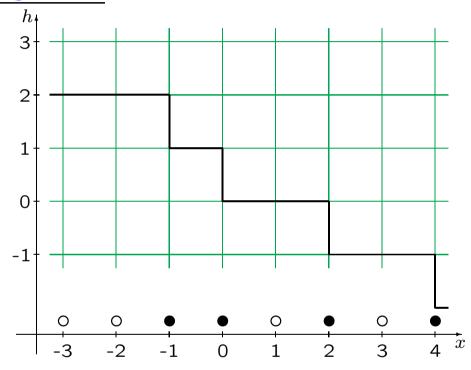


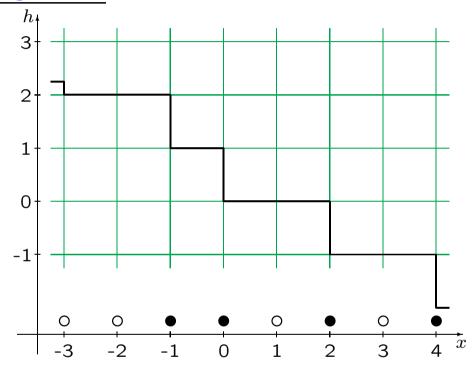




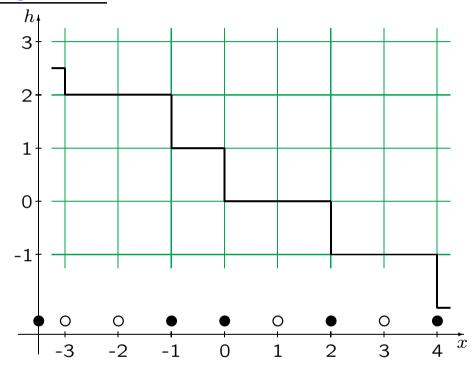


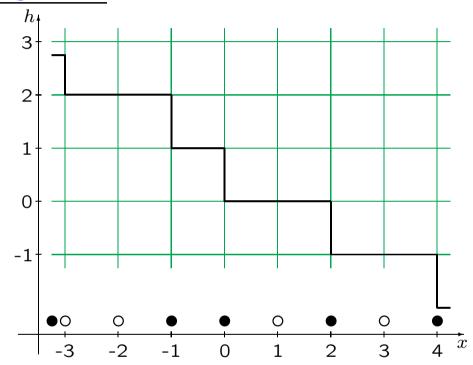


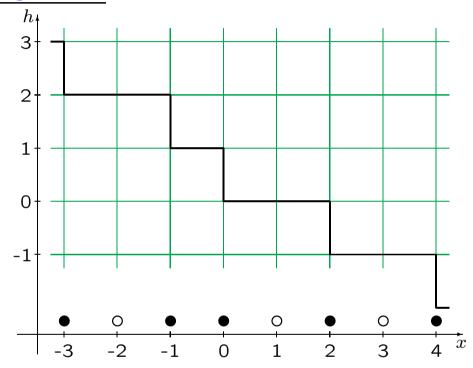


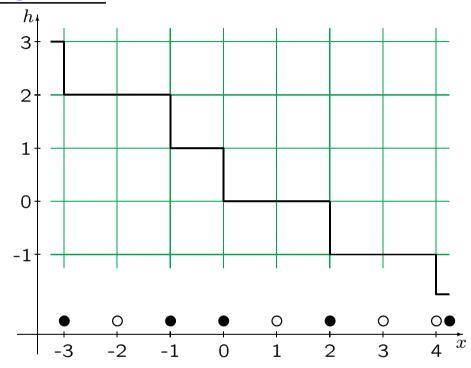


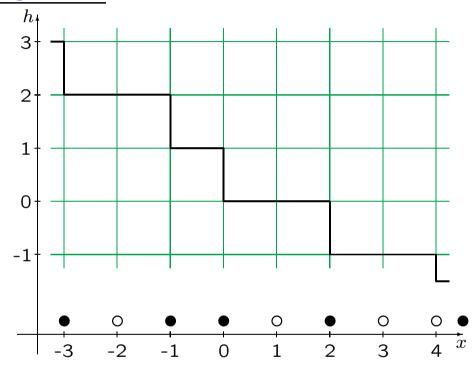
45

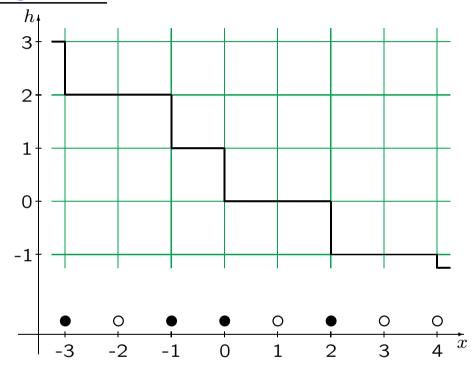


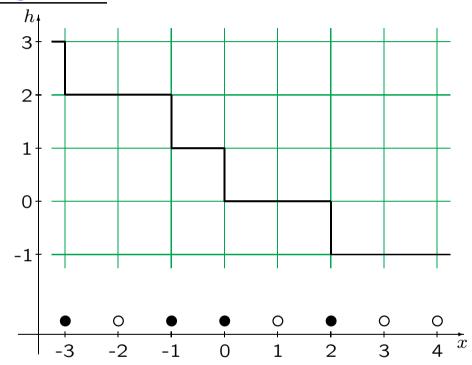




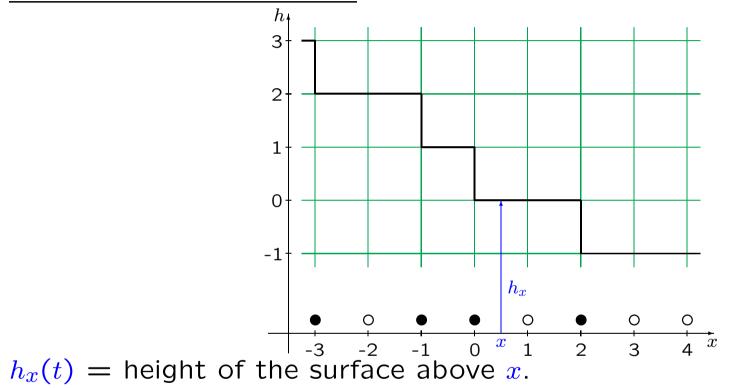




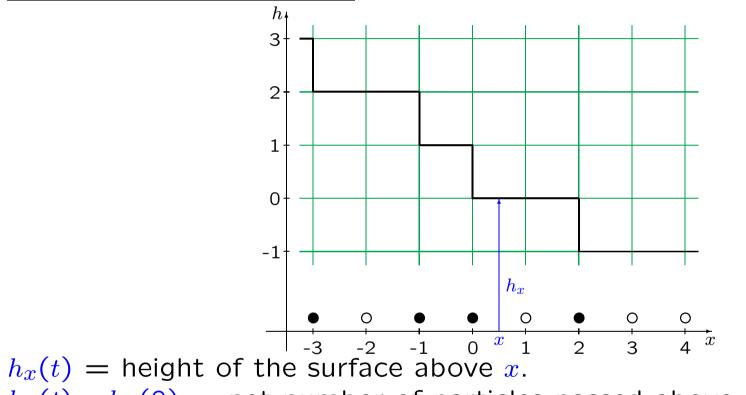




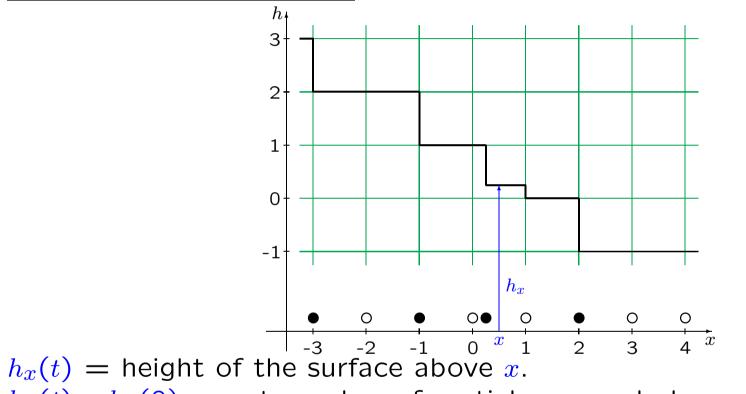




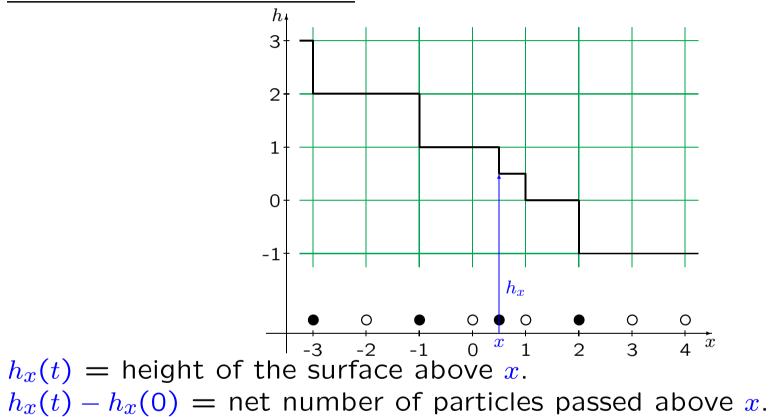




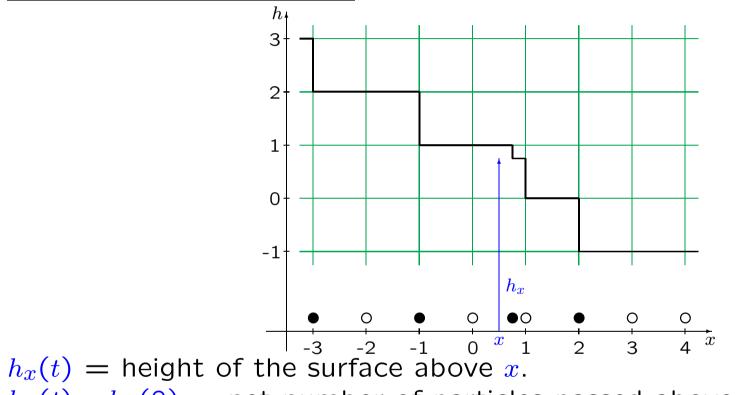




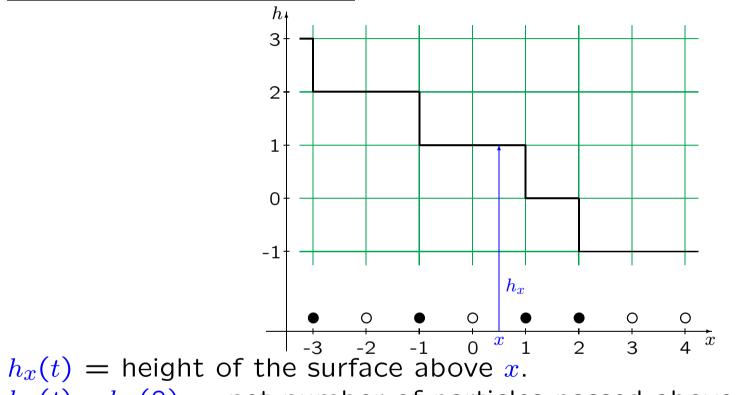




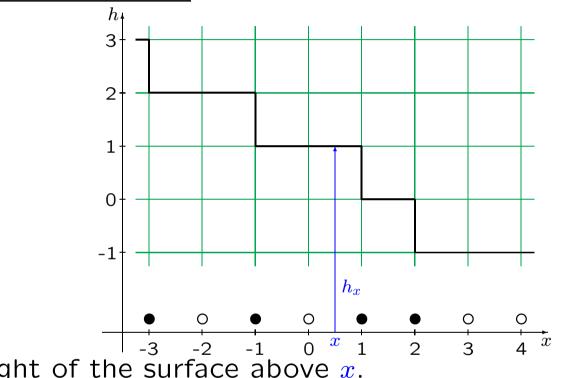






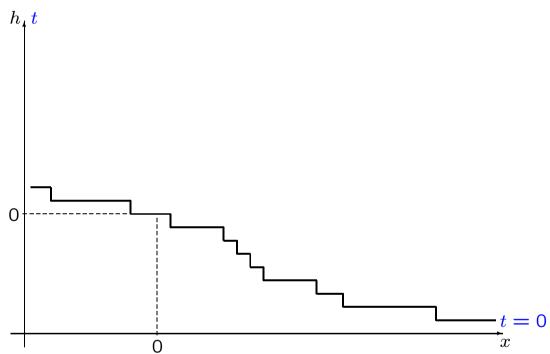




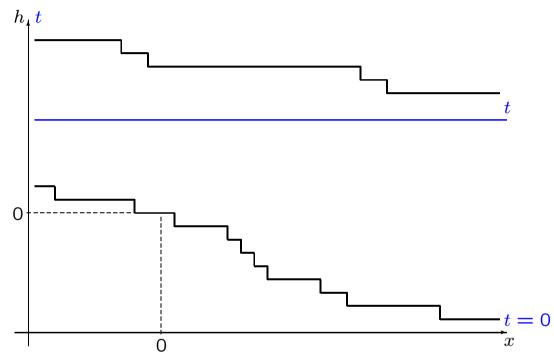


 $h_x(t) =$ height of the surface above x. $h_x(t) - h_x(0) =$ net number of particles passed above x. $h_{Vt}(t) =$ net number of particles passed through the moving window at Vt ($V \in \mathbb{R}$).

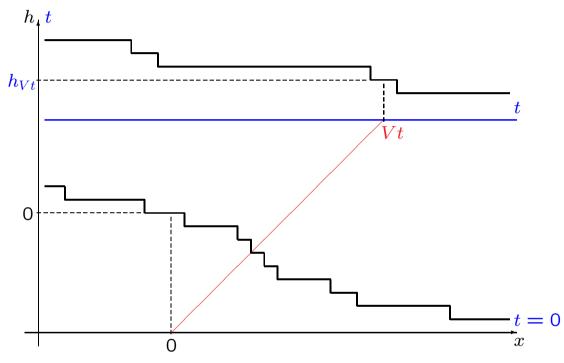




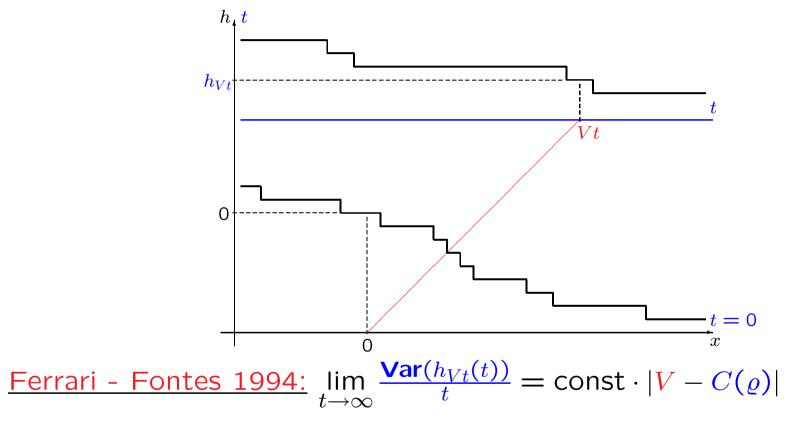




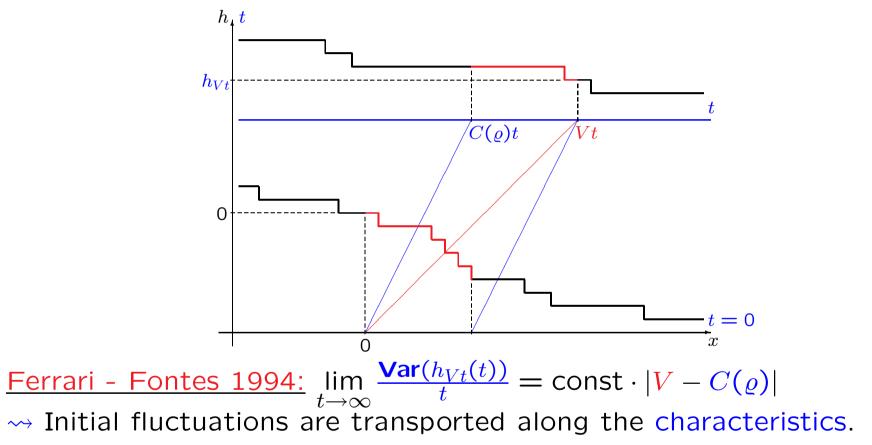


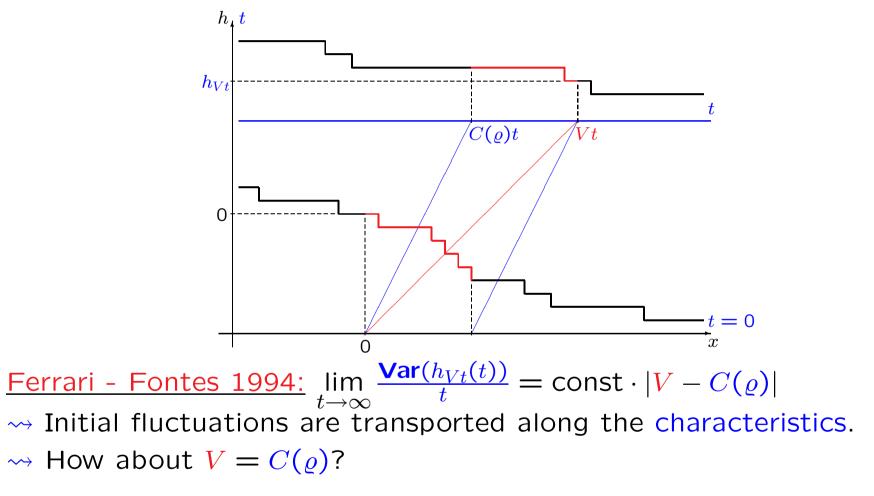












Conjecture:

$$\lim_{t \to \infty} \frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} = [\text{sg. non trivial}].$$

Conjecture:

$$\lim_{t\to\infty} \frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} = [\text{sg. non trivial}].$$

Theorem (B., Seppäläinen): For any $0 < \rho < 1$, and any q < p,

$$\begin{split} 0 &< \liminf_{t \to \infty} \frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} \\ &\leq \limsup_{t \to \infty} \frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} < \infty. \end{split}$$

Conjecture:

$$\lim_{t\to\infty} \frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} = [\operatorname{sg. non trivial}].$$

Theorem (B., Seppäläinen): For any $0 < \rho < 1$, and any q < p,

$$\begin{split} 0 &< \liminf_{t \to \infty} \frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} \\ &\leq \limsup_{t \to \infty} \frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} < \infty. \end{split}$$

Corollary: The corresponding scaling of the diffusivity is also proved.

Limit distributions (not yet controlling the second moment) in terms of the Tracy-Widom distribution (GUE random matrices) were found by Baik, Deift and Johansson 1999, Johansson 2000, and Ferrari and Spohn 2006 for the *totally* asymmetric exclusion (TASEP: p = 1, q = 0).

Limit distributions (not yet controlling the second moment) in terms of the Tracy-Widom distribution (GUE random matrices) were found by Baik, Deift and Johansson 1999, Johansson 2000, and Ferrari and Spohn 2006 for the *totally* asymmetric exclusion (TASEP: p = 1, q = 0).

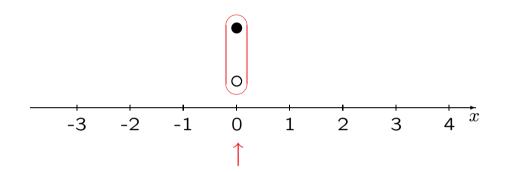
<u>Method was:</u> Last passage percolation, heavy combinatorics and a-symptotic analysis.

Limit distributions (not yet controlling the second moment) in terms of the Tracy-Widom distribution (GUE random matrices) were found by Baik, Deift and Johansson 1999, Johansson 2000, and Ferrari and Spohn 2006 for the *totally* asymmetric exclusion (TASEP: p = 1, q = 0).

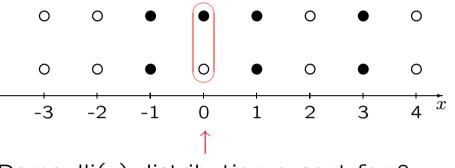
<u>Method was:</u> Last passage percolation, heavy combinatorics and a-symptotic analysis.

 → We needed to get rid of these tools. Premises: Cator and Groeneboom 2006 (Hammersley's process), B., Cator and Seppäläinen 2006 (TASEP, last passage).



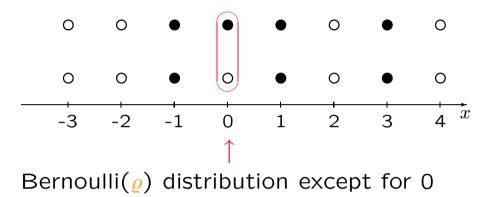


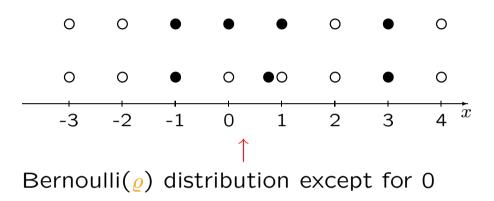


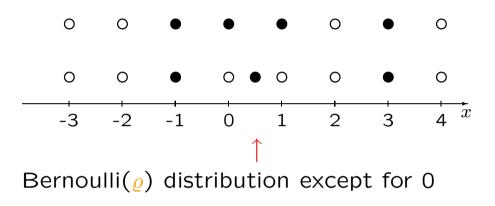


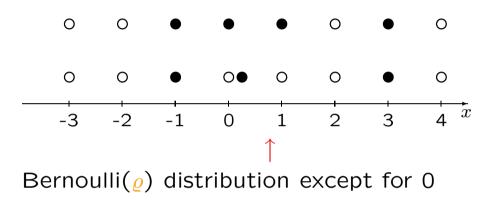
Bernoulli(ϱ) distribution except for 0



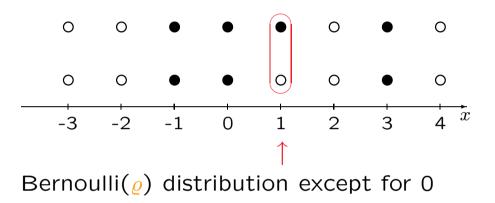


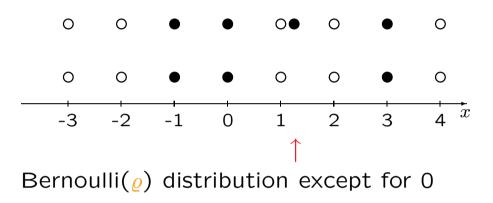


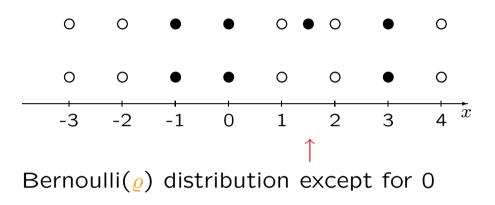


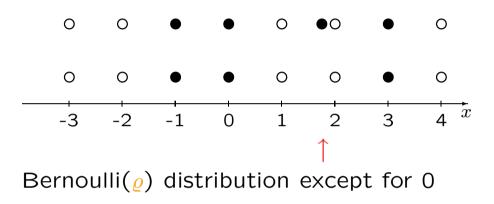




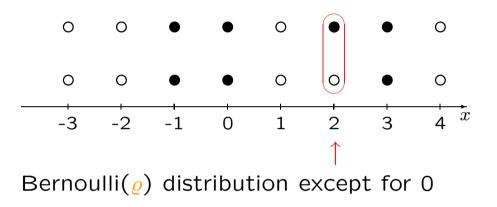




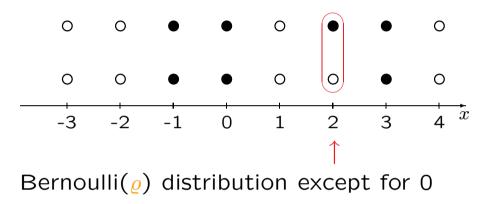






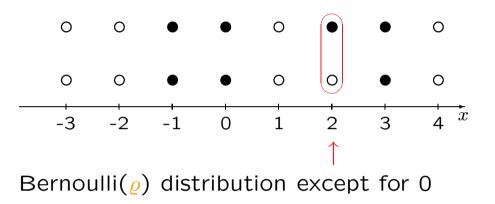






<u>Coupling</u>: A single discrepancy is always conserved = the second class particle. Its location at time t is Q(t).





<u>Coupling</u>: A single discrepancy is always conserved = the second class particle. Its location at time t is Q(t).

The second class particle is a highly nontrivial object. For example, the Bernoulli(ϱ) distribution is *not* stationary as seen by the second class particle.

Theorem:

 $\mathbf{E}(Q(t)) = C(\varrho)t$

(characteristic speed),

Theorem:

 $\mathbf{E}(Q(t)) = C(\varrho)t$

(characteristic speed), and

 $\operatorname{Var}(h_{Vt}(t)) = \operatorname{const} \cdot \operatorname{E}|Vt - Q(t)|.$

Theorem:

 $\mathbf{E}(Q(t)) = C(\varrho)t$

(characteristic speed), and

$$\operatorname{Var}(h_{Vt}(t)) = \operatorname{const} \cdot \mathbf{E}|Vt - Q(t)|.$$

Method of proof: martingale arguments, time-reversal, and conservation of particles.

Theorem:

 $\mathbf{E}(Q(t)) = C(\varrho)t$

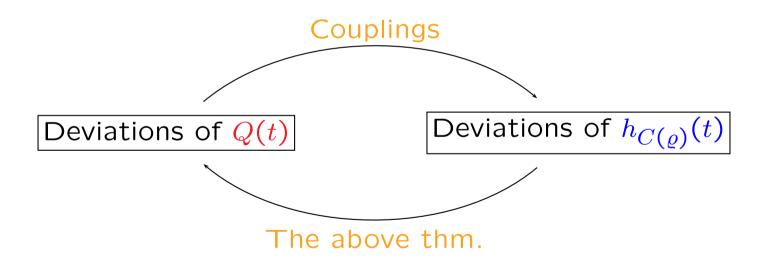
(characteristic speed), and

$$\operatorname{Var}(h_{Vt}(t)) = \operatorname{const} \cdot \operatorname{E}|Vt - Q(t)|.$$

Method of proof: martingale arguments, time-reversal, and conservation of particles.

The proof is based on ideas of Bálint Tóth, he said these ideas were standard.

Main idea for prooving $t^{1/3}$ scaling:



The coupling measure

Let $\lambda < \varrho$, and

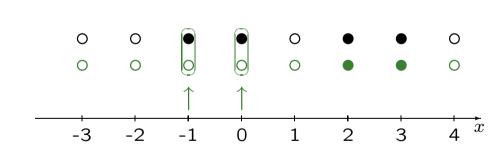
$$\mu \begin{pmatrix} \circ \\ \circ \end{pmatrix} = 1 - \varrho, \quad \mu \begin{pmatrix} \bullet \\ \circ \end{pmatrix} = \varrho - \lambda, \quad \mu \begin{pmatrix} \bullet \\ \bullet \end{pmatrix} = \lambda.$$

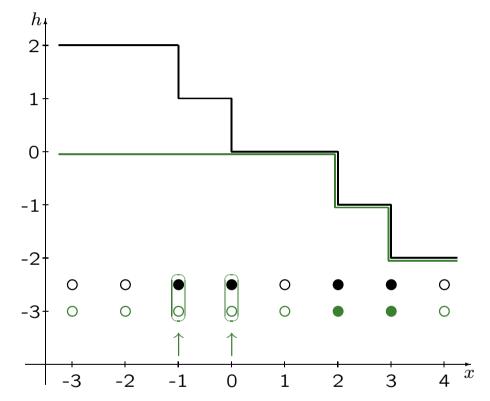
The coupling measure

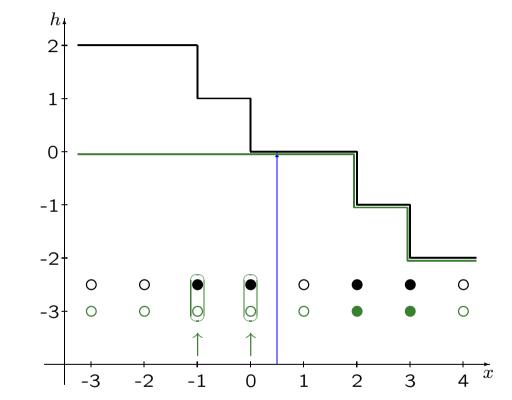
Let $\lambda < \varrho$, and

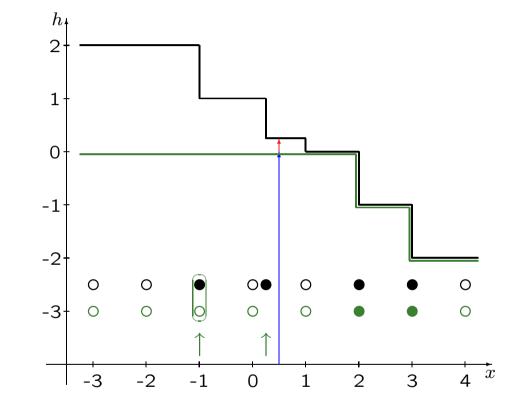
$$\mu \begin{pmatrix} \circ \\ \circ \end{pmatrix} = 1 - \varrho, \quad \mu \begin{pmatrix} \bullet \\ \circ \end{pmatrix} = \varrho - \lambda, \quad \mu \begin{pmatrix} \bullet \\ \bullet \end{pmatrix} = \lambda.$$

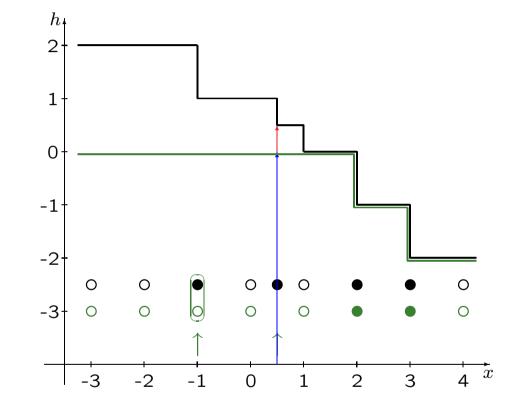
Then the "upper" marginal is Bernoulli(ϱ), and the "lower" marginal is Bernoulli(λ).

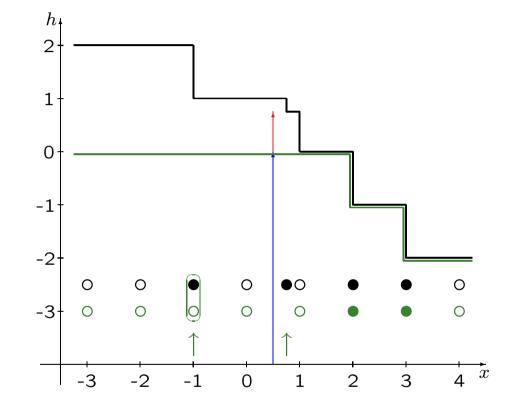


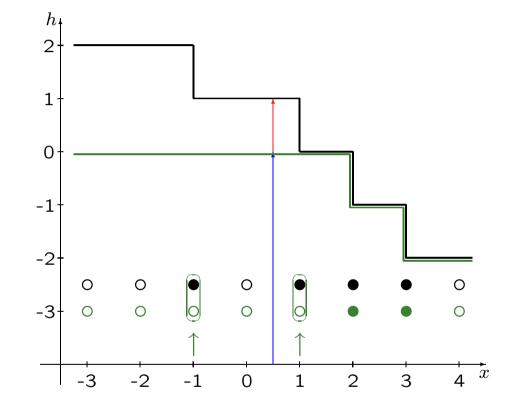


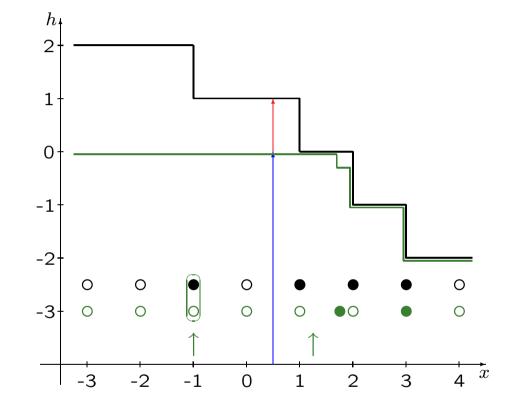


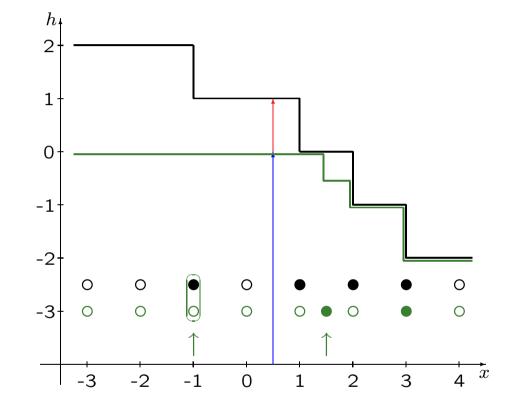


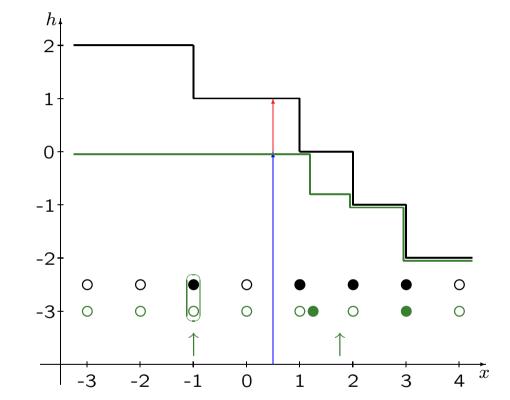


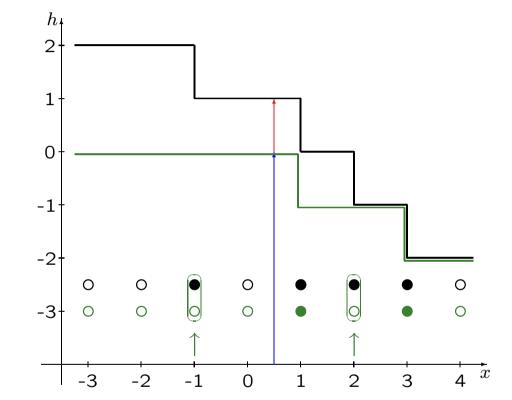




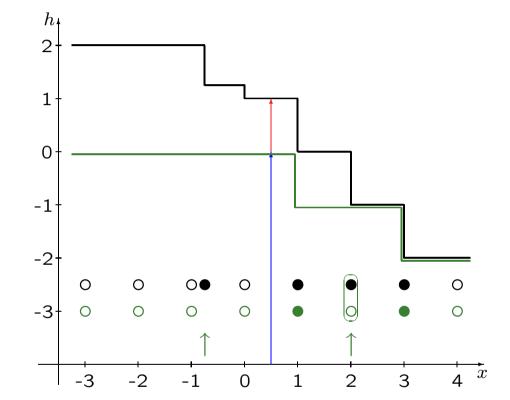


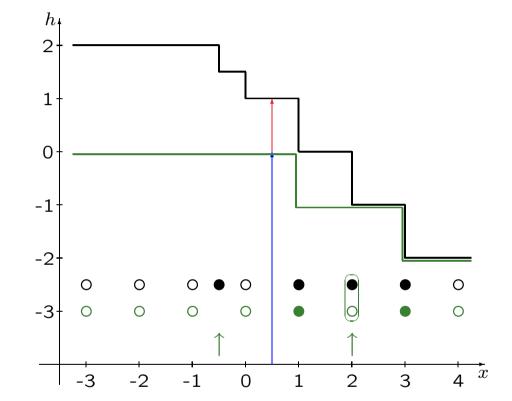




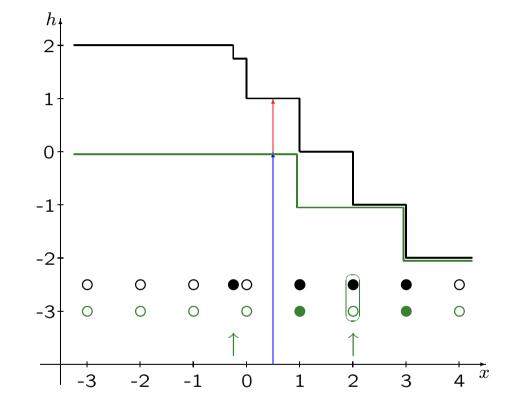


 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt ($V \in \mathbb{R}$).

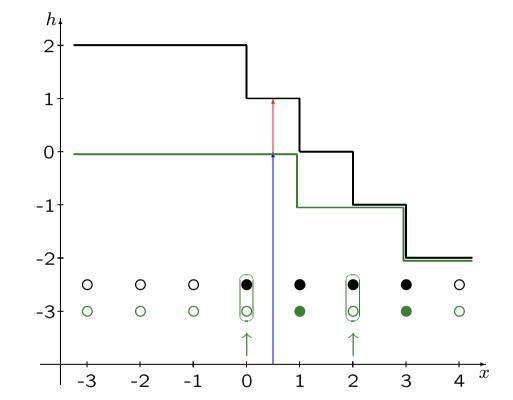




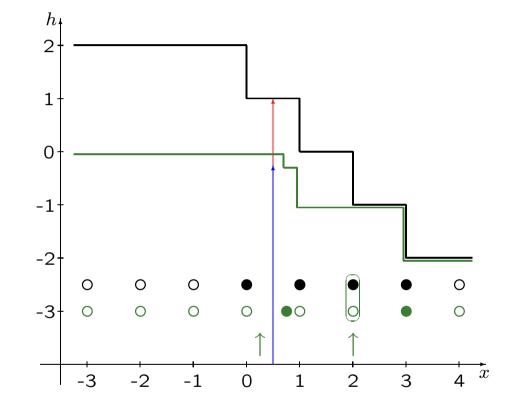
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt ($V \in \mathbb{R}$).



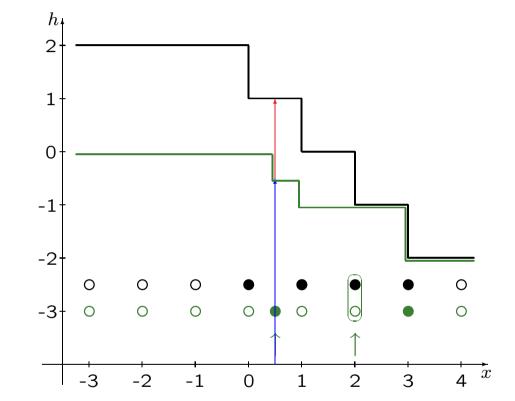
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt ($V \in \mathbb{R}$).



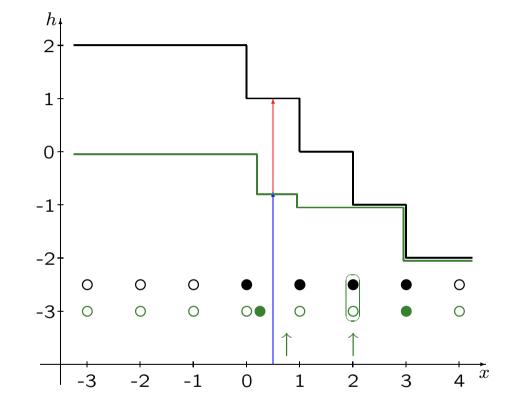
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt ($V \in \mathbb{R}$).



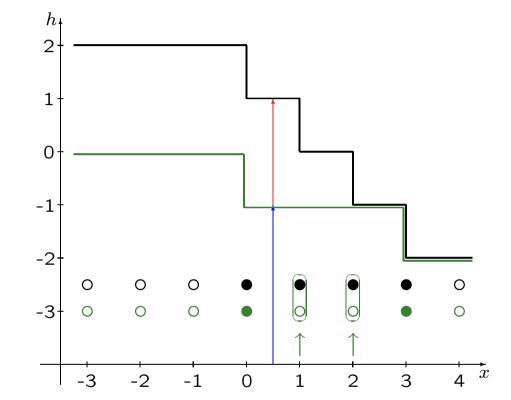
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt ($V \in \mathbb{R}$).



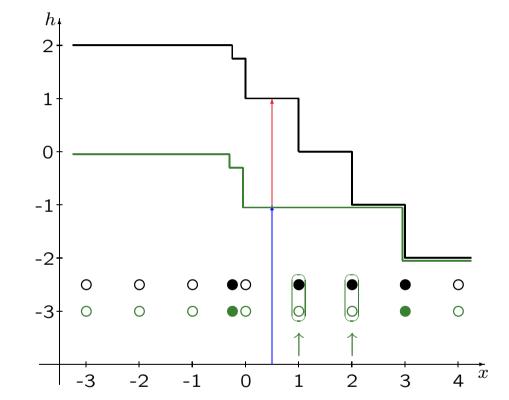
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt ($V \in \mathbb{R}$).

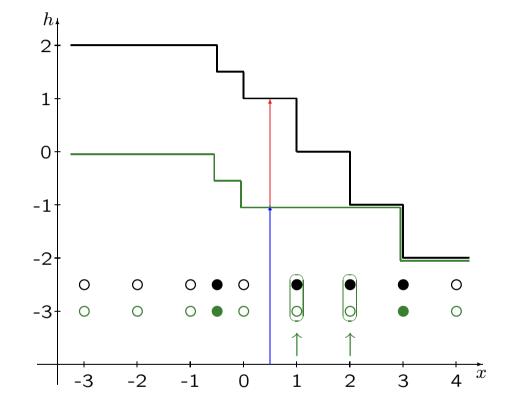


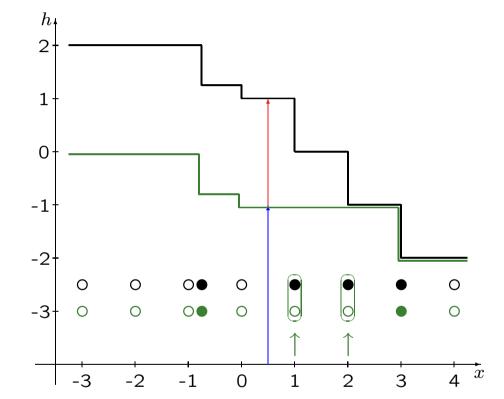
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt ($V \in \mathbb{R}$).

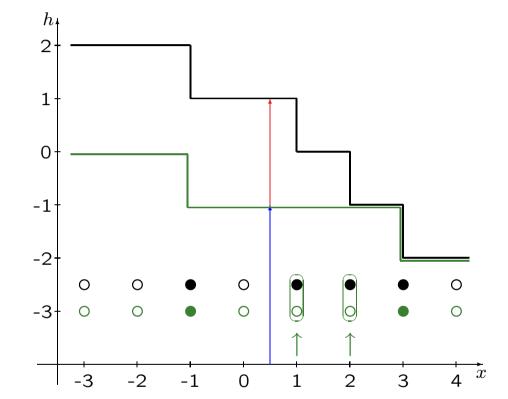


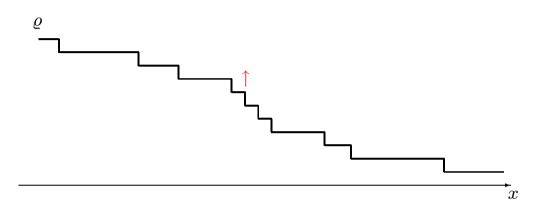
 $h_{Vt}(t) - h_{Vt}(t)$ = the net number of \uparrow 's passed through the moving window at Vt ($V \in \mathbb{R}$).

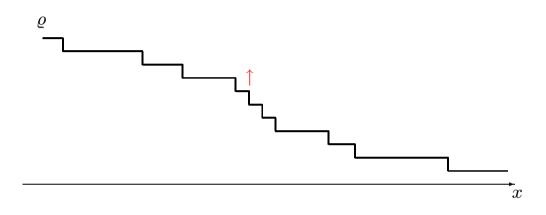




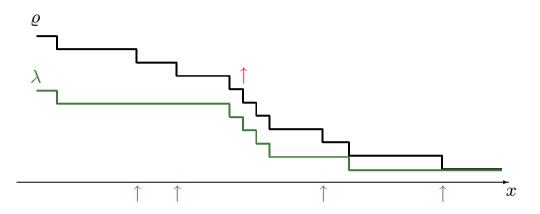




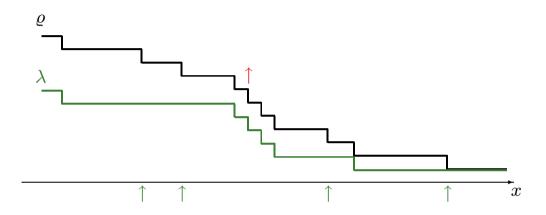




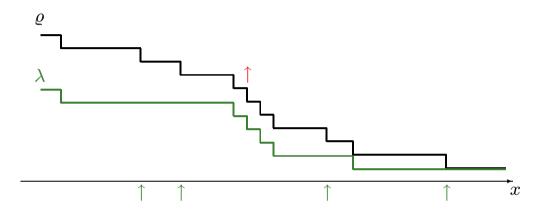
Connect Q(t)



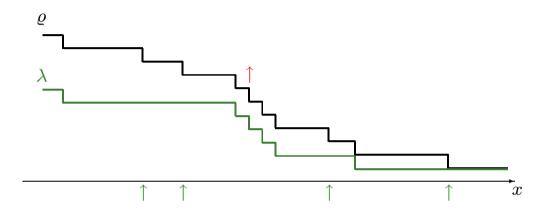
Connect Q(t) with the \uparrow 's



Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

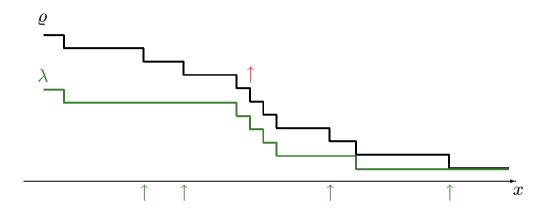


Connect Q(t) with the \uparrow 's (this needs nontrivial couplings): P{Q(t) is too large}



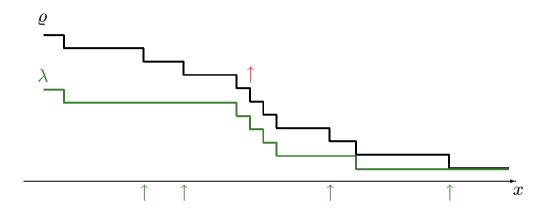
Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

 $P{Q(t) \text{ is too large}} \leq P{\text{too many } \uparrow \text{'s have crossed } C(\varrho)t}$



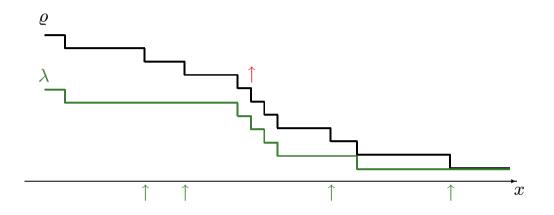
Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

$$\begin{split} \mathbf{P}\{Q(t) \text{ is too large}\} &\leq \mathbf{P}\{\text{too many }\uparrow\text{'s have crossed } C(\varrho)t\}\\ &\leq \mathbf{P}\{h_{C(\varrho)t}(t) - h_{C(\varrho)t}(t) \text{ is too large}\}. \end{split}$$



Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

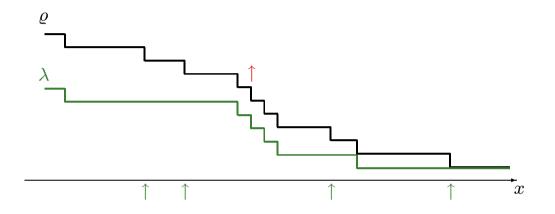
$$\begin{split} \mathbf{P}\{Q(t) \text{ is too large}\} &\leq \mathbf{P}\{\text{too many} \uparrow \text{'s have crossed } C(\varrho)t\} \\ &\leq \mathbf{P}\{h_{C(\varrho)t}(t) - h_{C(\varrho)t}(t) \text{ is too large}(\lambda)\}. \end{split}$$



Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

$$\begin{split} \mathbf{P}\{Q(t) \text{ is too large}\} &\leq \mathbf{P}\{\text{too many }\uparrow\text{'s have crossed } C(\varrho)t\}\\ &\leq \mathbf{P}\{h_{C(\varrho)t}(t) - h_{C(\varrho)t}(t) \text{ is too large}(\lambda)\}. \end{split}$$

Optimize "too large(λ)" in λ ,



Connect Q(t) with the \uparrow 's (this needs nontrivial couplings):

$$\begin{split} \mathbf{P}\{Q(t) \text{ is too large}\} &\leq \mathbf{P}\{\text{too many} \uparrow \text{'s have crossed } C(\varrho)t\} \\ &\leq \mathbf{P}\{h_{C(\varrho)t}(t) - h_{C(\varrho)t}(t) \text{ is too large}(\lambda)\}. \end{split}$$

Optimize "too large(λ)" in λ , use Chebyshev's inequality and relate $Var(h_{C(\varrho)t}(t))$ to $Var(h_{C(\varrho)t}(t))$.

The computations result in

$$\mathbf{P}\{\mathbf{Q}(t) - C(\varrho)t \ge u\} \le c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{C(\varrho)t}(t))$$

$$\mathbf{P}\{Q(t) - C(\varrho)t \ge u\} \le c \cdot \frac{t^2}{u^4} \cdot \mathsf{Var}(h_{C(\varrho)t}(t))$$

$$\begin{aligned} \mathbf{P}\{\mathbf{Q}(t) - C(\varrho)t \geq u\} &\leq c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{C(\varrho)t}(t)) \\ & \overset{\mathsf{Thm}}{=} c \cdot \frac{t^2}{u^4} \cdot \mathbf{E}[\mathbf{Q}(t) - C(\varrho)t]. \end{aligned}$$

$$\begin{aligned} \mathbf{P}\{\boldsymbol{Q(t)} - C(\varrho)t \geq u\} &\leq c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{C(\varrho)t}(t)) \\ & \overset{\mathsf{Thm}}{=} c \cdot \frac{t^2}{u^4} \cdot \mathbf{E}|\boldsymbol{Q(t)} - C(\varrho)t|. \end{aligned}$$

With

$$\widetilde{Q}(t) := Q(t) - C(\varrho)t$$
 and $E := \mathbf{E}|\widetilde{Q}(t)|,$

we have (with a similar lower deviation bound)

$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E.$$

$$\begin{split} \mathbf{P}\{\boldsymbol{Q(t)} - C(\varrho)t \geq u\} &\leq c \cdot \frac{t^2}{u^4} \cdot \mathbf{Var}(h_{C(\varrho)t}(t)) \\ & \overset{\mathsf{Thm}}{=} c \cdot \frac{t^2}{u^4} \cdot \mathbf{E}[\boldsymbol{Q(t)} - C(\varrho)t]. \end{split}$$

With

$$\widetilde{Q}(t) := Q(t) - C(\varrho)t$$
 and $E := \mathbf{E}|\widetilde{Q}(t)|,$

we have (with a similar lower deviation bound)

$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E.$$

Claim: this already implies the $t^{2/3}$ upper bound:

We had
$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E.$$

We had
$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.
$$E = \mathbf{E}|\tilde{Q}(t)| = \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > u\} \, \mathrm{d}u$$

We had
$$\mathbf{P}\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.

$$E = \mathbf{E}|\tilde{Q}(t)| = \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > u\} \, \mathrm{d}u$$

$$= E \int_0^\infty \mathbf{P}\{|\tilde{Q}(t)| > vE\} \, \mathrm{d}v$$

We had
$$P\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.

$$E = E|\tilde{Q}(t)| = \int_0^\infty P\{|\tilde{Q}(t)| > u\} du$$

$$= E \int_0^\infty P\{|\tilde{Q}(t)| > vE\} dv$$

$$\le E \int_{1/2}^\infty P\{|\tilde{Q}(t)| > vE\} dv + \frac{1}{2}E$$

We had
$$P\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.

$$E = E|\tilde{Q}(t)| = \int_0^\infty P\{|\tilde{Q}(t)| > u\} du$$

$$= E \int_0^\infty P\{|\tilde{Q}(t)| > vE\} dv$$

$$\le E \int_{1/2}^\infty P\{|\tilde{Q}(t)| > vE\} dv + \frac{1}{2}E$$

$$\le c \cdot \frac{t^2}{E^2} + \frac{1}{2}E,$$

that is, $E^3 \leq c \cdot t^2$.

We had
$$P\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.

$$E = E|\tilde{Q}(t)| = \int_0^\infty P\{|\tilde{Q}(t)| > u\} du$$

$$= E \int_0^\infty P\{|\tilde{Q}(t)| > vE\} dv$$

$$\le E \int_{1/2}^\infty P\{|\tilde{Q}(t)| > vE\} dv + \frac{1}{2}E$$

$$\le c \cdot \frac{t^2}{E^2} + \frac{1}{2}E,$$

that is, $E^3 \leq c \cdot t^2$.

$$\operatorname{Var}(h_{C(\varrho)t}(t)) \stackrel{\mathsf{Thm}}{=} \operatorname{const.} \cdot \operatorname{E}[Q(t) - C(\varrho)t]$$

We had
$$P\{|\tilde{Q}(t)| > u\} \le c \cdot \frac{t^2}{u^4} \cdot E$$
.

$$E = E|\tilde{Q}(t)| = \int_0^\infty P\{|\tilde{Q}(t)| > u\} du$$

$$= E \int_0^\infty P\{|\tilde{Q}(t)| > vE\} dv$$

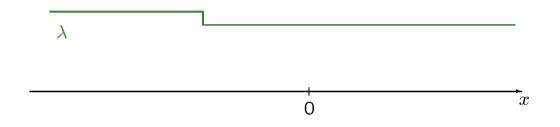
$$\le E \int_{1/2}^\infty P\{|\tilde{Q}(t)| > vE\} dv + \frac{1}{2}E$$

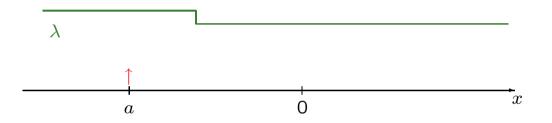
$$\le c \cdot \frac{t^2}{E^2} + \frac{1}{2}E,$$

that is, $E^3 \leq c \cdot t^2$.

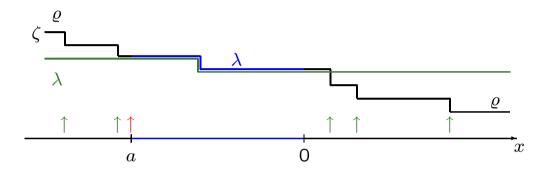
$$Var(h_{C(\varrho)t}(t)) \stackrel{\mathsf{Thm}}{=} const. \cdot \mathbf{E}|Q(t) - C(\varrho)t|$$
$$= const. \cdot E \le c \cdot t^{2/3}.$$

136

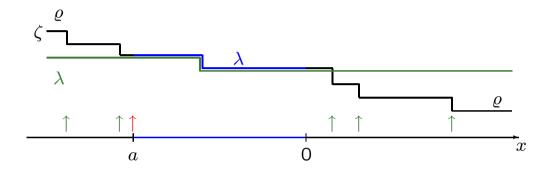




Let
$$Q^{a}(0) = a < 0$$
.

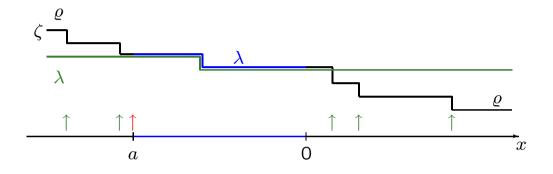


Let
$$Q^a(0) = a < 0$$
.



Let $Q^a(0) = a < 0$. If $Q^a(t) \le C(\varrho)t$, then the \uparrow 's have not crossed the path $C(\varrho)t$ from left to right:

 $\mathbf{P}\{Q^{a}(t) \leq C(\varrho)t\} \leq \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$

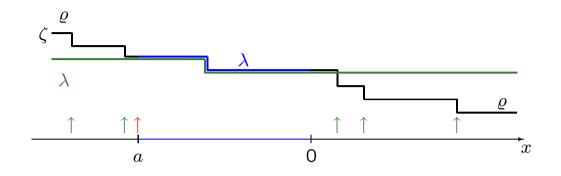


Let $Q^a(0) = a < 0$. If $Q^a(t) \le C(\varrho)t$, then the \uparrow 's have not crossed the path $C(\varrho)t$ from left to right:

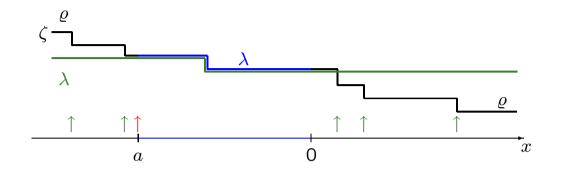
$$\mathbf{P}\{Q^{a}(t) \leq C(\varrho)t\} \leq \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

Therefore:

$$1 \leq \mathbf{P}\{Q^{a}(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}.$$

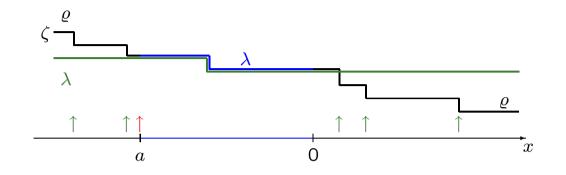


 $1 \leq \mathbf{P}\{Q^{a}(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$



 $1 \leq \mathbf{P}\{Q^{a}(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$

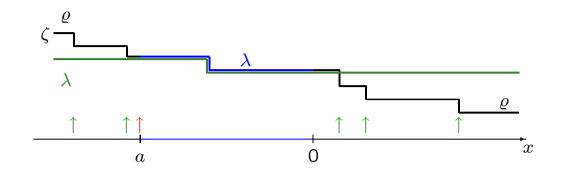
 \rightsquigarrow Set a so that $E(Q^a(t)) < C(\varrho)t$,



 $1 \leq \mathbf{P}\{Q^{a}(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$

 \rightsquigarrow Set a so that $E(Q^a(t)) < C(\varrho)t$,

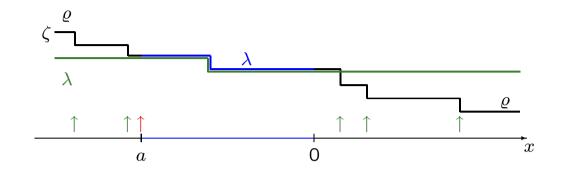
 $\rightarrow E(h_{C(\varrho)t}(t)) - E(h_{C(\varrho)t}(t)) \sim t(\varrho - \lambda)^2 > 0$ would be the case, if ζ was Bernoulli(ϱ) distributed.



 $1 \leq \mathbf{P}\{Q^{a}(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$

 \rightsquigarrow Set a so that $E(Q^a(t)) < C(\varrho)t$,

 $\rightarrow E(h_{C(\varrho)t}(t)) - E(h_{C(\varrho)t}(t)) \sim t(\varrho - \lambda)^2 > 0$ would be the case, if ζ was Bernoulli(ϱ) distributed. Instead, $E(h_{C(\varrho)t}(t))$ will have a harmless Radon-Nikodym factor.



 $1 \leq \mathbf{P}\{Q^{a}(t) > C(\varrho)t\} + \mathbf{P}\{h_{C(\varrho)t}(t) < h_{C(\varrho)t}(t)\}$

 \rightsquigarrow Set a so that $E(Q^a(t)) < C(\varrho)t$,

 $\rightarrow E(h_{C(\varrho)t}(t)) - E(h_{C(\varrho)t}(t)) \sim t(\varrho - \lambda)^2 > 0$ would be the case, if ζ was Bernoulli(ϱ) distributed. Instead, $E(h_{C(\varrho)t}(t))$ will have a harmless Radon-Nikodym factor.

 \Rightarrow Both probabilities are deviation probabilities.

Apply Markov's inequality on the first, Chebyshev's on the second probability (use again the connection between $Var(h_{C(\varrho)t}(t))$ and $Var(h_{C(\varrho)t}(t))$).

Apply Markov's inequality on the first, Chebyshev's on the second probability (use again the connection between $Var(h_{C(\rho)t}(t))$ and $Var(h_{C(\rho)t}(t))$).

The correct scaling of the parameters is: $\varrho - \lambda \sim t^{-1/3}$, $a \sim -t^{2/3}$.

Apply Markov's inequality on the first, Chebyshev's on the second probability (use again the connection between $Var(h_{C(\rho)t}(t))$ and $Var(h_{C(\rho)t}(t))$).

The correct scaling of the parameters is: $\varrho - \lambda \sim t^{-1/3}$, $a \sim -t^{2/3}$. In this case

$$1 \leq c_1 \cdot \frac{\mathbf{E}(|\widetilde{Q^a}(t)|)}{t^{2/3}} + c_2 \cdot \frac{\mathbf{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}$$

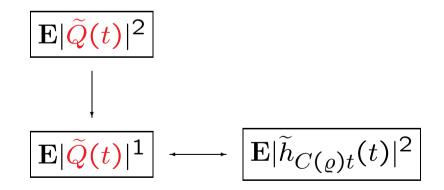
Apply Markov's inequality on the first, Chebyshev's on the second probability (use again the connection between $Var(h_{C(\rho)t}(t))$ and $Var(h_{C(\rho)t}(t))$).

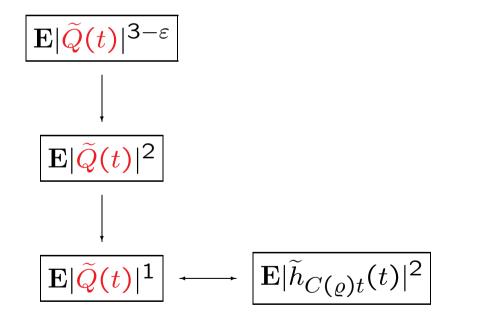
The correct scaling of the parameters is: $\varrho - \lambda \sim t^{-1/3}$, $a \sim -t^{2/3}$. In this case

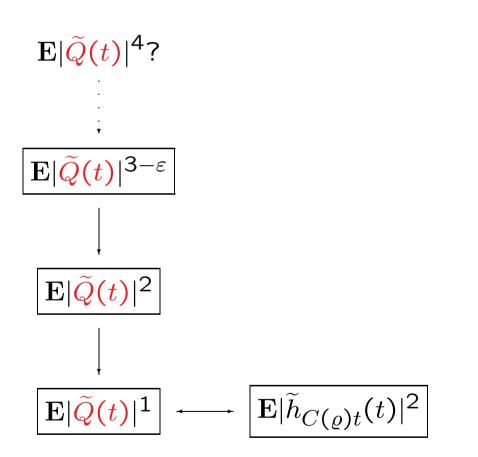
$$1 \leq c_1 \cdot \frac{\mathbf{E}(|\widetilde{Q^a}(t)|)}{t^{2/3}} + c_2 \cdot \frac{\mathbf{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}$$
$$\underset{c}{\overset{\mathsf{Thm}}{=}} c \cdot \frac{\mathbf{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}.$$

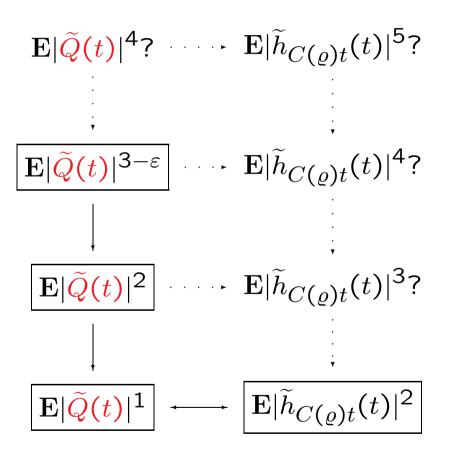
151

$$\mathbf{E}|\tilde{Q}(t)|^{1} \longleftrightarrow \mathbf{E}|\tilde{h}_{C(\varrho)t}(t)|^{2}$$









 \rightarrow What is the limit $\lim_{t\rightarrow\infty}\frac{\mathrm{Var}(h_{C(\varrho)t}(t))}{t^{2/3}}=?$ What does it have to do with Gaussian random matrices? (Difficult.)

 \rightarrow What is the limit $\lim_{t\to\infty} \frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} = ?$ What does it have to do with Gaussian random matrices? (Difficult.)

 \rightarrow Other processes (zero range, *Bricklayers*', ...)?

 \rightarrow What is the limit $\lim_{t\to\infty} \frac{\operatorname{Var}(h_{C(\varrho)t}(t))}{t^{2/3}} = ?$ What does it have to do with Gaussian random matrices? (Difficult.)

→ Other processes (zero range, *Bricklayers*', ...)?

 \rightarrow Some processes (e.g. symmetric simple exclusion, linear rate zero range) show $t^{1/4}$ scaling (with Gaussian limits), rather than $t^{1/3}$. Where is the borderline? Are there other scalings as well?

Thank you.